# Reality of the Quantum State: A Stronger $\psi$ -ontology Theorem

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# The Quantum State $\psi$ — Real or Phenomenal?

- Assume some space  $\Lambda$  of *ontic states*
- Preparation of *quantum states* ψ, φ ∈ ℋ induce probability distributions μ<sub>ψ</sub>, μ<sub>φ</sub> over Λ, etc.



- If distributions can overlap  $\rightarrow \psi$ -epistemic
- If distributions never overlap  $\rightarrow$ Each  $\lambda \in \Lambda$  encodes a *unique* quantum state, so  $\psi$ -ontic

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## The PBR Theorem\*

The following assumptions

- 1. systems have an objective physical state
- 2. quantum predictions are correct
- 3. preparation independence

imply  $\psi$ -ontic.





\*Pusey, Barrett & Rudolph, arXiv:1111.3328 [quant-ph]

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### **Preparation Independence**

The only reasonable option?



 $\mu(\lambda_A, \lambda_B \mid p_A, p_B) = \mu(\lambda_A \mid p_A) \times \mu(\lambda_B \mid p_B)$ 

## Comparison with Bell Locality

#### An intuitive notion in measurement scenarios



 $p(o_A, o_B \mid m_A, m_B, \lambda) = p(o_A \mid m_A, \lambda) \times p(o_A \mid m_B, \lambda)$ 

(Ruled out by Bell's Theorem)

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### Weakening Preparation Independence

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# No-signalling



$$p(o_A \mid m_A, m_B) = p(o_A \mid m_A)$$
$$p(o_B \mid m_A, m_B) = p(o_B \mid m_B)$$

- Allows good notion of subsystem
- Consistent with SR

## An Alternative to Preparation Independence

*Idea:* make the minimum assumption that will allow a reasonable notion of subsystem



$$\mu(\lambda_A \mid p_A, p_B) = \mu(\lambda_A \mid p_A)$$
$$\mu(\lambda_B \mid p_A, p_B) = \mu(\lambda_B \mid p_B)$$

#### **Escaping PBR's Conclusion**

#### A $\psi$ -epistemic model realising PBR statistics:

Define  $\mu_{00}, \mu_{0+}, \mu_{+0}, \mu_{++}$  by the table below and measurement response functions as on the right



$$\begin{split} \boldsymbol{\xi}_{\neg(\boldsymbol{\psi},\boldsymbol{\psi})}(\boldsymbol{\lambda}) &:= \begin{cases} 1 & \text{if } \boldsymbol{\lambda} = (\lambda_{\boldsymbol{\phi}-\boldsymbol{\psi}},\lambda_{\boldsymbol{\phi}-\boldsymbol{\psi}}) \\ 1/2 & \text{if } \boldsymbol{\lambda} \in \{(\lambda_{\boldsymbol{\psi}\wedge\boldsymbol{\phi}},\lambda_{\boldsymbol{\phi}-\boldsymbol{\psi}}),(\lambda_{\boldsymbol{\phi}-\boldsymbol{\psi}},\lambda_{\boldsymbol{\psi}\wedge\boldsymbol{\phi}})\} \\ 0 & \text{otherwise} \end{cases} \\ \boldsymbol{\xi}_{\neg(\boldsymbol{\psi},\boldsymbol{\phi})}(\boldsymbol{\lambda}) &:= \begin{cases} 1 & \text{if } \boldsymbol{\lambda} = (\lambda_{\boldsymbol{\phi}-\boldsymbol{\psi}},\lambda_{\boldsymbol{\psi}-\boldsymbol{\phi}}) \\ 1/2 & \text{if } \boldsymbol{\lambda} \in \{(\lambda_{\boldsymbol{\psi}\wedge\boldsymbol{\phi}},\lambda_{\boldsymbol{\psi}-\boldsymbol{\phi}}),(\lambda_{\boldsymbol{\phi}-\boldsymbol{\psi}},\lambda_{\boldsymbol{\psi}\wedge\boldsymbol{\phi}})\}\} \\ 0 & \text{otherwise} \end{cases} \\ \boldsymbol{\xi}_{\neg(\boldsymbol{\phi},\boldsymbol{\psi})}(\boldsymbol{\lambda}) &:= \begin{cases} 1 & \text{if } \boldsymbol{\lambda} = (\lambda_{\boldsymbol{\psi}-\boldsymbol{\phi}},\lambda_{\boldsymbol{\phi}-\boldsymbol{\psi}}) \\ 1/2 & \text{if } \boldsymbol{\lambda} \in \{(\lambda_{\boldsymbol{\psi}-\boldsymbol{\phi}},\lambda_{\boldsymbol{\psi}\wedge\boldsymbol{\phi}}),(\lambda_{\boldsymbol{\psi}\wedge\boldsymbol{\phi}},\lambda_{\boldsymbol{\phi}-\boldsymbol{\psi}})\}\} \\ 0 & \text{otherwise} \end{cases} \\ \boldsymbol{\xi}_{\neg(\boldsymbol{\phi},\boldsymbol{\phi})}(\boldsymbol{\lambda}) &:= \begin{cases} 1 & \text{if } \boldsymbol{\lambda} = (\lambda_{\boldsymbol{\psi}-\boldsymbol{\phi}},\lambda_{\boldsymbol{\psi}\wedge\boldsymbol{\phi}}),(\lambda_{\boldsymbol{\psi}\wedge\boldsymbol{\phi}},\lambda_{\boldsymbol{\phi}-\boldsymbol{\psi}})\} \\ 1/2 & \text{if } \boldsymbol{\lambda} \in \{(\lambda_{\boldsymbol{\psi}-\boldsymbol{\phi}},\lambda_{\boldsymbol{\psi}\wedge\boldsymbol{\phi}}),(\lambda_{\boldsymbol{\psi}\wedge\boldsymbol{\phi}},\lambda_{\boldsymbol{\psi}-\boldsymbol{\phi}})\}\} \\ 0 & \text{otherwise} \end{cases} \end{cases}$$

## The Crux of the Matter (Outline)

One step of the PBR argument:

- Suppose μ<sub>ψ</sub> and μ<sub>φ</sub> have overlapping supports, such that a device that can prepare p ∈ {ψ, φ} results in an ontic state in the overlap region with probability at least q
- Then with two such preparation devices, there is probability at least  $q^2$  that *both* ontic states lie in the overlap region



This does not hold under our weakened notion of independence.

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## A Stronger $\psi$ -ontology Theorem

• Uniformly sample *n* preparations from *m* preparation devices



• Even allowing for the loophole in PBR, we prove a bound on the probability of being in the overlap region

$$q \leq \frac{1}{\left\lceil \frac{n}{m-1} \right\rceil}$$

• 
$$q \rightarrow 0$$
 as  $m \rightarrow \infty$ 

# Conclusion

- Take a reasonable weaker notion of independence
- $\psi$  can be interpreted statistically

#### BUT!!!

- The degree to which systems may be composed limits the degree to which  $\psi$  may be statistical
- Fine for simple toy theories, but not for fully fledged physical theories
- $\psi$  is still real!

# Appendix: The Quantum State $\psi$ — Real or Phenomenal?

#### $\psi$ -ontic:

- A real physical wave (on configuration space?)
- Easiest way to think about interference
- PBR theorem



#### $\psi$ -epistemic:

- $\psi$  gives probabilistic information
- Collapse  $\rightarrow$  Bayesian updating
- Can't reliably distinguish non-orthogonal  $\psi, \phi$
- *ψ* is exponential in the number of systems
- Can't be cloned
- Can be teleported