All-versus-Nothing Arguments, Stabiliser Groups, and the "AvN Triple" Conjecture

Samson Abramsky Joint work with Kohei Kishida and Ray Lal

Department of Computer Science, University of Oxford

This style of argument was first conceptualised by Mermin.

This style of argument was first conceptualised by Mermin.

See in particular his paper "A Simple Unified Form For the Major No-Hidden-Variables Theorems" (PRL 1990)

This style of argument was first conceptualised by Mermin.

See in particular his paper "A Simple Unified Form For the Major No-Hidden-Variables Theorems" (PRL 1990)

Many papers subsequently, with many examples.

This style of argument was first conceptualised by Mermin.

See in particular his paper "A Simple Unified Form For the Major No-Hidden-Variables Theorems" (PRL 1990)

Many papers subsequently, with many examples.

However, no general definition of what an AvN argument is.

This style of argument was first conceptualised by Mermin.

See in particular his paper "A Simple Unified Form For the Major No-Hidden-Variables Theorems" (PRL 1990)

Many papers subsequently, with many examples.

However, no general definition of what an AvN argument is.

We shall provide such a definition, and formulate a conjecture of a simple characterisation of when such arguments can be made.

This style of argument was first conceptualised by Mermin.

See in particular his paper "A Simple Unified Form For the Major No-Hidden-Variables Theorems" (PRL 1990)

Many papers subsequently, with many examples.

However, no general definition of what an AvN argument is.

We shall provide such a definition, and formulate a conjecture of a simple characterisation of when such arguments can be made.

Motivation:

This style of argument was first conceptualised by Mermin.

See in particular his paper "A Simple Unified Form For the Major No-Hidden-Variables Theorems" (PRL 1990)

Many papers subsequently, with many examples.

However, no general definition of what an AvN argument is.

We shall provide such a definition, and formulate a conjecture of a simple characterisation of when such arguments can be made.

Motivation:

• Understand where AvN sits in the hierarchy of contextuality properties

This style of argument was first conceptualised by Mermin.

See in particular his paper "A Simple Unified Form For the Major No-Hidden-Variables Theorems" (PRL 1990)

Many papers subsequently, with many examples.

However, no general definition of what an AvN argument is.

We shall provide such a definition, and formulate a conjecture of a simple characterisation of when such arguments can be made.

Motivation:

- Understand where AvN sits in the hierarchy of contextuality properties
- Characterise the quantum states which give rise to maximal degrees of non-locality/contextuality.

The XOR Game



$$\mathsf{GHZ} \;=\; \frac{|\uparrow\uparrow\uparrow\rangle + \;|\downarrow\downarrow\downarrow\rangle}{\sqrt{2}}$$

$$\mathsf{GHZ} = \frac{\left|\uparrow\uparrow\uparrow\right\rangle + \left|\downarrow\downarrow\downarrow\right\rangle}{\sqrt{2}}$$

	+++	+ + -	+ - +	+	-++	-+-	+	
XXX	1	0	0	1	0	1	1	0
XYY	0	1	1	0	1	0	0	1
YXY	0	1	1	0	1	0	0	1
YYX	0	1	1	0	1	0	0	1

$$\mathsf{GHZ} \ = \ \frac{\left|\uparrow\uparrow\uparrow\right\rangle + \left|\downarrow\downarrow\downarrow\right\rangle}{\sqrt{2}}$$

	+++	+ + -	+ - +	+	-++	-+-	+	
XXX	1	0	0	1	0	1	1	0
XYY	0	1	1	0	1	0	0	1
YXY	0	1	1	0	1	0	0	1
YYX	0	1	1	0	1	0	0	1

Strongly contextual: no assignment

$$\{X_1, Y_1, X_2, Y_2, X_3, Y_3\} \longrightarrow \{+1, -1\}$$

consistent with this support.

$$\mathsf{GHZ} = \frac{|\uparrow\uparrow\uparrow\rangle + |\downarrow\downarrow\downarrow\rangle}{\sqrt{2}}$$

	+++	+ + -	+ - +	+	-++	-+-	+	
XXX	1	0	0	1	0	1	1	0
XYY	0	1	1	0	1	0	0	1
YXY	0	1	1	0	1	0	0	1
YYX	0	1	1	0	1	0	0	1

Strongly contextual: no assignment

$$\{X_1, Y_1, X_2, Y_2, X_3, Y_3\} \longrightarrow \{+1, -1\}$$

consistent with this support.

