(In)Completeness results for the zx-calculus

Miriam Backens

Department of Computer Science, University of Oxford

10 Years of Categorical Quantum Mechanics, Oxford, 2014

< □ > < 同 > < 三 > < 三 > < 三 > < ○ < ○ </p>

Outline

Introduction

The zx-calculus is incomplete [Schröder & Zamdzhiev 2014]

◆□▶ ◆□▶ ▲□▶ ▲□▶ □ のQ@

Completeness results

Stabilizer quantum mechanics The single-qubit Clifford+T group

Outlook & Conclusions

Outline

Introduction

The zx-calculus is incomplete [Schröder & Zamdzhiev 2014]

◆□▶ ◆□▶ ▲□▶ ▲□▶ □ のQ@

Completeness results

Stabilizer quantum mechanics The single-qubit Clifford+T group

Outlook & Conclusions

Elements of zx-calculus diagrams

• green nodes with *n* inputs and *m* outputs, $\alpha \in (-\pi, \pi]$



▶ red nodes with *n* inputs and *m* outputs, $\beta \in (-\pi, \pi]$



Hadamard nodes with one input and one output

・ロト ・ 同 ・ ・ ヨ ・ ・ ヨ ・ うへつ

Elements of zx-calculus diagrams

• green nodes with *n* inputs and *m* outputs, $\alpha \in (-\pi, \pi]$

$$\left[\begin{array}{c} m\\ \hline \cdots\\ \hline \alpha\\ \hline n\end{array}\right] := |0\rangle^{\otimes m} \langle 0|^{\otimes n} + e^{i\alpha} |1\rangle^{\otimes m} \langle 1|^{\otimes n},$$

▶ red nodes with *n* inputs and *m* outputs, $\beta \in (-\pi, \pi]$

$$\left[\begin{array}{c} \underbrace{m} \\ \vdots \\ \vdots \\ n \end{array}\right] := |+\rangle^{\otimes m} \langle +|^{\otimes n} + e^{i\beta} |-\rangle^{\otimes m} \langle -|^{\otimes n},$$

Hadamard nodes with one input and one output

$$\left[\begin{array}{c} \mathbf{\underline{\mu}} \end{array} \right] := \left| + \right\rangle \left\langle \mathbf{0} \right| + \left| - \right\rangle \left\langle \mathbf{1} \right|$$

Properties of the interpretation functor

parallel composition:

serial composition:



to get the Hermitian adjoint of a diagram, turn it upside-down and flip all the phases: e.g.



a phase label of zero is usually left out: e.g. • := • 0

▲□▶ ▲□▶ ▲□▶ ▲□▶ ▲□ ● のへぐ

- a phase label of zero is usually left out: e.g. := 0
- diagrams with no inputs are states: e.g.

- a phase label of zero is usually left out: e.g. := 0
- diagrams with no inputs are states: e.g.

・ロト・日本・日本・日本・日本

▶ diagrams with no outputs are *effects*: e.g. ♥ *m*

- a phase label of zero is usually left out: e.g. $\mathbf{b} := \mathbf{b} \mathbf{0}$
- diagrams with no inputs are states: e.g.
- diagrams with no outputs are *effects*: e.g. $\phi\pi$
- diagrams with no inputs or outputs are scalars: e.g. $\frac{1}{2}\pi/2$

・ロト・日本・日本・日本・日本

- a phase label of zero is usually left out: e.g. := 0
- diagrams with no inputs are states: e.g.
- ▶ diagrams with no outputs are *effects*: e.g. ♥ *m*
- diagrams with no inputs or outputs are scalars: e.g. $\frac{1}{2}\pi/2$
- ► a red or green node with one input and one output is also called *phase shift*: e.g. φ -π/2

▲□▶▲□▶▲□▶▲□▶ □ のへで

Rules of the zx-calculus

- ignore non-zero scalar factors $(\llbracket \bullet \pi \rrbracket \doteq 0 \doteq \llbracket \bullet \pi \rrbracket)$
- only the topology matters



Universality, soundness, and completeness

universality: any pure state, post-selected pure projective measurement, and unitary operation can be expressed: to see this note that [[] = |0⟩, [[] = ⟨0|,

$$\llbracket \phi \alpha \rrbracket \doteq \begin{pmatrix} 1 & 0 \\ 0 & e^{i\alpha} \end{pmatrix}, \quad \llbracket \phi - \phi \rrbracket \doteq \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 0 & 1 \\ 0 & 0 & 1 & 0 \end{pmatrix}$$