Note that the eigenvalues of the operators XXX etc. are +1 and -1.

$$\mathsf{GHZ} = \frac{|\uparrow\uparrow\uparrow\rangle + |\downarrow\downarrow\downarrow\rangle}{\sqrt{2}}$$

	+++	+ + -	+ - +	+	-++	-+-	+	
XXX	1	0	0	1	0	1	1	0
XYY	0	1	1	0	1	0	0	1
YXY	0	1	1	0	1	0	0	1
YYX	0	1	1	0	1	0	0	1

Strongly contextual: no assignment

$$\{X_1, Y_1, X_2, Y_2, X_3, Y_3\} \longrightarrow \{+1, -1\}$$

consistent with this support.

Note that the eigenvalues of the operators XXX etc. are +1 and -1.

The **expected values** of these measurements give information about the **parity** of the support.

The XYY, YXY and YYX operators all stabilise the GHZ state, *i.e.* leave it fixed.

The XYY, YXY and YYX operators all stabilise the GHZ state, *i.e.* leave it fixed.

Note that

$$\langle A \rangle_{v} = \langle v | A | v \rangle, \qquad \langle v | A | v \rangle = 1 \iff A | v \rangle = | v \rangle.$$

The XYY, YXY and YYX operators all stabilise the GHZ state, *i.e.* leave it fixed.

Note that

$$\langle A \rangle_{v} = \langle v | A | v \rangle, \qquad \langle v | A | v \rangle = 1 \iff A | v \rangle = | v \rangle.$$

Thus the expected value of measuring any of these operators on GHZ is +1.

The XYY, YXY and YYX operators all **stabilise** the GHZ state, *i.e.* leave it fixed.

Note that

$$\langle A \rangle_{v} = \langle v | A | v \rangle, \qquad \langle v | A | v \rangle = 1 \iff A | v \rangle = | v \rangle.$$

Thus the expected value of measuring any of these operators on GHZ is +1.

This says that the support of the outcomes of measuring XXX on GHZ should have **even parity**.

The *XYY*, *YXY* and *YYX* operators all **stabilise** the GHZ state, *i.e.* leave it fixed. Note that

$$\langle A \rangle_{v} = \langle v | A | v \rangle, \qquad \langle v | A | v \rangle = 1 \iff A | v \rangle = | v \rangle.$$

Thus the expected value of measuring any of these operators on GHZ is +1.

This says that the support of the outcomes of measuring XXX on GHZ should have **even parity**.

However, their product is -XXX, which also stabilises GHZ.

The *XYY*, *YXY* and *YYX* operators all **stabilise** the GHZ state, *i.e.* leave it fixed. Note that

$$\langle A \rangle_{v} = \langle v | A | v \rangle, \qquad \langle v | A | v \rangle = 1 \iff A | v \rangle = | v \rangle.$$

Thus the expected value of measuring any of these operators on GHZ is +1.

This says that the support of the outcomes of measuring XXX on GHZ should have **even parity**.

However, their product is -XXX, which also stabilises GHZ.

X_1	Y_2	Y_3	=	1
Y_1	X_2	Y_3	=	1
Y_1	Y_2	<i>X</i> ₃	=	1
X_1	X_2	X_3	=	$^{-1}$

However, this can never be the case for any assignment

$$\{X_1, Y_1, X_2, Y_2, X_3, Y_3\} \longrightarrow \{+1, -1\}$$

Use the isomorphism

$$(\{+1,-1\},\times) \cong (\{0,1\},\oplus)$$

Use the isomorphism

$$(\{+1,-1\},\times) \cong (\{0,1\},\oplus)$$

We can translate the stabilisers into parity assertions:

$$X_1 \oplus Y_2 \oplus Y_3 = 0$$

$$Y_1 \oplus X_2 \oplus Y_3 = 0$$

$$Y_1 \oplus Y_2 \oplus X_3 = 0$$

$$X_1 \oplus X_2 \oplus X_3 = 1$$

Use the isomorphism

$$(\{+1,-1\},\times) \cong (\{0,1\},\oplus)$$

We can translate the stabilisers into parity assertions:

$$X_1 \oplus Y_2 \oplus Y_3 = 0$$

$$Y_1 \oplus X_2 \oplus Y_3 = 0$$

$$Y_1 \oplus Y_2 \oplus X_3 = 0$$

$$X_1 \oplus X_2 \oplus X_3 = 1$$

Clearly, these are inconsistent.

We can define everything for general empirical models (*i.e.* "generalized probability tables") over a measurement scenario (X, \mathcal{M}) (with dichotomic measurements).

We can define everything for general empirical models (*i.e.* "generalized probability tables") over a measurement scenario (X, \mathcal{M}) (with dichotomic measurements).

To each such model e, we can associate an **XOR theory** $\mathbb{T}_{\oplus}(e)$.

We can define everything for general empirical models (*i.e.* "generalized probability tables") over a measurement scenario (X, \mathcal{M}) (with dichotomic measurements).

To each such model e, we can associate an **XOR theory** $\mathbb{T}_{\oplus}(e)$.

For each measurement context $C \in \mathcal{M}$, this will have the assertion

$$\bigoplus_{x\in C} x = 0$$

when the support of e_C is even, and

$$\bigoplus_{x\in C} x = 1$$

when the support is odd.

We can define everything for general empirical models (*i.e.* "generalized probability tables") over a measurement scenario (X, \mathcal{M}) (with dichotomic measurements).

To each such model e, we can associate an **XOR theory** $\mathbb{T}_{\oplus}(e)$.

For each measurement context $C \in \mathcal{M}$, this will have the assertion

$$\bigoplus_{x\in C} x = 0$$

when the support of e_C is even, and

$$\bigoplus_{x\in C} x = 1$$

when the support is odd.

We say that the model is AvN if this theory is inconsistent.

We can define everything for general empirical models (*i.e.* "generalized probability tables") over a measurement scenario (X, M) (with dichotomic measurements).

To each such model e, we can associate an **XOR theory** $\mathbb{T}_{\oplus}(e)$.

For each measurement context $C \in \mathcal{M}$, this will have the assertion

$$\bigoplus_{x\in C} x = 0$$

when the support of e_C is even, and

$$\bigoplus_{x\in C} x = 1$$

when the support is odd.

We say that the model is AvN if this theory is inconsistent.

Proposition

If an empirical model e is AvN, then it is strongly contextual.

To see how such AvN models can arise from quantum mechanics, we generalise Mermin's argument.

To see how such AvN models can arise from quantum mechanics, we generalise Mermin's argument.

The natural setting for this is stabilisers.

To see how such AvN models can arise from quantum mechanics, we generalise Mermin's argument.

The natural setting for this is stabilisers.

The Pauli *n*-group \mathcal{P}_n : a list of *n* Pauli operators (from $\{X, Y, Z, I\}$), with a global phase from $\{\pm 1, \pm i\}$.

To see how such AvN models can arise from quantum mechanics, we generalise Mermin's argument.

The natural setting for this is stabilisers.

The Pauli *n*-group \mathcal{P}_n : a list of *n* Pauli operators (from $\{X, Y, Z, I\}$), with a global phase from $\{\pm 1, \pm i\}$.

A Galois correspondence between Pauli operators and states/vectors in the Hilbert space \mathbb{C}^n :

$$gRv \iff gv = v.$$

Closure operators on sets of group elements and of vectors:

$$S^{\perp} := \{ v \mid \forall g \in S. gRv \}, \qquad V^{\perp} := \{ g \mid \forall v \in V. gRv \}.$$

To see how such AvN models can arise from quantum mechanics, we generalise Mermin's argument.