・ロト ・ 同 ・ ・ ヨ ・ ・ ヨ ・ うへつ

Universality, soundness, and completeness

universality: any pure state, post-selected pure projective measurement, and unitary operation can be expressed: to see this note that []●] = |0⟩, [[●]] = ⟨0|,

$$\llbracket \phi \alpha \rrbracket \doteq \begin{pmatrix} 1 & 0 \\ 0 & e^{i\alpha} \end{pmatrix}, \quad \llbracket \phi - \phi \rrbracket \doteq \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 0 & 1 \\ 0 & 0 & 1 & 0 \end{pmatrix}$$

 soundness: any equality that can be derived graphically can also be derived using matrices; to see this note that

 $[\![LHS]\!] \doteq [\![RHS]\!] \quad \text{for all rewrite rules}$

・ロト ・ 同 ・ ・ ヨ ・ ・ ヨ ・ うへつ

Universality, soundness, and completeness

universality: any pure state, post-selected pure projective measurement, and unitary operation can be expressed: to see this note that []●] = |0⟩, [[●]] = ⟨0|,

$$\llbracket \phi \alpha \rrbracket \doteq \begin{pmatrix} 1 & 0 \\ 0 & e^{i\alpha} \end{pmatrix}, \quad \llbracket \phi - \phi \rrbracket \doteq \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 0 & 1 \\ 0 & 0 & 1 & 0 \end{pmatrix}$$

 soundness: any equality that can be derived graphically can also be derived using matrices; to see this note that

 $\llbracket LHS \rrbracket \doteq \llbracket RHS \rrbracket$ for all rewrite rules

 completeness: can any equality that can be derived using matrices also be derived graphically, i.e.

does $\llbracket D_1 \rrbracket \doteq \llbracket D_2 \rrbracket$ imply $D_1 = D_2$?

Outline

Introduction

The zx-calculus is incomplete [Schröder & Zamdzhiev 2014]

◆□▶ ◆□▶ ▲□▶ ▲□▶ □ のQ@

Completeness results

Stabilizer quantum mechanics The single-qubit Clifford+T group

Outlook & Conclusions

Euler decomposition of arbitrary unitaries

Any 2x2 unitary U can be written as

$$\boldsymbol{U} \doteq \begin{pmatrix} \boldsymbol{1} & \boldsymbol{0} \\ \boldsymbol{0} & \boldsymbol{e}^{i\alpha} \end{pmatrix} \begin{pmatrix} \cos\beta & -i\sin\beta \\ -i\sin\beta & \cos\beta \end{pmatrix} \begin{pmatrix} \boldsymbol{1} & \boldsymbol{0} \\ \boldsymbol{0} & \boldsymbol{e}^{i\gamma} \end{pmatrix}$$

So if the zx-calculus is complete, then for any U we have to be able to find α, β, γ such that

$$\boxed{U} = \phi \alpha \\ \phi \beta \\ \phi \gamma$$

・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・

Alternative interpretation functors

• define $\llbracket - \rrbracket_k$ by $\llbracket \stackrel{\bullet}{\mathbf{H}} \rrbracket_k := \llbracket \stackrel{\bullet}{\mathbf{H}} \rrbracket$ and



◆□▶ ◆□▶ ▲□▶ ▲□▶ ■ ののの

▶ if k is an odd integer, the interpretation [[-]]_k is sound



 By Euler decomposition of arbitrary unitaries, can find α, β, γ such that $\llbracket D_1 \rrbracket = \llbracket D_2 \rrbracket$

◆□▶ ◆□▶ ▲□▶ ▲□▶ ■ ののの



► By Euler decomposition of arbitrary unitaries, can find α, β, γ such that $\left[\begin{bmatrix} D_1 \\ D_1 \end{bmatrix} \right] \doteq \left[\begin{bmatrix} D_2 \\ D_2 \end{bmatrix} \right]$

◆□▶ ◆□▶ ▲□▶ ▲□▶ ■ ののの

• Thus if zx-calculus is complete, $\vec{D}_1 = \vec{D}_2$



 By Euler decomposition of arbitrary unitaries, can find α, β, γ such that $\llbracket D_1 \rrbracket \doteq \llbracket D_2 \rrbracket$

◆□▶ ◆□▶ ▲□▶ ▲□▶ ■ ののの

- Thus if zx-calculus is complete, $\vec{D}_1 = \vec{D}_2$
- ▶ Then must also have $\left\| \begin{bmatrix} D_1 \\ D_1 \end{bmatrix} \right\|_{2} \doteq \left\| \begin{bmatrix} D_2 \\ D_2 \end{bmatrix} \right\|_{2}$



- By Euler decomposition of arbitrary unitaries, can find α, β, γ such that $\llbracket D_1 \rrbracket \doteq \llbracket D_2 \rrbracket$
- Thus if zx-calculus is complete, $\vec{D}_1 = \vec{D}_2$
- ▶ Then must also have $\left\| \begin{bmatrix} D_1 \\ D_1 \end{bmatrix} \right\|_{2} \doteq \left\| \begin{bmatrix} D_2 \\ D_2 \end{bmatrix} \right\|_{2}$

• But
$$\left[\left[\overrightarrow{D_1}\right]\right]_{-3} \doteq I$$
 and $\left[\left[\left[\overrightarrow{D_2}\right]\right]_{-3}$ is non-trivial \implies contradiction!

Further notes on incompleteness

- counterexample involves only line graphs, i.e. very simple structure
- can find similar counterexamples for other types of graphs
- might have to add infinitely many new rewrite rules to complete the general zx-calculus

But:

 Euler decomposition of general unitaries relies on arbitrary phases being allowed

 can get completeness results for current ruleset by restricting phases

Outline

Introduction

The zx-calculus is incomplete [Schröder & Zamdzhiev 2014]

◆□▶ ◆□▶ ▲□▶ ▲□▶ ■ ののの

Completeness results

Stabilizer quantum mechanics The single-qubit Clifford+T group

Outlook & Conclusions



Introduction

The zx-calculus is incomplete [Schröder & Zamdzhiev 2014]

◆□▶ ◆□▶ ▲□▶ ▲□▶ ■ ののの

Completeness results Stabilizer quantum mechanics The single-qubit Clifford+T group

Outlook & Conclusions

Stabilizer quantum mechanics

Stabilizer operations:

- preparation of qubits in state $|0\rangle$
- Clifford unitaries, generated by

$$S = \begin{pmatrix} 1 & 0 \\ 0 & i \end{pmatrix}, \ H = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 & 1 \\ 1 & -1 \end{pmatrix}, \ \Lambda X = \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 0 & 1 \\ 0 & 0 & 1 & 0 \end{pmatrix}$$

◆□▶ ◆□▶ ▲□▶ ▲□▶ □ のQ@

measurements in computational basis

Stabilizer quantum mechanics

Stabilizer operations:

- \blacktriangleright preparation of qubits in state $|0\rangle$
- Clifford unitaries, generated by

$$S = \begin{pmatrix} 1 & 0 \\ 0 & i \end{pmatrix}, \ H = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 & 1 \\ 1 & -1 \end{pmatrix}, \ \Lambda X = \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 0 & 1 \\ 0 & 0 & 1 & 0 \end{pmatrix}$$

measurements in computational basis

ZX-calculus: $S \doteq \llbracket \phi \pi/2 \rrbracket$, $H \doteq \llbracket \mu \rrbracket$, $\Lambda X \doteq \llbracket \phi - \phi \rrbracket$ i.e. need to restrict phases to integer multiples of $\pi/2$

$$\begin{array}{ccc}
\underline{m} & \underline{m} \\
\underline{n} & \underline{\beta} \\
\underline{n} & \underline{n}
\end{array}$$
where $\alpha, \beta \in \{-\pi/2, 0, \pi/2, \pi\}$

◆□▶ ◆□▶ ▲□▶ ▲□▶ ■ ののの

Graph states in the zx-calculus

Definition

Let G be a finite simple undirected graph. The zx-calculus diagram for the corresponding graph state consists of:

- ▶ for each node in *G*, a green node with one output, and
- ▶ for each edge in *G*, an edge with a Hadamard node on it.

◆□▶ ◆□▶ ▲□▶ ▲□▶ ■ ののの





Graph states in the zx-calculus

Definition

Let *G* be a finite simple undirected graph. The zx-calculus diagram for the corresponding graph state consists of:

- ▶ for each node in *G*, a green node with one output, and
- for each edge in G, an edge with a Hadamard node on it.



・ コット (雪) (小田) (コット 日)

Graph states in the zx-calculus

Definition

Let *G* be a finite simple undirected graph. The zx-calculus diagram for the corresponding graph state consists of:

- ▶ for each node in *G*, a green node with one output, and
- ▶ for each edge in *G*, an edge with a Hadamard node on it.





・ロット (雪) ・ (日) ・ (日)

The local Clifford group consists of all tensor products of the single-qubit Clifford operators $\langle S, H \rangle$.

Theorem (Van den Nest et. al, 2004)

Any stabilizer state is local Clifford-equivalent to some graph state.

< □ > < 同 > < 三 > < 三 > < 三 > < ○ < ○ </p>

The local Clifford group consists of all tensor products of the single-qubit Clifford operators $\langle S, H \rangle$.

Theorem (Van den Nest et. al, 2004)

Any stabilizer state is local Clifford-equivalent to some graph state.

Theorem (Van den Nest et. al, 2004)

Two graph states are local Clifford-equivalent if and only if they are related by a sequence of local complementations.

A local complementation about a vertex v inverts the subgraph generated by the neighbourhood of v: e.g.

(ロ) (同) (三) (三) (三) (○) (○)



The local Clifford group consists of all tensor products of the single-qubit Clifford operators $\langle S, H \rangle$.

Theorem (Van den Nest et. al, 2004)

Any stabilizer state is local Clifford-equivalent to some graph state.

Theorem (Van den Nest et. al, 2004)

Two graph states are local Clifford-equivalent if and only if they are related by a sequence of local complementations.

A local complementation about a vertex v inverts the subgraph generated by the neighbourhood of v: e.g.

・ロト ・ 同 ・ ・ ヨ ・ ・ ヨ ・ うへつ



The local Clifford group consists of all tensor products of the single-qubit Clifford operators $\langle S, H \rangle$.

Theorem (Van den Nest et. al, 2004)

Any stabilizer state is local Clifford-equivalent to some graph state.

Theorem (Van den Nest et. al, 2004)

Two graph states are local Clifford-equivalent if and only if they are related by a sequence of local complementations.

A local complementation about a vertex v inverts the subgraph generated by the neighbourhood of v: e.g.



◆□▶ ◆□▶ ▲□▶ ▲□▶ ■ ののの

Local complementations in the zx-calculus

Theorem (Duncan & Perdrix, 2009)

Denote the result of a local complementation about the vertex v in the graph G by $G \star v$. The graph state diagrams for G and $G \star v$ satisfy



where $\alpha_u = -\pi/2$ if $\{u, v\}$ is an edge, $\alpha_u = 0$ otherwise.

E.g. a local complementation about the 3rd qubit of the previous graph state:



Local complementations in the zx-calculus

Theorem (Duncan & Perdrix, 2009)

Denote the result of a local complementation about the vertex v in the graph G by $G \star v$. The graph state diagrams for G and $G \star v$ satisfy



where $\alpha_u = -\pi/2$ if $\{u, v\}$ is an edge, $\alpha_u = 0$ otherwise.

E.g. a local complementation about the 3rd qubit of the previous graph state:



GS-LC diagrams

Definition

A diagram in the stabilizer zx-calculus is called a *GS-LC diagram* if it consists of a graph state diagram with arbitrary single-qubit Clifford unitaries (i.e. combinations of phase shifts and Hadamards) applied to each output.

E.g.



◆□▶ ◆□▶ ▲□▶ ▲□▶ ■ ののの

Every state diagram is equal to some GS-LC diagram

Theorem

Any stabilizer ZX-calculus diagram with no inputs and at least one output can be rewritten to a GS-LC diagram.

Proof.

Decompose the diagram into basic spiders

and single-qubit Clifford unitaries.

- For each basic element, applying it to a GS-LC diagram yields a diagram that can be rewritten into GS-LC form.
- Thus, by induction, the theorem holds.

Reduced GS-LC diagrams

Definition

A diagram in the zx-calculus is called a *reduced GS-LC diagram* if it is a GS-LC diagram and satisfies the following two conditions:

1. All the single-qubit Clifford unitaries belong to the set

$$\left| \begin{array}{c} \bullet \pi/2 \\ \bullet \pi/2 \\ \bullet \pi \end{array} \right| \left| \begin{array}{c} \bullet \pi/2 \\ \bullet \pi/2 \\ \bullet \pi/2 \\ \bullet \pi/2 \\ \bullet \pi/2 \end{array} \right| \left| \begin{array}{c} \bullet \pi/2 \\ \bullet \pi/2 \\ \bullet \pi/2 \\ \bullet \pi/2 \end{array} \right| \left| \begin{array}{c} \bullet \pi/2 \\ \bullet \pi/2 \\ \bullet \pi/2 \\ \bullet \pi/2 \end{array} \right| \left| \begin{array}{c} \bullet \pi/2 \\ \bullet \pi/2 \\ \bullet \pi/2 \\ \bullet \pi/2 \end{array} \right| \left| \begin{array}{c} \bullet \pi/2 \\ \bullet \pi/2 \\ \bullet \pi/2 \\ \bullet \pi/2 \end{array} \right| \left| \begin{array}{c} \bullet \pi/2 \\ \bullet \pi/2 \\ \bullet \pi/2 \\ \bullet \pi/2 \end{array} \right| \left| \begin{array}{c} \bullet \pi/2 \\ \bullet \pi/2 \\ \bullet \pi/2 \\ \bullet \pi/2 \end{array} \right| \left| \begin{array}{c} \bullet \pi/2 \\ \bullet \pi/2 \\ \bullet \pi/2 \\ \bullet \pi/2 \end{array} \right| \left| \begin{array}{c} \bullet \pi/2 \\ \bullet \pi/2 \\ \bullet \pi/2 \\ \bullet \pi/2 \end{array} \right| \left| \begin{array}{c} \bullet \pi/2 \\ \bullet \pi/2 \\ \bullet \pi/2 \\ \bullet \pi/2 \end{array} \right| \left| \begin{array}{c} \bullet \pi/2 \\ \bullet \pi/2 \\ \bullet \pi/2 \\ \bullet \pi/2 \end{array} \right| \left| \begin{array}{c} \bullet \pi/2 \\ \bullet \pi/2 \\ \bullet \pi/2 \\ \bullet \pi/2 \end{array} \right| \left| \begin{array}{c} \bullet \pi/2 \\ \bullet \pi/2 \\ \bullet \pi/2 \\ \bullet \pi/2 \end{array} \right| \left| \begin{array}{c} \bullet \pi/2 \\ \bullet \pi/2 \\ \bullet \pi/2 \\ \bullet \pi/2 \end{array} \right| \left| \begin{array}{c} \bullet \pi/2 \\ \bullet \pi/2 \\ \bullet \pi/2 \\ \bullet \pi/2 \end{array} \right| \left| \begin{array}{c} \bullet \pi/2 \\ \bullet \pi/2 \\ \bullet \pi/2 \\ \bullet \pi/2 \end{array} \right| \left| \begin{array}{c} \bullet \pi/2 \\ \bullet \pi/2 \\ \bullet \pi/2 \\ \bullet \pi/2 \end{array} \right| \left| \begin{array}{c} \bullet \pi/2 \\ \bullet \pi/2 \\ \bullet \pi/2 \\ \bullet \pi/2 \end{array} \right| \left| \begin{array}{c} \bullet \pi/2 \\ \bullet \pi/2 \\ \bullet \pi/2 \\ \bullet \pi/2 \end{array} \right| \left| \begin{array}{c} \bullet \pi/2 \\ \bullet \pi/2 \\ \bullet \pi/2 \\ \bullet \pi/2 \end{array} \right| \left| \begin{array}{c} \bullet \pi/2 \\ \bullet \pi/2 \\ \bullet \pi/2 \\ \bullet \pi/2 \end{array} \right| \left| \begin{array}{c} \bullet \pi/2 \\ \bullet \pi/2 \\ \bullet \pi/2 \end{array} \right| \left| \begin{array}{c} \bullet \pi/2 \\ \bullet \pi/2 \\ \bullet \pi/2 \\ \bullet \pi/2 \end{array} \right| \left| \begin{array}{c} \bullet \pi/2 \\ \bullet \pi/2 \\ \bullet \pi/2 \end{array} \right| \left| \begin{array}{c} \bullet \pi/2 \\ \bullet \pi/2 \\ \bullet \pi/2 \\ \bullet \pi/2 \end{array} \right| \left| \begin{array}{c} \bullet \pi/2 \\ \bullet \pi/2 \\ \bullet \pi/2 \\ \bullet \pi/2 \end{array} \right| \left| \begin{array}{c} \bullet \pi/2 \\ \bullet \pi/2 \\ \bullet \pi/2 \\ \bullet \pi/2 \\ \bullet \pi/2 \end{array} \right| \left| \begin{array}{c} \bullet \pi/2 \\ \bullet$$

2. Two adjacent vertices must not both have single-qubit Clifford unitaries that include red nodes.

Theorem

Any stabilizer ZX-calculus state diagram can be rewritten to a reduced GS-LC diagram.

Comparing reduced GS-LC diagrams

Theorem (inspired by Elliott et al., 2008)

There exists a terminating algorithm that, given a pair of reduced GS-LC diagrams on the same number of qubits, rewrites them to a pair of identical diagrams if and only if the two diagrams represent the same state.

Theorem

The zx-calculus is complete for stabilizer state diagrams.

・ロト ・ 同 ・ ・ ヨ ・ ・ ヨ ・ うへつ

The Choi-Jamiołkowski isomorphism

Theorem (Choi-Jamiołkowski isomorphism)

For any operator A from n to m qubits and for any n + m-qubit state B,



◆□▶ ◆□▶ ▲□▶ ▲□▶ ■ ののの

The stabilizer zx-calculus is complete

Theorem

The zx-calculus is complete for all stabilizer diagrams.

Proof.

Given two zx-calculus diagrams with *n* inputs and *m* outputs each:

- Apply the Choi-Jamiołkowski isomorphism to get two diagrams with n + m outputs each.
- Bring the diagrams into reduced GS-LC form.
- Apply the comparison algorithm for reduced GS-LC diagrams.
- If this yields a pair of identical diagrams, use the Choi-Jamiołkowski isomorphism to transform the sequence of equal state diagrams back into operators.

This yields a sequence of rewrites which transforms one of the original diagrams into the other.

Outline

Introduction

The zx-calculus is incomplete [Schröder & Zamdzhiev 2014]

◆□▶ ◆□▶ ▲□▶ ▲□▶ ■ ののの

Completeness results Stabilizer quantum mechanics The single-qubit Clifford+T group

Outlook & Conclusions

The single-qubit Clifford+T group

An approximately universal group, generated by:

• single-qubit Clifford group $C_1 = \langle S, H \rangle$, where

$$S = \begin{pmatrix} 1 & 0 \\ 0 & i \end{pmatrix}$$
 and $H = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 & 1 \\ 1 & -1 \end{pmatrix}$

• T gate (note $T^2 = S$)

$$T = \begin{pmatrix} 1 & 0 \\ 0 & e^{i\pi/4} \end{pmatrix}$$

・ロト ・ 同 ・ ・ ヨ ・ ・ ヨ ・ うへつ

The single-qubit Clifford+T group

An approximately universal group, generated by:

• single-qubit Clifford group $C_1 = \langle S, H \rangle$, where

$$S = \begin{pmatrix} 1 & 0 \\ 0 & i \end{pmatrix}$$
 and $H = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 & 1 \\ 1 & -1 \end{pmatrix}$

T gate (note *T*² = *S*)

$$T=egin{pmatrix} 1 & 0 \ 0 & e^{i\pi/4} \end{pmatrix}$$

zx-calculus: $T \doteq \llbracket \phi \pi/4 \rrbracket$ and $H \doteq \llbracket \mu \rrbracket$

- diagrams are restricted to line graphs (each node has one input and one output)
- phases are restricted to integer multiples of $\pi/4$

The normal form for single-qubit Clifford+T diagrams

Following [Matsumoto & Amano 2008], any single-qubit Clifford+T diagram is either pure Clifford or it can be written as

$$\begin{array}{c}
\overbrace{V_{1}}{V_{n}} \\
\overbrace{V_{1}}{V_{1}} \\
\overbrace{V}{V_{1}} \\
\overbrace{V}{V} \\
\overbrace{V}{V_{1}} \\
\overbrace{V}$$

with *n* a non-negative integer and $\alpha, \beta, \gamma \in \{0, \pi/2, \pi, -\pi/2\}$.

Rewriting diagrams into normal form

- write diagram in terms of φ α where α is a multiple of π/4, and φ β with β a multiple of π/2
- diagrams are rewritten into normal form by pushing phase shift towards the bottom of the diagram
- ▶ we say we can "push A past B" if there is A' such that



・ロト ・ 同 ・ ・ ヨ ・ ・ ヨ ・ うへつ

- can push $\phi \alpha$ and $\phi \pi$ past $\phi \pi/4$
- can push $\phi \alpha$ and $\phi \pi$ past $\phi \pi/2$
- can push $\phi \pi$ and $\phi \pi$ past $V \in \left\{ \begin{array}{c} \phi \pi/4 \\ \phi \pi/2 \end{array} \right\}$

Normal form diagrams act non-trivially

Lemma

No normal form diagram represents the identity operator.

- ▶ write any single-qubit density operator as xX + yY + zZ, where $x, y, z \in \mathbb{R}$ and X, Y, Z are the Pauli matrices
- Clifford unitaries act on the vectors (x, y, z) by permuting the elements and adding minus signs; T sends

$$(x,y,z)\mapsto rac{1}{\sqrt{2}}\left(x-y,x+y,z\sqrt{2}
ight).$$

• if *D* is a normal form operator, $D|0\rangle$ has vector

$$\frac{1}{\sqrt{2^m}} \left(x_1 + x_2 \sqrt{2}, y_1 + y_2 \sqrt{2}, z_1 + z_2 \sqrt{2} \right)$$

where $x_1, x_2, y_1, y_2, z_1, z_2 \in \mathbb{Z}$

 \blacktriangleright by parity arguments, can show none of them represent $|0\rangle$

(ロ) (同) (三) (三) (三) (○) (○)

Normal forms are unique

Lemma

The normal form of the inverse of a normal form diagram has the same number of $\pi/4$ nodes as the original diagram.

◆□▶ ◆□▶ ▲□▶ ▲□▶ □ のQ@

Normal forms are unique

Lemma

The normal form of the inverse of a normal form diagram has the same number of $\pi/4$ nodes as the original diagram.

Theorem

The normal form is unique.

Suppose D and D' are non-identical normal form diagrams

・ロト ・ 同 ・ ・ ヨ ・ ・ ヨ ・ うへつ

- Can then show that $D^{-1} \circ D'$ has non-trivial normal-form
- Thus $\llbracket D \rrbracket \neq \llbracket D' \rrbracket$, and by soundness $D \neq D'$

Normal forms are unique

Lemma

The normal form of the inverse of a normal form diagram has the same number of $\pi/4$ nodes as the original diagram.

Theorem

The normal form is unique.

Suppose D and D' are non-identical normal form diagrams

(ロ) (同) (三) (三) (三) (○) (○)

- Can then show that $D^{-1} \circ D'$ has non-trivial normal-form
- Thus $\llbracket D \rrbracket \neq \llbracket D' \rrbracket$, and by soundness $D \neq D'$

Corollary

The zx-calculus is complete for the single-qubit Clifford+T group.

Outline

Introduction

The zx-calculus is incomplete [Schröder & Zamdzhiev 2014]

Completeness results

Stabilizer quantum mechanics The single-qubit Clifford+T group

Outlook & Conclusions

◆□ ▶ ◆□ ▶ ◆ □ ▶ ◆ □ ▶ ◆ □ ▶ ◆ □ ▶

Related work and open questions

- Stabilizer completeness proof carries over to a ZX-like graphical calculus for Spekkens' toy theory (joint work with Ali Nabi Duman)
- can the two completeness results be combined into a completeness proof for the multi-qubit Clifford+T group?
- zero-diagram completeness?
- what happens if we put the scalars back in (possibly still ignoring complex phases)?

(ロ) (同) (三) (三) (三) (○) (○)

can the completeness results be implemented in Quantomatic?

Summary

- ZX-calculus is not complete in general but fragments of it are complete, e.g.
 - ► line graphs where all phases are multiples of π/4 (single-qubit Clifford+T group)
 - ► diagrams where all phases are multiples of π/2 (stabilizer quantum mechanics)

(ロ) (同) (三) (三) (三) (○) (○)

ongoing work to extend completeness results

Summary

- ZX-calculus is not complete in general but fragments of it are complete, e.g.
 - ► line graphs where all phases are multiples of π/4 (single-qubit Clifford+T group)
 - diagrams where all phases are multiples of π/2 (stabilizer quantum mechanics)

(ロ) (同) (三) (三) (三) (○) (○)

ongoing work to extend completeness results

Thank you!