The natural setting for this is stabilisers.

The Pauli *n*-group \mathcal{P}_n : a list of *n* Pauli operators (from $\{X, Y, Z, I\}$), with a global phase from $\{\pm 1, \pm i\}$.

A Galois correspondence between Pauli operators and states/vectors in the Hilbert space \mathbb{C}^n :

$$gRv \iff gv = v.$$

Closure operators on sets of group elements and of vectors:

$$S^{\perp} := \{ v \mid \forall g \in S. gRv \}, \qquad V^{\perp} := \{ g \mid \forall v \in V. gRv \}.$$

The closed sets $(X = X^{\perp \perp})$ are subgroups and subspaces respectively.

To see how such AvN models can arise from quantum mechanics, we generalise Mermin's argument.

The natural setting for this is stabilisers.

The Pauli *n*-group \mathcal{P}_n : a list of *n* Pauli operators (from $\{X, Y, Z, I\}$), with a global phase from $\{\pm 1, \pm i\}$.

A Galois correspondence between Pauli operators and states/vectors in the Hilbert space \mathbb{C}^n :

$$gRv \iff gv = v.$$

Closure operators on sets of group elements and of vectors:

$$S^{\perp} := \{ v \mid \forall g \in S. gRv \}, \qquad V^{\perp} := \{ g \mid \forall v \in V. gRv \}.$$

The closed sets $(X = X^{\perp \perp})$ are subgroups and subspaces respectively.

The subgroups of \mathcal{P}_n which stabilise non-trivial subspaces must be commutative, and only contain elements with global phases ± 1 .





The subgroups are **constraints** on states: the more constraints, the fewer states satisfy them.



The subgroups are **constraints** on states: the more constraints, the fewer states satisfy them.

Akin to the Galois correspondence of theories and models in logic.



The subgroups are **constraints** on states: the more constraints, the fewer states satisfy them.

Akin to the Galois correspondence of theories and models in logic.

Note that the correspondence is tight: a rank k subgroup determines a dimension 2^{n-k} subspace.

We can associate an XOR theory $\mathbb{T}_{\oplus}(S)$ to each stabiliser subgroup S.

We can associate an XOR theory $\mathbb{T}_{\oplus}(S)$ to each stabiliser subgroup S.

For each element $P_1 \cdots P_n$ of *S*, $P_i \in \{X, Y, Z, I\}$, with global phase +1, we have the formula

$$\bigoplus_{i=1}^{n} P_i = 0$$

and for each such element with global phase -1, we have the formula

$$\bigoplus_{i=1}^n P_i = 1$$

We can associate an XOR theory $\mathbb{T}_{\oplus}(S)$ to each stabiliser subgroup S.

For each element $P_1 \cdots P_n$ of *S*, $P_i \in \{X, Y, Z, I\}$, with global phase +1, we have the formula

$$\bigoplus_{i=1}^{n} P_i = 0$$

and for each such element with global phase -1, we have the formula

$$\bigoplus_{i=1}^n P_i = 1$$

We say that S is AvN if $\mathbb{T}_{\oplus}(S)$ is inconsistent.

We can associate an XOR theory $\mathbb{T}_{\oplus}(S)$ to each stabiliser subgroup S.

For each element $P_1 \cdots P_n$ of *S*, $P_i \in \{X, Y, Z, I\}$, with global phase +1, we have the formula

$$\bigoplus_{i=1}^{n} P_i = 0$$

and for each such element with global phase -1, we have the formula

$$\bigoplus_{i=1}^n P_i = 1$$

We say that S is AvN if $\mathbb{T}_{\oplus}(S)$ is inconsistent.

Question:

How can we characterise when this happens?

Define an AvN **triple** in \mathcal{P}_n to be (e, f, g) (order is important) with global phases +1, which pairwise commute, and additionally satisfy the following conditions:

- (A1) For all i = 1, ..., n at least two of e_i , f_i , g_i are the same.
- (A2) The number of *i* such that $e_i = g_i \neq f_i$, all distinct from *I*, is odd.

Define an AvN **triple** in \mathcal{P}_n to be (e, f, g) (order is important) with global phases +1, which pairwise commute, and additionally satisfy the following conditions:

- (A1) For all i = 1, ..., n at least two of e_i , f_i , g_i are the same.
- (A2) The number of *i* such that $e_i = g_i \neq f_i$, all distinct from *I*, is odd.

So in (A2) these are triples PQP of Pauli matrices, all distinct from I, $Q \neq P$.

Define an AvN **triple** in \mathcal{P}_n to be (e, f, g) (order is important) with global phases +1, which pairwise commute, and additionally satisfy the following conditions:

- (A1) For all i = 1, ..., n at least two of e_i , f_i , g_i are the same.
- (A2) The number of *i* such that $e_i = g_i \neq f_i$, all distinct from *I*, is odd.

So in (A2) these are triples PQP of Pauli matrices, all distinct from I, $Q \neq P$. Now the claim is that such a triple yields an AvN argument.

Define an AvN **triple** in \mathcal{P}_n to be (e, f, g) (order is important) with global phases +1, which pairwise commute, and additionally satisfy the following conditions:

- (A1) For all i = 1, ..., n at least two of e_i , f_i , g_i are the same.
- (A2) The number of *i* such that $e_i = g_i \neq f_i$, all distinct from *I*, is odd.

So in (A2) these are triples PQP of Pauli matrices, all distinct from I, $Q \neq P$. Now the claim is that such a triple yields an AvN argument.

Note that the conditions imply that the product e.f.g = -h, which translates into a condition of odd parity on the support of any state stabilised by these operators for the measurement h.

Define an AvN **triple** in \mathcal{P}_n to be (e, f, g) (order is important) with global phases +1, which pairwise commute, and additionally satisfy the following conditions:

- (A1) For all i = 1, ..., n at least two of e_i , f_i , g_i are the same.
- (A2) The number of *i* such that $e_i = g_i \neq f_i$, all distinct from *I*, is odd.

So in (A2) these are triples PQP of Pauli matrices, all distinct from I, $Q \neq P$. Now the claim is that such a triple yields an AvN argument.

Note that the conditions imply that the product e.f.g = -h, which translates into a condition of odd parity on the support of any state stabilised by these operators for the measurement h.

On the other hand, condition (A1) implies that under any global assignment/section on the variables, we can cancel the repeated items in each column, and deduce an even parity for h.

Define an AvN **triple** in \mathcal{P}_n to be (e, f, g) (order is important) with global phases +1, which pairwise commute, and additionally satisfy the following conditions:

- (A1) For all i = 1, ..., n at least two of e_i , f_i , g_i are the same.
- (A2) The number of *i* such that $e_i = g_i \neq f_i$, all distinct from *I*, is odd.

So in (A2) these are triples PQP of Pauli matrices, all distinct from I, $Q \neq P$. Now the claim is that such a triple yields an AvN argument.

- Note that the conditions imply that the product e.f.g = -h, which translates into a condition of odd parity on the support of any state stabilised by these operators for the measurement h.
- On the other hand, condition (A1) implies that under any global assignment/section on the variables, we can cancel the repeated items in each column, and deduce an even parity for h.
- This means that any state in V_S , where S is the subgroup generated by $\{e, f, g\}$, admits an AvN argument. Note that this is a 2^{n-3} -dimensional space, assuming e, f, g are independent.

The further conjecture is that having an AvN triple is **necessary** as well as sufficient for an AvN argument.

The further conjecture is that having an AvN triple is **necessary** as well as sufficient for an AvN argument.

More precisely, any AvN subgroup S must contain an AvN triple.

The further conjecture is that having an AvN triple is **necessary** as well as sufficient for an AvN argument.

More precisely, any AvN subgroup S must contain an AvN triple.

Example from Mermin, yielding a GHZ argument:

X Y Y Y X Y Y Y X

Example of 1-dimensional cluster state, n = 4: