Categorical Quantum Mechanics: The “Monoidal” Approach

Samson Abramsky

Introduction

Categorical Approaches to Physics

- Crane, Baez, Dolan et al. Higher-dimensional categories, TQFT's, categorification, etc.

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We aim to give an introduction to the "monoidal" approach, which could as well be called "linear".

Some distinctive features

Comparison with topos approach:

monoidal	*vs.*	cartesian
linear	*vs.*	intuitionistic
processes	*vs.*	propositions
geometry of proofs	*vs.*	geometric logic

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Comparison with n-categories approach. We emphasize:

- operational aspects
- interplay of quantum and classical
- compositionality
- open *vs.* closed systems.

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- operational aspects
- interplay of quantum and classical
- compositionality
- open *vs.* closed systems.

These are important for applications to quantum informatics, but also of foundational significance.

Quantum Operations/Quantum Processes

Bits and Qubits: Classical *vs.* Quantum Operations

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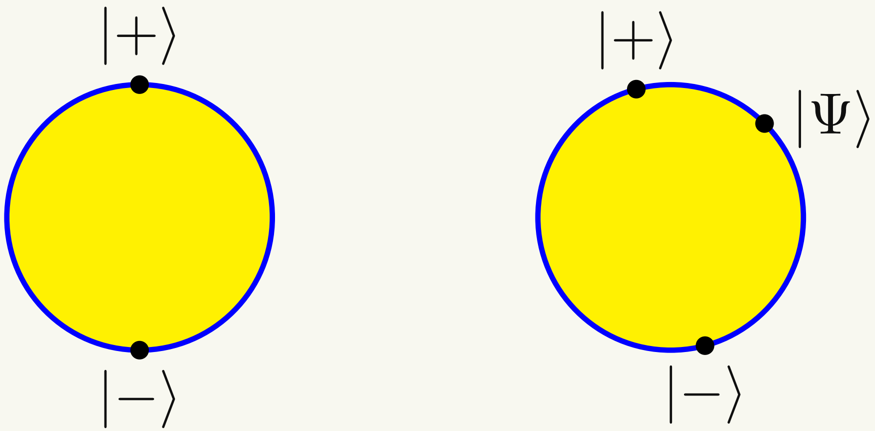

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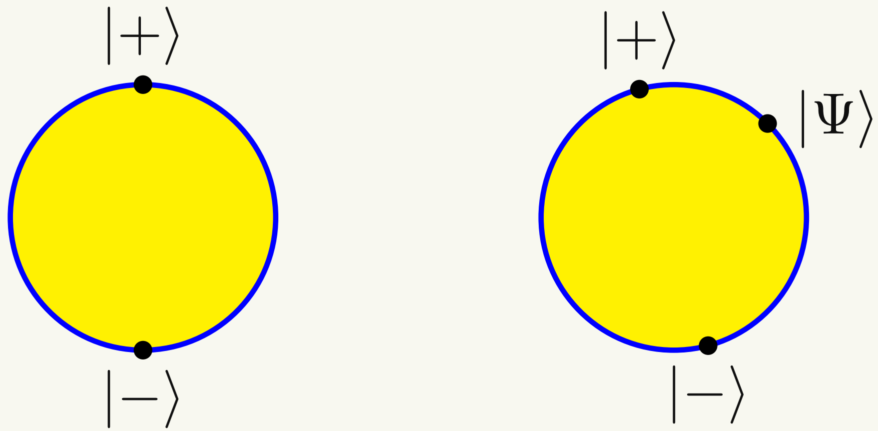

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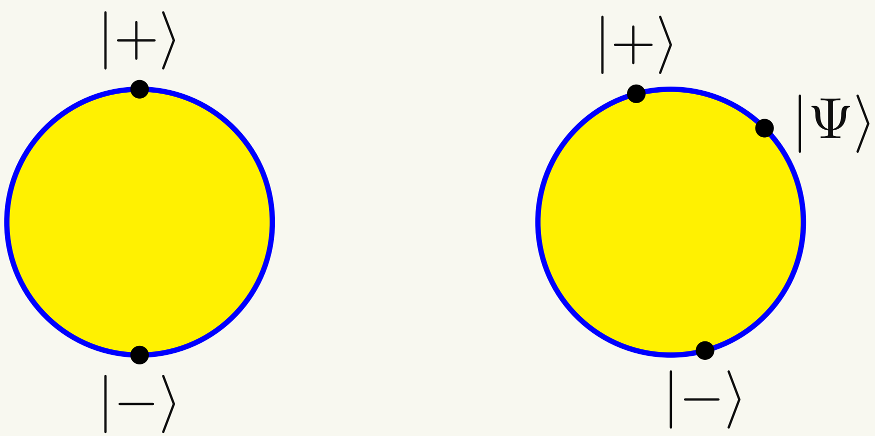
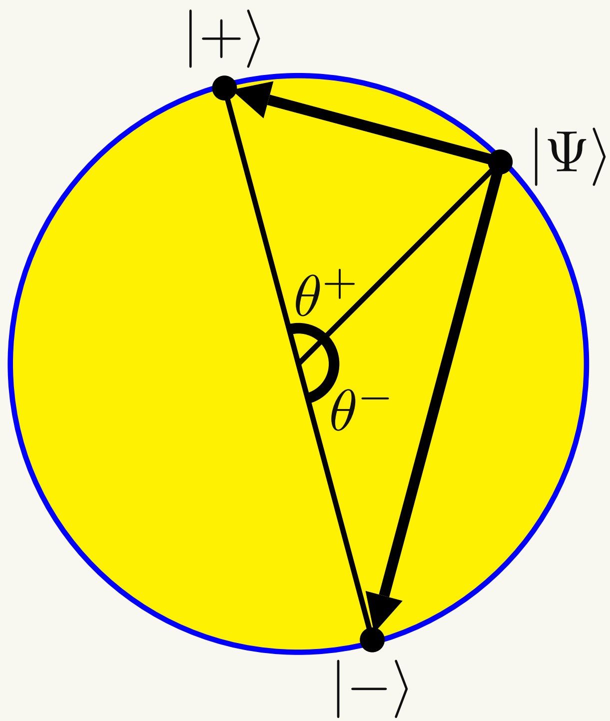


- measurements of qubits
  - have two outcomes $|-\rangle$, $|+\rangle$
  - change the value $|\psi\rangle$
- admit **unitary transformations**, *i.e.* antipodes and angles are preserved.



We have **partial constant maps** $Q \longrightarrow Q$ on the sphere Q

$$\mathcal{P}_+ : |\psi\rangle \mapsto |+\rangle \qquad \mathcal{P}_- : |\psi\rangle \mapsto |-\rangle$$

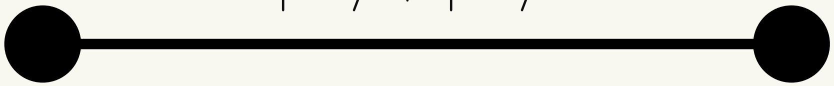
which have chance prob$(\theta^\sharp)$ for $\sharp \in \{+, -\}$.

We know the value **after** the measurement, but not what the qubit was **before** the measurement!

So measurements change the state, and apparently lose information. This seems like bad news! However . . .

Bell state:

$$|00\rangle + |11\rangle$$



EPR state:

$$|01\rangle + |10\rangle$$

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Compound systems are represented by **tensor product**: $\mathcal{H}_1 \otimes \mathcal{H}_2$.
Typical element:

$$\sum_i \lambda_i \cdot \phi_i \otimes \psi_i$$

Superposition encodes **correlation**. Einstein's 'spooky action at a distance'. Even if the particles are spatially separated, measuring one has an effect on the state of the other.
Bell's theorem: QM is **essentially non-local**.

Towards a general formalism for describing physical processes

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- Basic reflection of space: must be able to describe **compound systems**, operations **localized** to part of a compound system, and operations performed independently on different parts of a compound system — **parallel composition**.

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- Basic reflection of space: must be able to describe **compound systems**, operations **localized** to part of a compound system, and operations performed independently on different parts of a compound system — **parallel composition**.

So we want a general setting in which we can describe processes (of whatever kind) closed under sequential and parallel composition.

Basic Setting: Symmetric Monoidal Categories

A category $\mathcal{C}$ has **objects** (types) A, B, C, . . . , and for each pair of objects A, B a set of **morphisms** $\mathcal{C}(A, B)$. (Notation: $f : A \rightarrow B$). It also has identities $\mathrm{id}_A : A \rightarrow A$, and composition $g \circ f$ when types match:

$$A \xrightarrow{f} B \xrightarrow{g} C$$

Categories allow us to deal explictly with **typed processes**, e.g.

Logic	**Programming**	**Computation**
Propositions	Data Types	States
Proofs	Programs	Transitions

For QM:

Types of system (e.g. qubits Q) $\rightsquigarrow$ objects
Processes/Operations (e.g. measurements) $\rightsquigarrow$ morphisms

Symmetric Monoidal Categories

A **symmetric monoidal category** comes equipped with an associative operation which acts on **both** objects **and** morphisms — a **bifunctor**:

$$A \otimes B \qquad f_1 \otimes f_2 : A_1 \otimes A_2 \longrightarrow B_1 \otimes B_2$$

There is also a symmetry operation

$$\sigma_{A,B} : A \otimes B \longrightarrow B \otimes A$$

which satisfies some 'obvious' rules, e.g. naturality:

$$\begin{tikzcd}
A_1 \otimes A_2 \arrow[r, "f_1 \otimes f_2"] \arrow[d, "\sigma_{A_1,A_2}"'] & B_1 \otimes B_2 \arrow[d, "\sigma_{B_1,B_2}"] \\
A_2 \otimes A_1 \arrow[r, "f_1 \otimes f_2"] & B_2 \otimes B_1
\end{tikzcd}$$

The Logic of Tensor Product

Tensor can express **independent** or **concurrent** actions (mathematically: bifunctoriality):

$$\begin{tikzcd} A_1 \otimes A_2 \arrow[r, "f_1 \otimes \mathsf{id}"] \arrow[d, "\mathsf{id} \otimes f_2"'] & B_1 \otimes A_2 \arrow[d, "\mathsf{id} \otimes f_2"] \\ A_1 \otimes B_2 \arrow[r, "f_1 \otimes \mathsf{id}"'] & B_1 \otimes B_2 \end{tikzcd}$$

But tensor is **not** a cartesian product, in the sense that **we cannot reconstruct an 'element' of the tensor from its components**.

This turns out to comprise the **absence** of

$$\begin{array}{lll} & A \xrightarrow{\Delta} A \otimes A & A_1 \otimes A_2 \xrightarrow{\pi_i} A_i \\ \text{Cf.} & A \vdash A \wedge A & A_1 \wedge A_2 \vdash A_i \end{array}$$

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No-Cloning and No-Deleting are built in!

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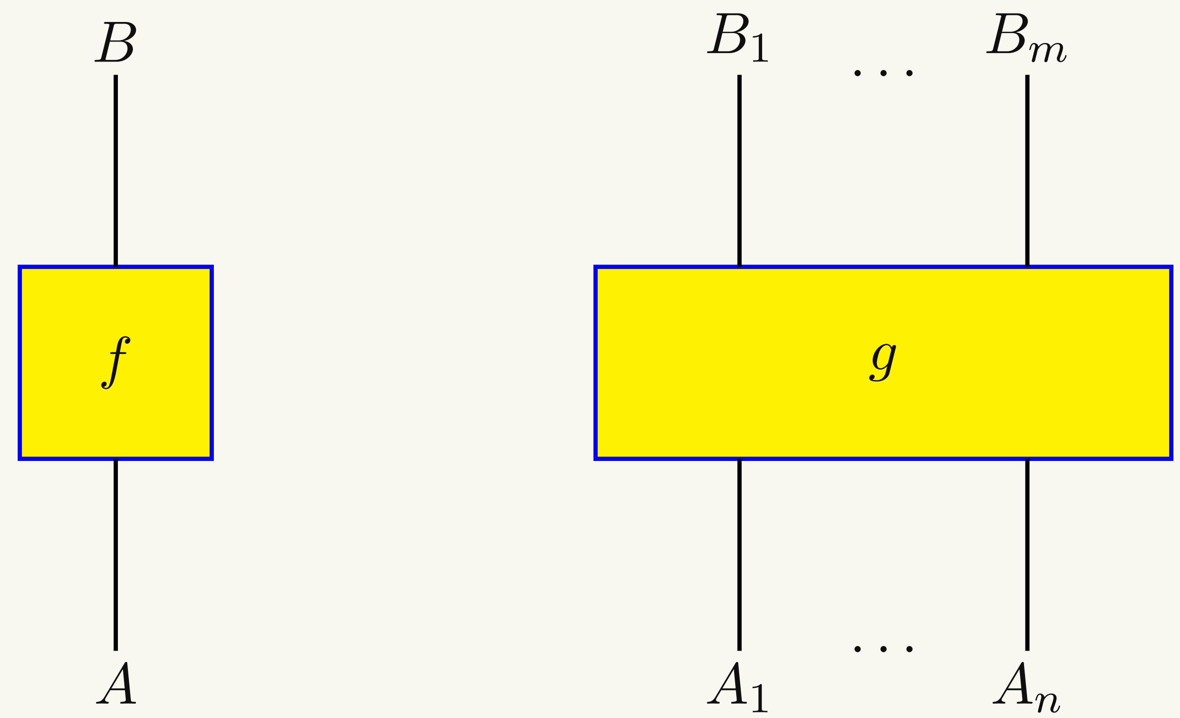
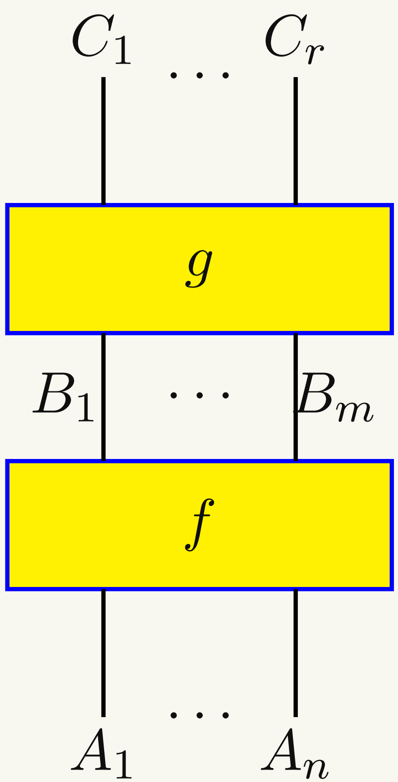
There is a natural objection to this, that this is too abstract, and we lose all grip of what is going on.

But this objection does not really hold water! Monoidal categories, quite generally, admit a beautiful graphical or diagrammatic calculus (Joyal, Street et al.) which makes equational proofs perspicuous, and is sound and complete for equational reasoning in monoidal categories. It also supports links with Logic (e.g. Proof Nets) and with Geometry (Knots, Braids, Temperley-Lieb algebra etc.)

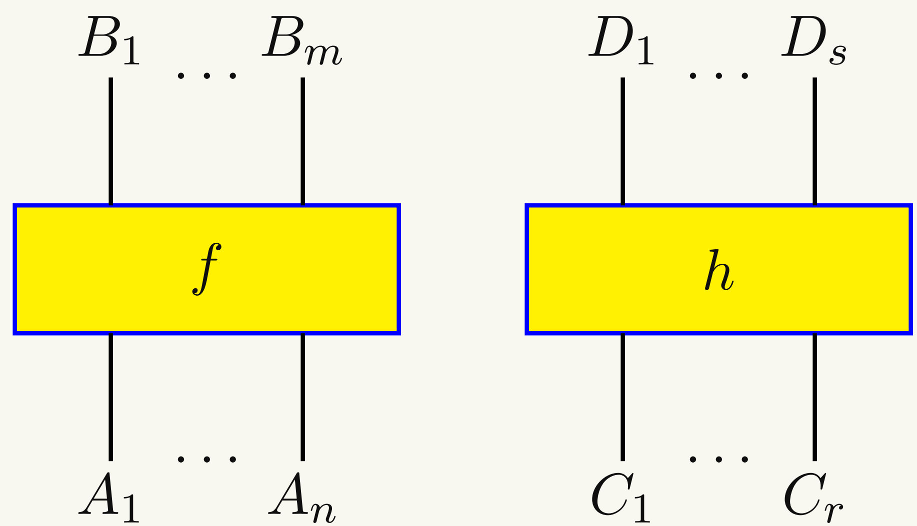


$$f : A \longrightarrow B \qquad g : A_1 \otimes \cdots \otimes A_n \longrightarrow B_1 \otimes \cdots \otimes B_m.$$

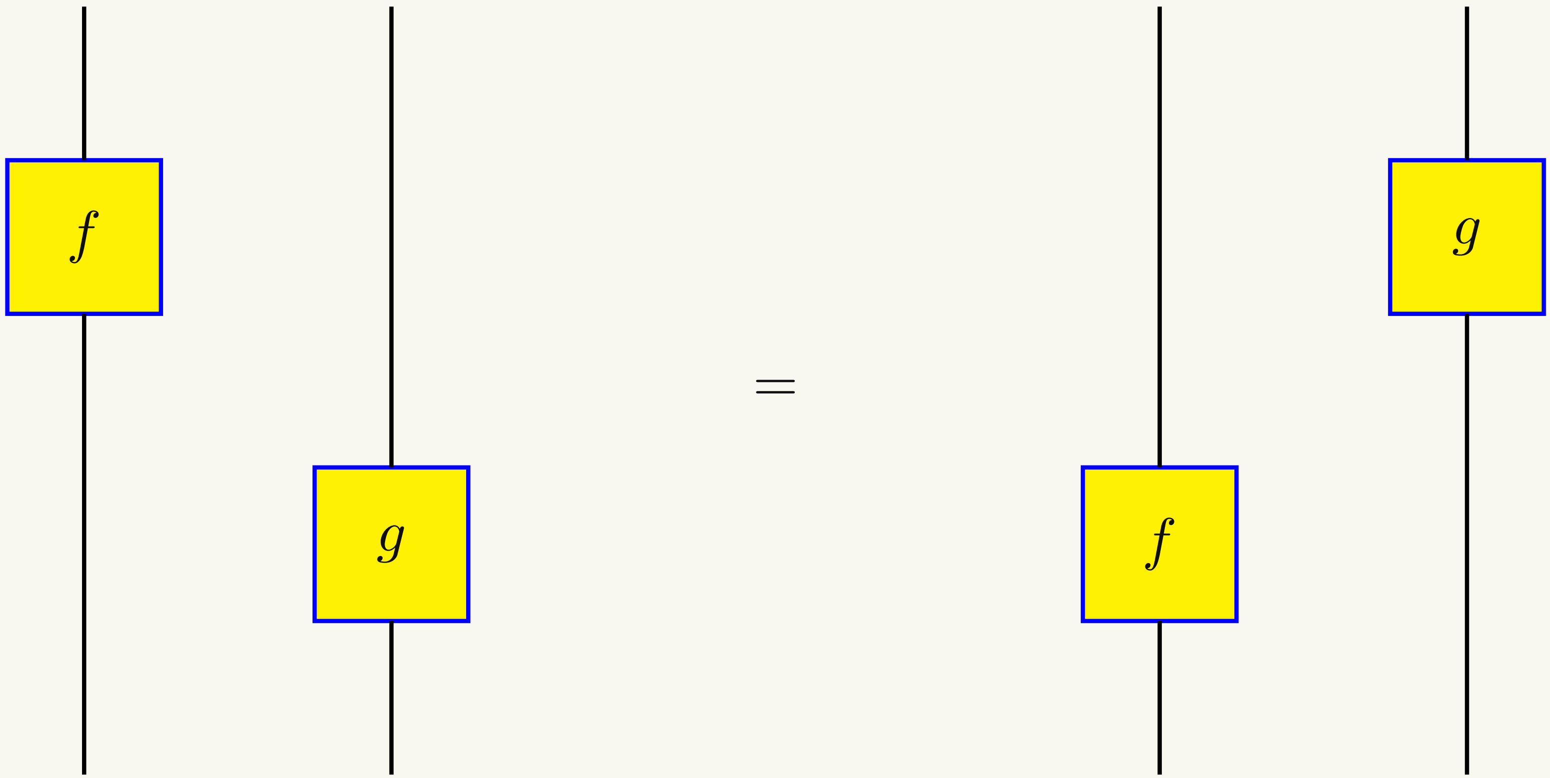
Series and Parallel Composition



$g \circ f$

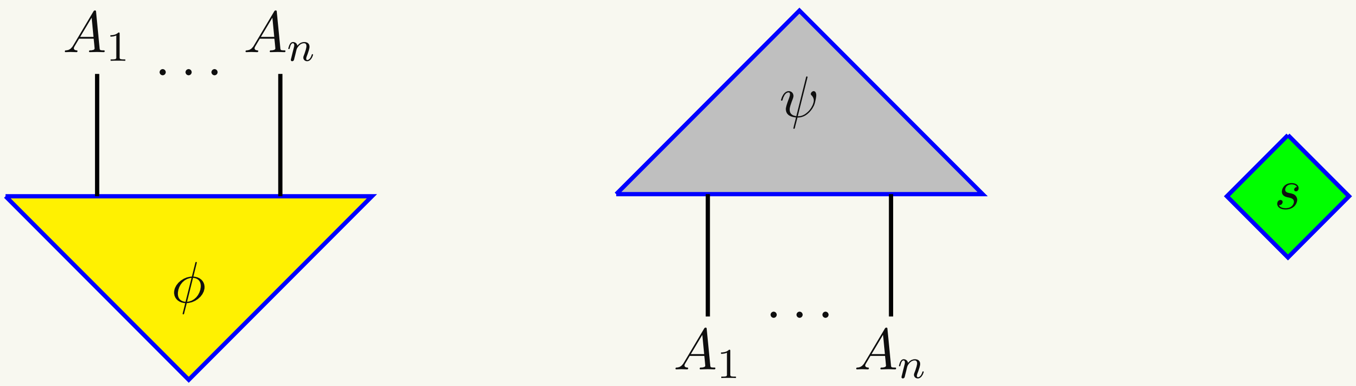


$f \otimes h$



$$(f \otimes 1) \circ (1 \otimes g) \quad = \quad f \otimes g \quad = \quad (1 \otimes g) \circ (f \otimes 1)$$

Bras, Kets and Scalars



$$\phi : A_1 \otimes \cdots \otimes A_n \longrightarrow I \qquad \psi : I \longrightarrow A_1 \otimes \cdots \otimes A_n \qquad s : I \longrightarrow I.$$

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Kets:	no inputs
Scalars:	no inputs or outputs.

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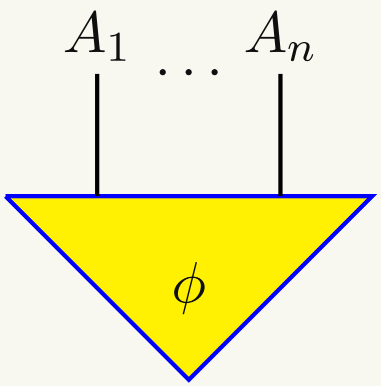

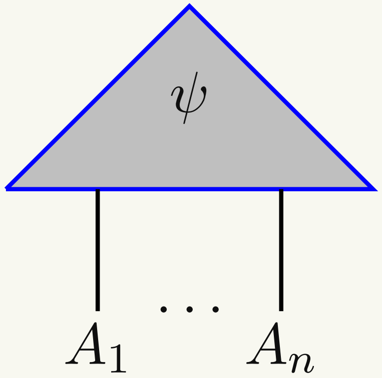

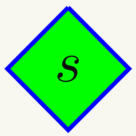

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Two-dimensional generalization of **Dirac notation**!

$$\langle \phi \mid \qquad \mid \psi \rangle \qquad \langle \phi \mid \psi \rangle$$

Interlude: The Miracle of Scalars

A **scalar** in a monoidal category is a morphism $s : I \longrightarrow I$.

Examples: $(\mathbf{FdVec}_{\mathbb{K}}, \otimes)$, $(\mathbf{Rel}, \times)$.

Scalars in monoidal categories

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(1) $\mathcal{C}(I, I)$ is a commutative monoid

$$\begin{tikzcd}
I \arrow[r, "\rho_I^{-1}"] & I \otimes I \arrow[r, equal] & I \otimes I \arrow[r, "\lambda_I"] \arrow[d, "1\otimes t"] & I \arrow[d, "t"] \\
I \arrow[r, "\rho_I^{-1}"] \arrow[u, "s"] \arrow[d, "t"'] & I \otimes I \arrow[r, "s\otimes t"] \arrow[u, "s\otimes 1"'] \arrow[d, "1\otimes t"'] & I \otimes I \arrow[r, "\lambda_I"] & I \\
I \arrow[r, "\lambda_I^{-1}"'] & I \otimes I \arrow[r, equal] & I \otimes I \arrow[r, "\rho_I"'] \arrow[u, "s\otimes 1"] & I \arrow[u, "s"']
\end{tikzcd}$$

using the coherence equation $\lambda_I = \rho_I$.

Scalar Product

(2) Each scalar $s : I \to I$ induces a natural transformation

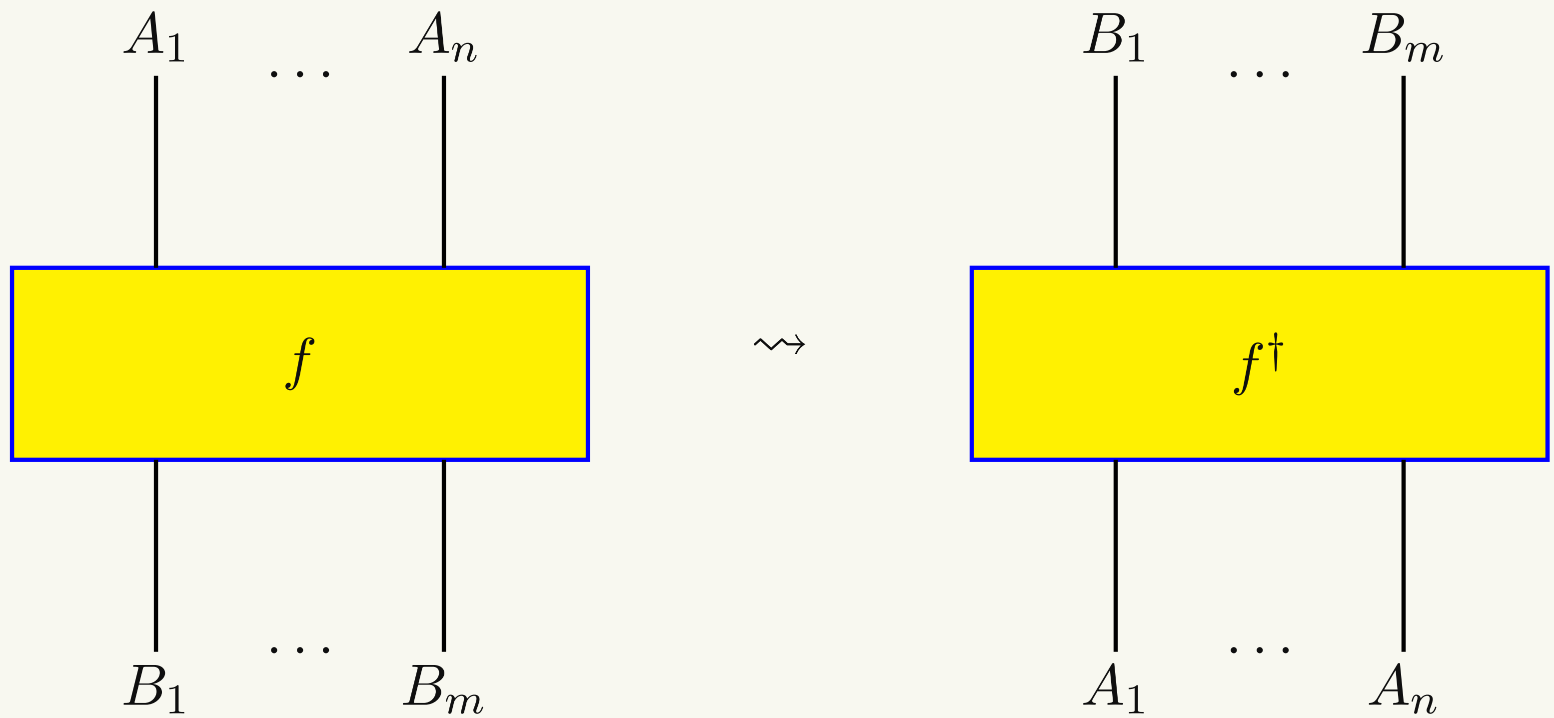
$$s_A : A \xrightarrow{\simeq} I \otimes A \xrightarrow{s \otimes 1_A} I \otimes A \xrightarrow{\simeq} A\,.$$

$$\begin{tikzcd} A \arrow[r, "s_A"] \arrow[d, "f"'] & A \arrow[d, "f"] \\ B \arrow[r, "s_B"'] & B \end{tikzcd}$$

We write $s \bullet f$ for $f \circ s_A = s_B \circ f$. Note that

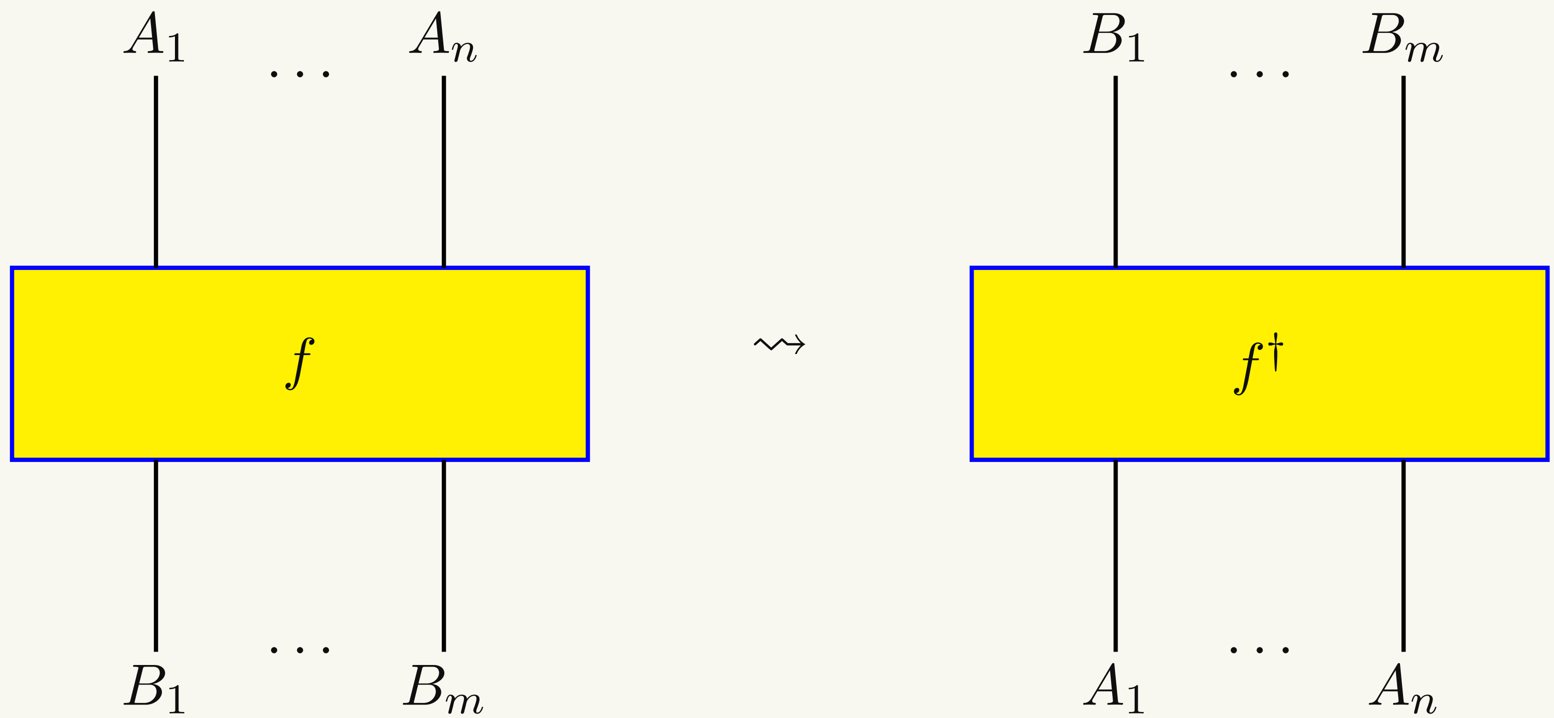
$$\begin{array}{rcl} s \bullet (t \bullet f) & = & (s \circ t) \bullet f \\ (s \bullet g) \circ (r \bullet f) & = & (s \circ r) \bullet (g \circ f) \\ (s \bullet f) \otimes (t \bullet g) & = & (s \circ t) \bullet (f \otimes g) \end{array}$$

Dagger Monoidal Categories



$$\frac{f : A \longrightarrow B}{f^\dagger : B \longrightarrow A} \qquad f^{\dagger\dagger} = f \qquad (g \circ f)^\dagger = f^\dagger \circ g^\dagger$$

Adjoint Arrows — reflection in the x-axis



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We can turn kets into bras and vice versa — full scale Dirac notation! Given $\phi, \psi : I \longrightarrow A$,

$$\langle \phi \mid \psi \rangle = \phi^\dagger \circ \psi : I \longrightarrow I$$

which is indeed a scalar!

Unitaries

We can also define **unitaries**. An isomorphism $U : A \xrightarrow{\cong} B$ is unitary if $U^{-1} = U^{\dagger}$.

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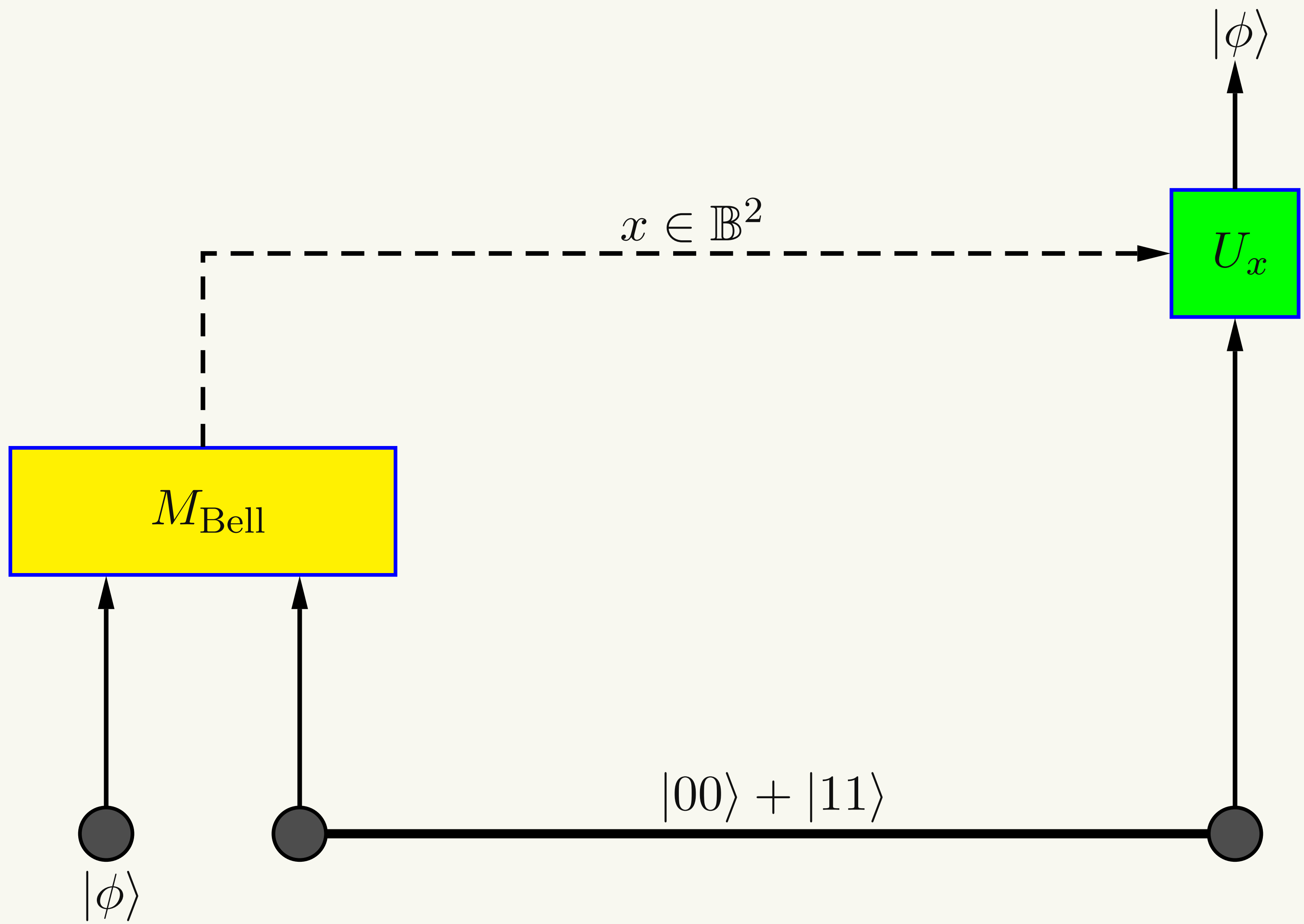
Examples: **Hilb**, **Rel**.

Entanglement, Bell States and Compact Categories

From ‘paradox’ to ‘feature’: Teleportation



Entangled states as linear maps

Indeed, $\mathcal{H}_1 \otimes \mathcal{H}_2$ is spanned by

$$\begin{array}{ccc} |11\rangle & \cdots & |1m\rangle \\ \vdots & \ddots & \vdots \\ |n1\rangle & \cdots & |nm\rangle \end{array}$$

hence

$$\sum_{i,j} \alpha_{ij} |ij\rangle \quad \longleftrightarrow \quad \begin{pmatrix} \alpha_{11} & \cdots & \alpha_{1m} \\ \vdots & \ddots & \vdots \\ \alpha_{n1} & \cdots & \alpha_{nm} \end{pmatrix} \quad \longleftrightarrow \quad |i\rangle \mapsto \sum_j \alpha_{ij} |j\rangle$$

Pairs $|\psi_1, \psi_2\rangle$ are a special case — $|ij\rangle$ in a well-chosen basis.

This is **Map-State Duality**.

A projector onto the 1-dimensional subspace spanned by a vector $|\psi\rangle$ will be written $\mathcal{P}_\psi$. It is essentially (up to scalar multiples) a "partial constant map"

$$\mathcal{P}_\psi : |\phi\rangle \mapsto |\psi\rangle.$$

This will correspond e.g. to a branch of a (projective, non-degenerate) measurement, or to a preparation.

Some notation for projectors

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We combine this notation with Map-State Duality: we write a projector $\mathcal{P}_\psi$ on a tensor product space $\mathcal{H}_1 \otimes \mathcal{H}_2$ as $\mathcal{P}_f$, where f is the linear map $\mathcal{H}_1 \longrightarrow \mathcal{H}_2$ associated to ψ under Map-State Duality.

On the trail of structure

The identity map

$$|\mathsf{id} : Q \to Q\rangle \in Q \otimes Q \qquad |11\rangle + \cdots + |nn\rangle \;\longleftrightarrow\; |i\rangle \mapsto |i\rangle$$

is the **Bell state**.

A measurement of $Q \otimes Q$ has four outcomes

$$|f_1\rangle, |f_2\rangle, |f_3\rangle, |f_4\rangle \qquad (\textit{cf.}\ |00\rangle, |01\rangle, |10\rangle, |11\rangle)$$

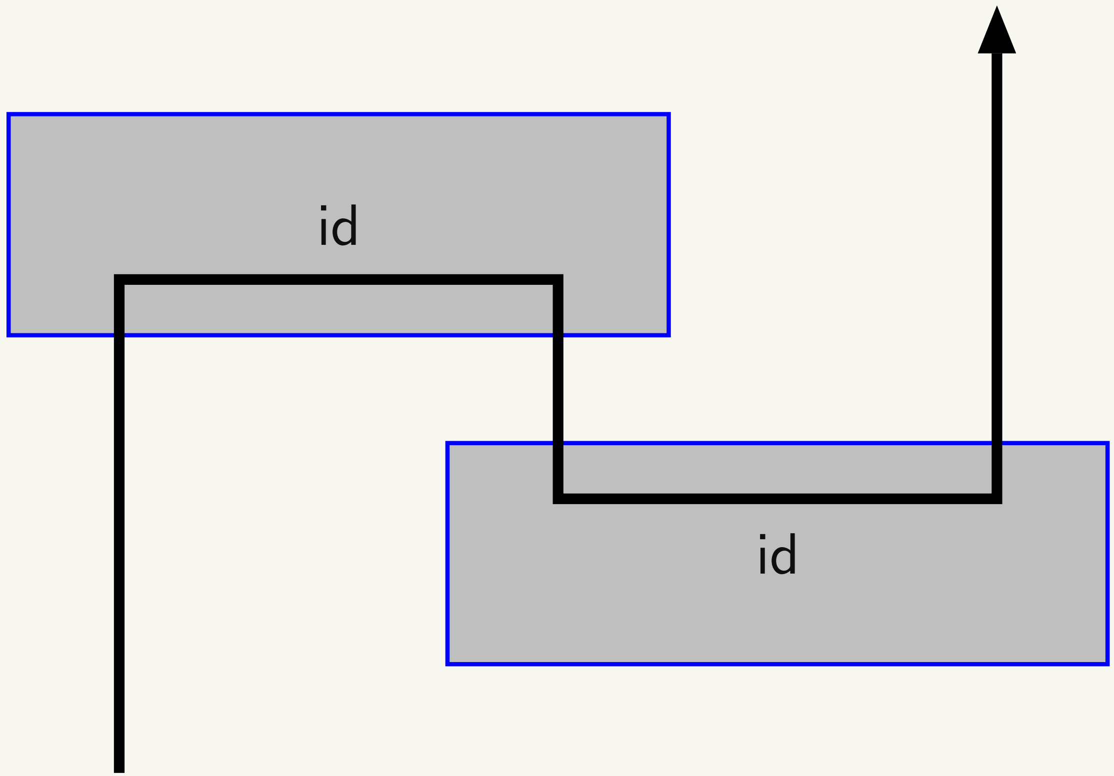
and corresponding projectors

$$\mathcal{P}_f : Q \otimes Q \to Q \otimes Q :: |g\rangle \mapsto |f\rangle$$

E.g. the Bell state is produced by

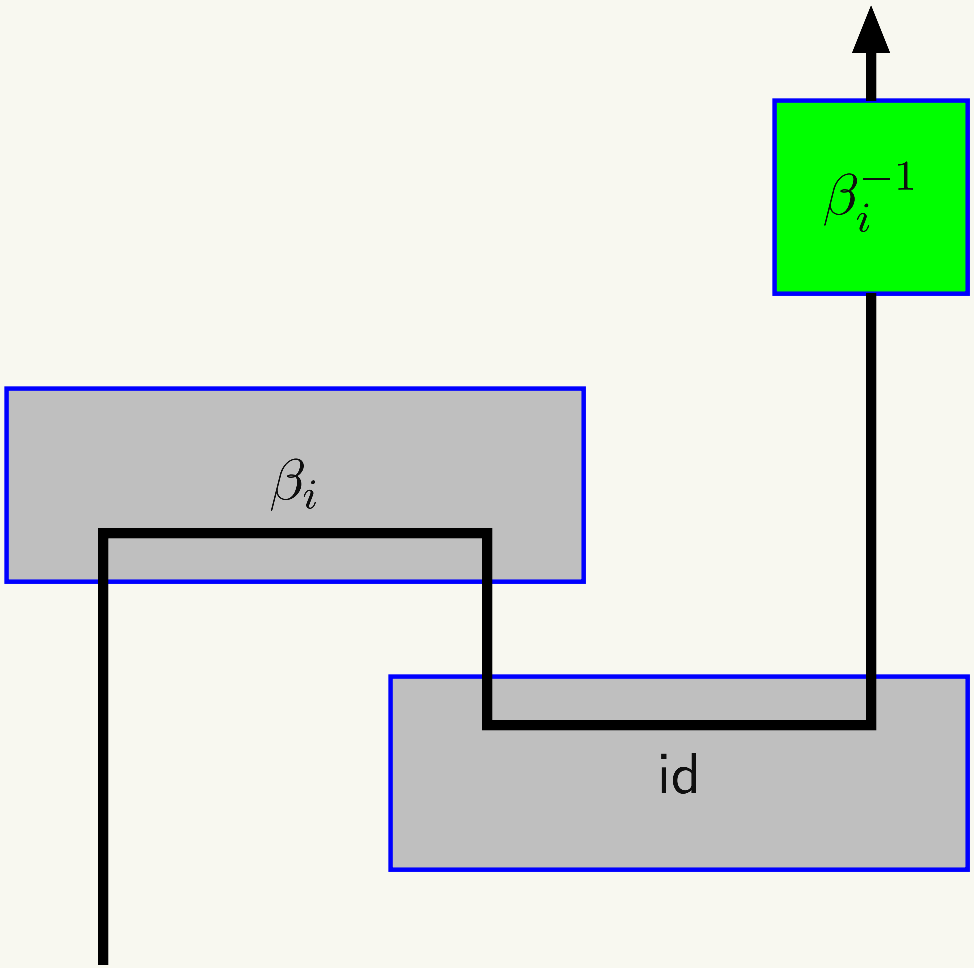
$$\mathcal{P}_{\mathsf{id}} : Q \otimes Q \to Q \otimes Q :: |g\rangle \mapsto |\mathsf{id}\rangle$$

Key Question: Do entangled states *qua* functions compose (somehow)?



$$\mathrm{id} \circ \mathrm{id} = \mathrm{id}$$

Teleportation: general case



$$\beta_i^{-1} \circ \mathsf{id} \circ \beta_i = \mathsf{id}, \quad 1 \leq i \leq 4$$

Categorical Axiomatics

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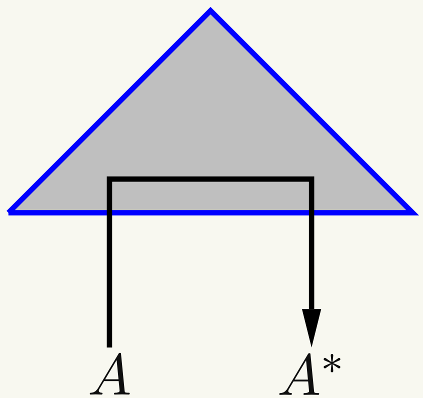
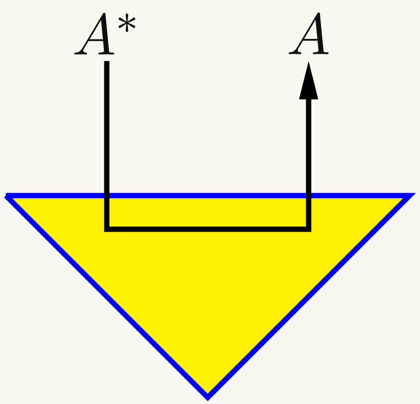
- Diagrammatically, the (partial) trace closes (part of) the system; when we close the whole system we get **loops** — *i.e.* scalars!

Cups and Caps

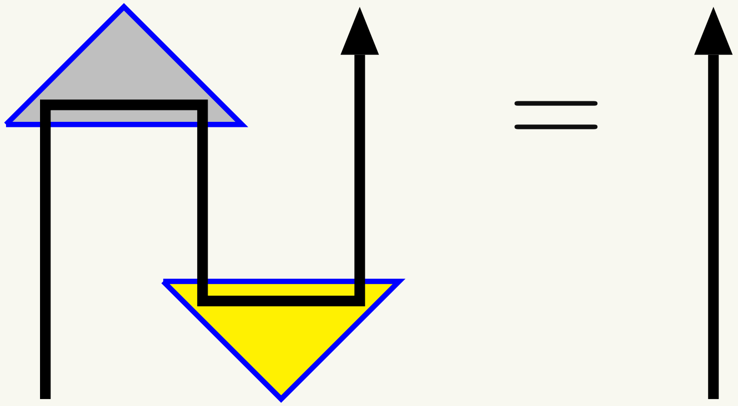


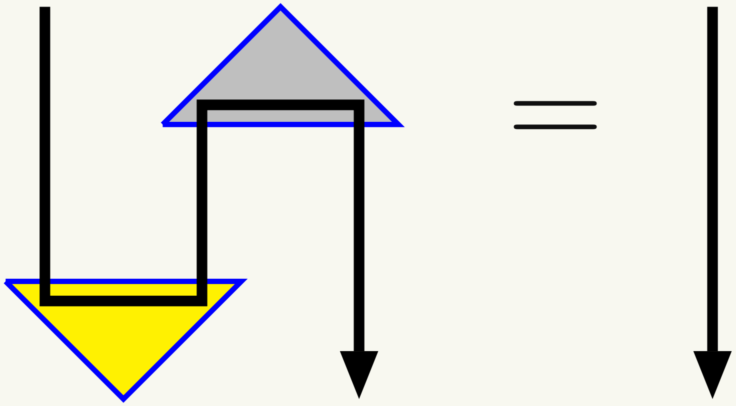


$$\epsilon_A : A \otimes A^* \longrightarrow I \qquad\qquad \eta_A : I \longrightarrow A^* \otimes A.$$

Caps = Bell States; Cups = Bell Tests.

Compact Closure: The basic algebraic laws for units and counits.

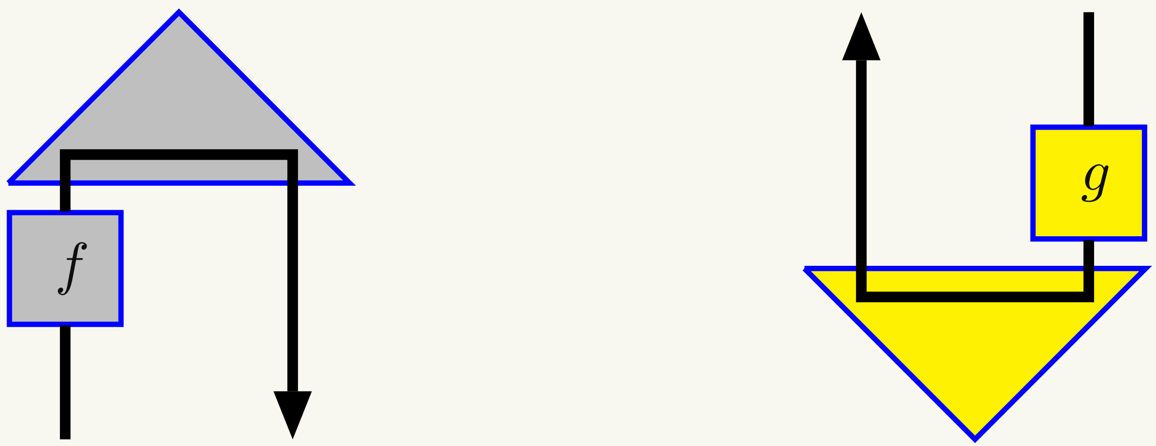




$(\epsilon_A \otimes 1_A) \circ (1_A \otimes \eta_A) = 1_A$

$(1_{A^*} \otimes \epsilon_A) \circ (\eta_A \otimes 1_{A^*}) = 1_{A^*}$

For coherence with the dagger structure, we require that $\epsilon_A = \eta_A^\dagger$.

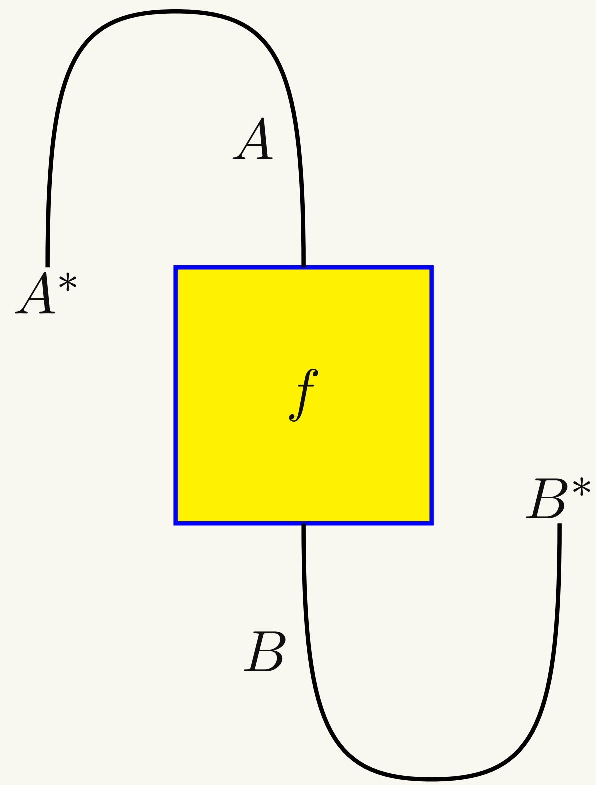


$\llcorner f \lrcorner : A \otimes B^* \longrightarrow I$ $\quad\quad$ $\ulcorner f \urcorner : I \longrightarrow A^* \otimes B$

$$\mathcal{C}(A \otimes B^*, I) \simeq \mathcal{C}(A, B) \simeq \mathcal{C}(I, A^* \otimes B).$$

This is the general form of Map-State duality.

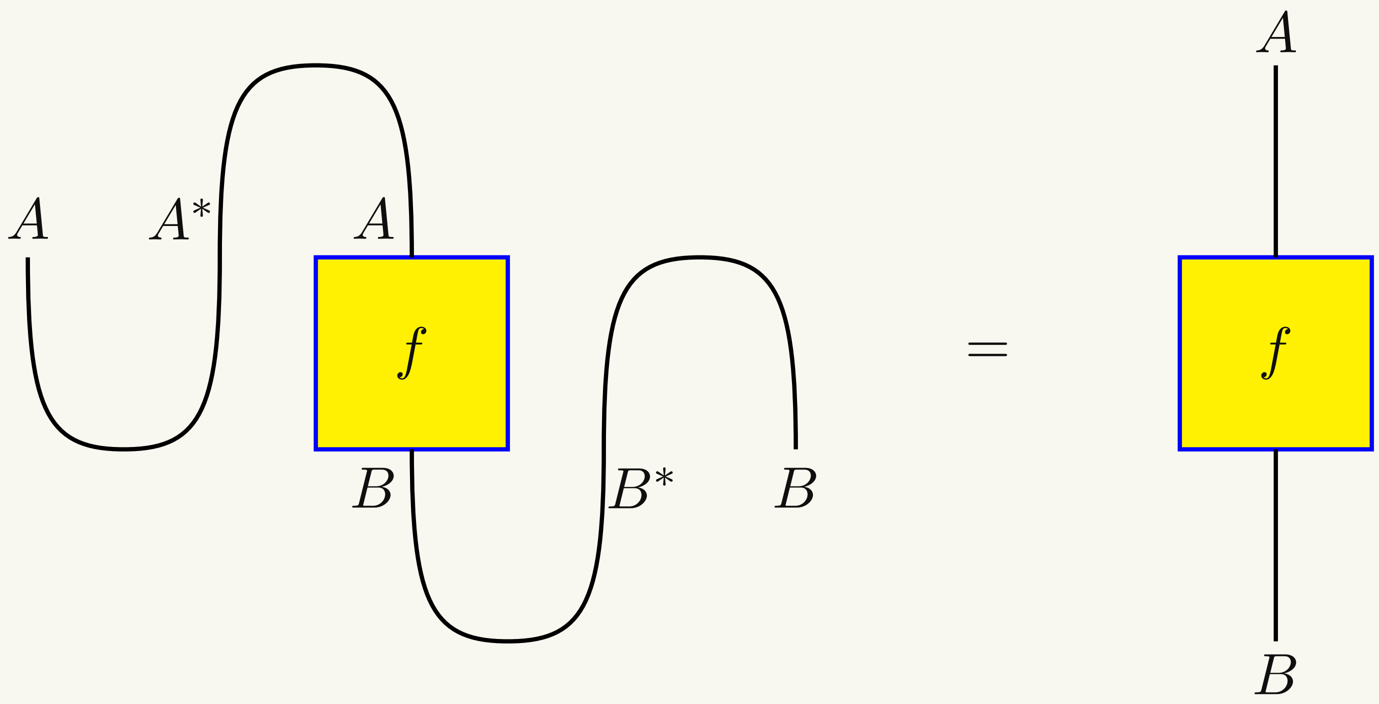
Definition of Duality



$$f^* = (1 \otimes \epsilon_B) \circ (1 \otimes f \otimes 1) \circ (\eta_A \otimes 1).$$

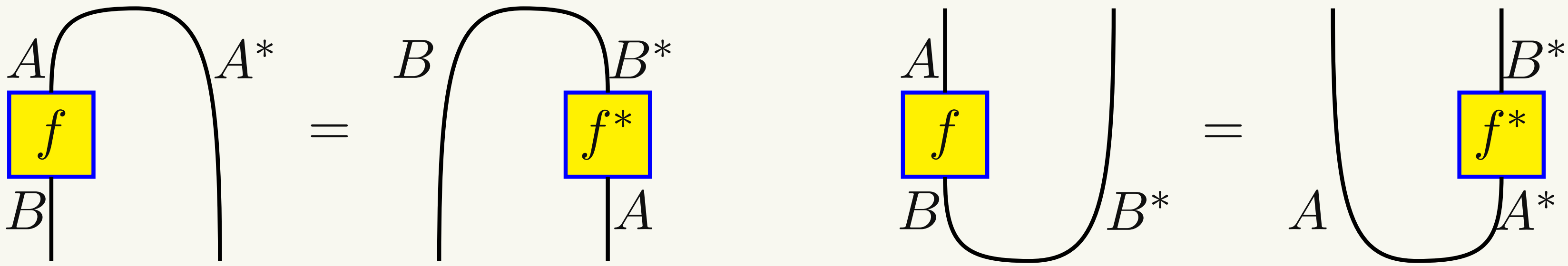
A Little Taste of Diagrammatic Reasoning

Duality is Involutive

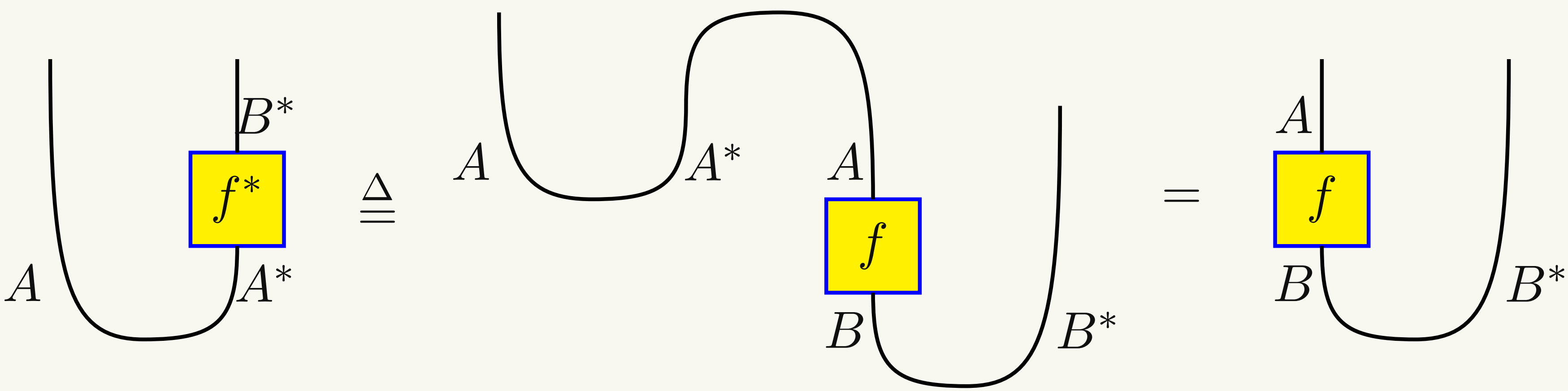


$$f^{**} = f.$$

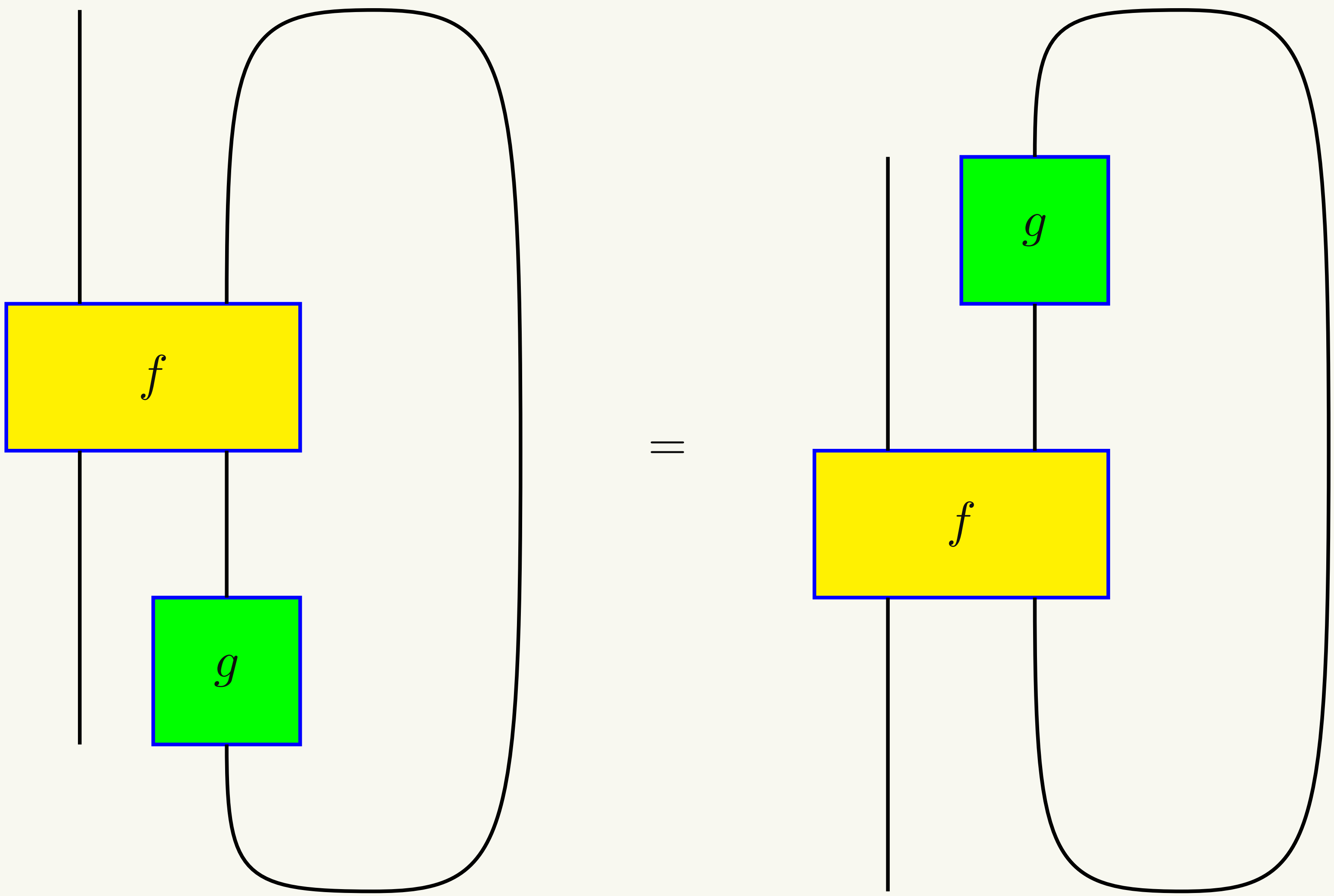
Moving Boxes round Cups and Caps

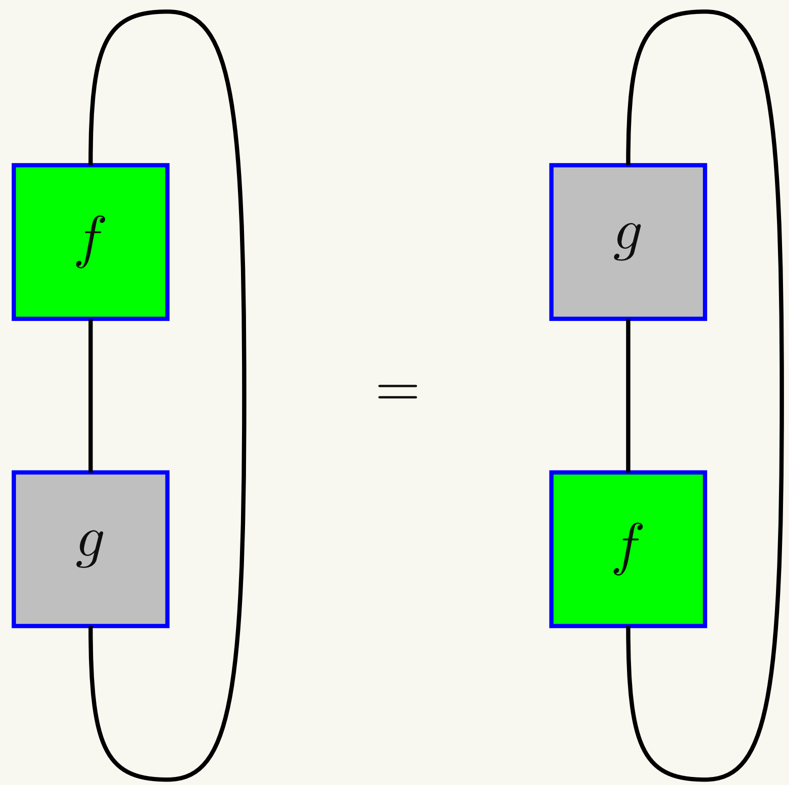


Diagrammatic Proof

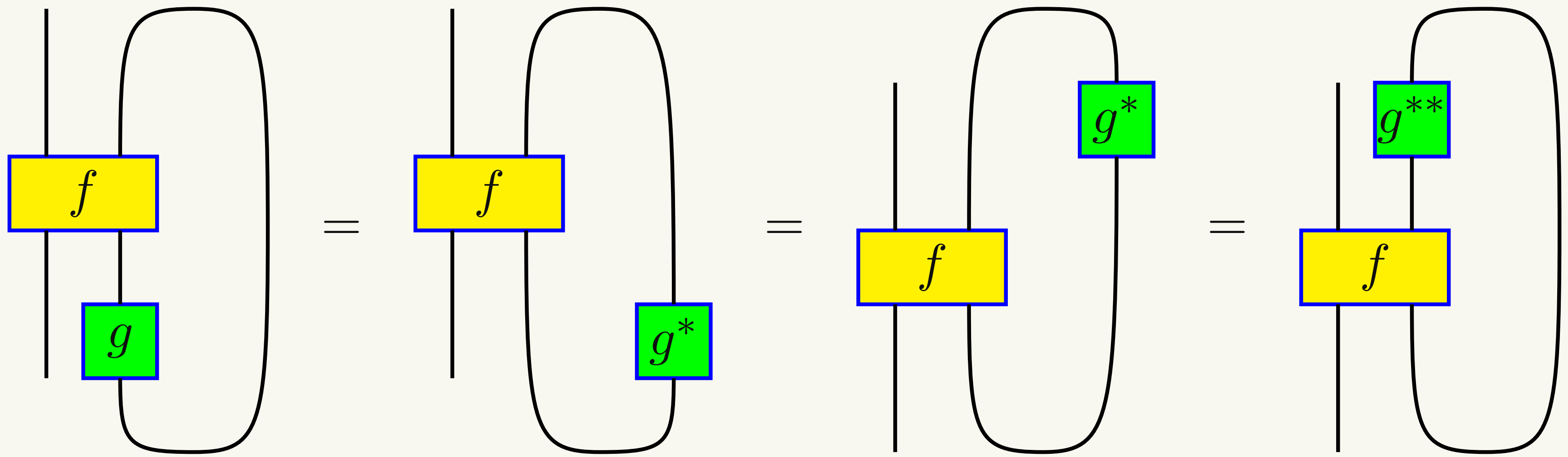


Feedback Dinaturality




f
g
$=$
g
f

Graphical Proof of Feedback Dinaturality



We use $g^{**} = g$ to conclude.

Illuminating Quantum Information Flow in Entangled Systems

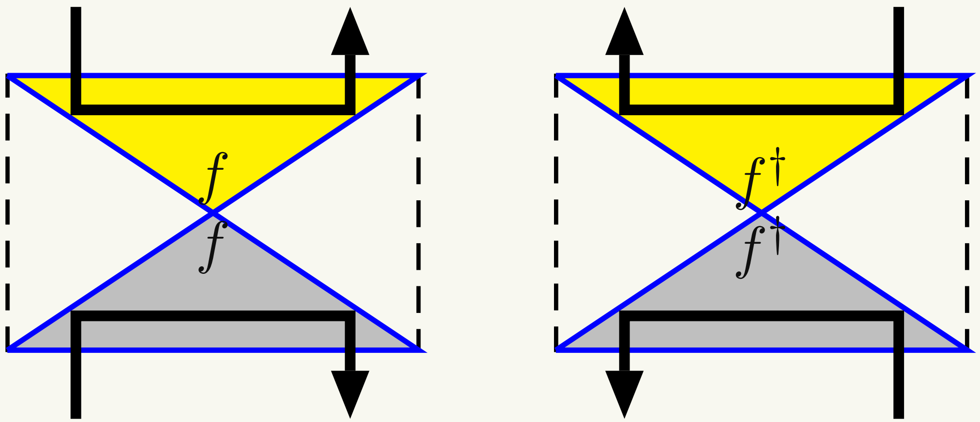
Bipartite Projectors

Information flow in entangled states can be captured mathematically by the isomorphism

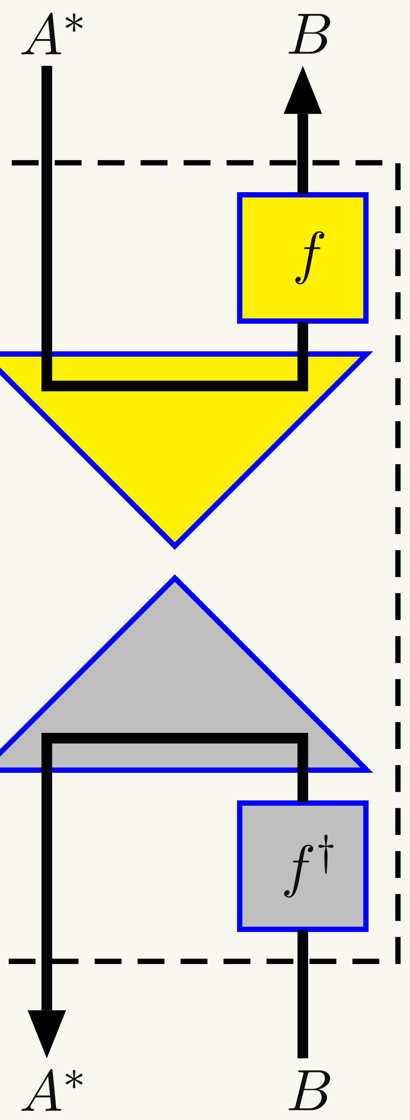
$$\mathbf{Hom}(A, B) \cong A^* \otimes B.$$

This leads to a **decomposition** of bipartite projectors into "names" (preparations) and "conames" (measurements).

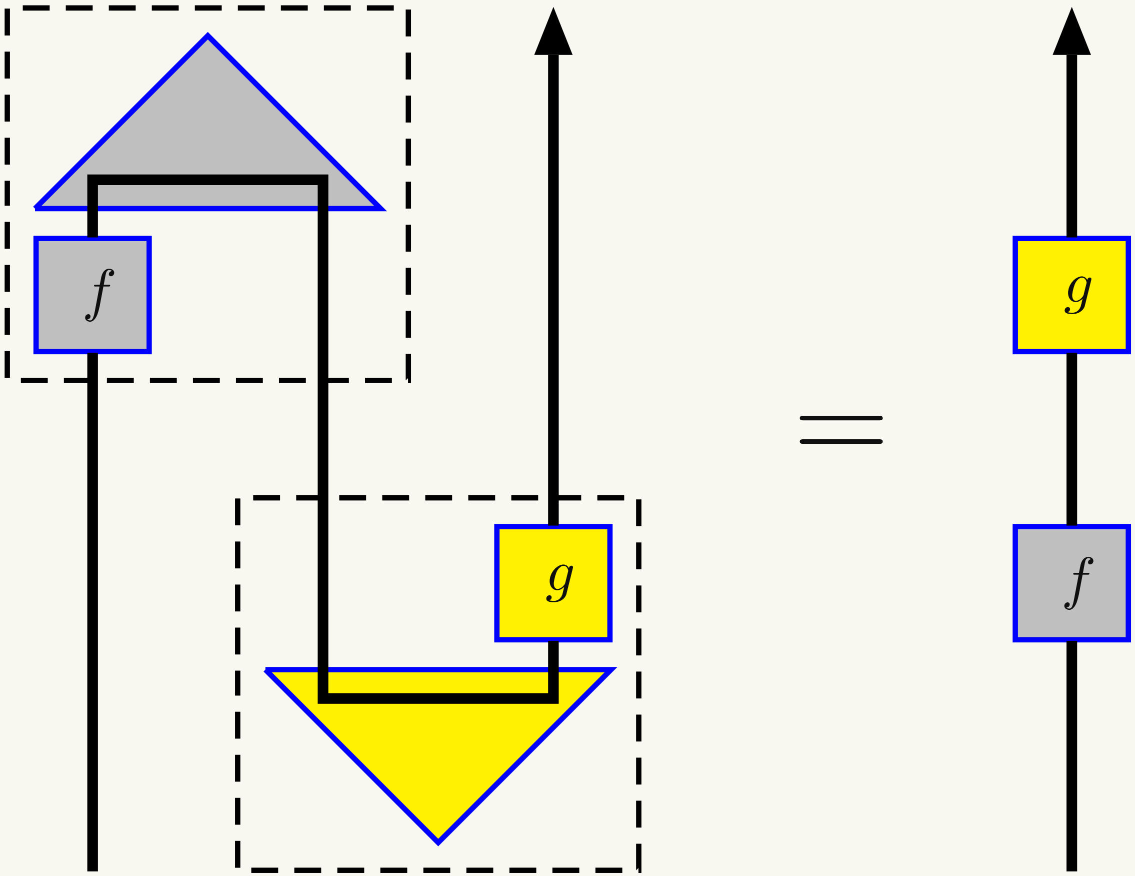
In graphical notation:

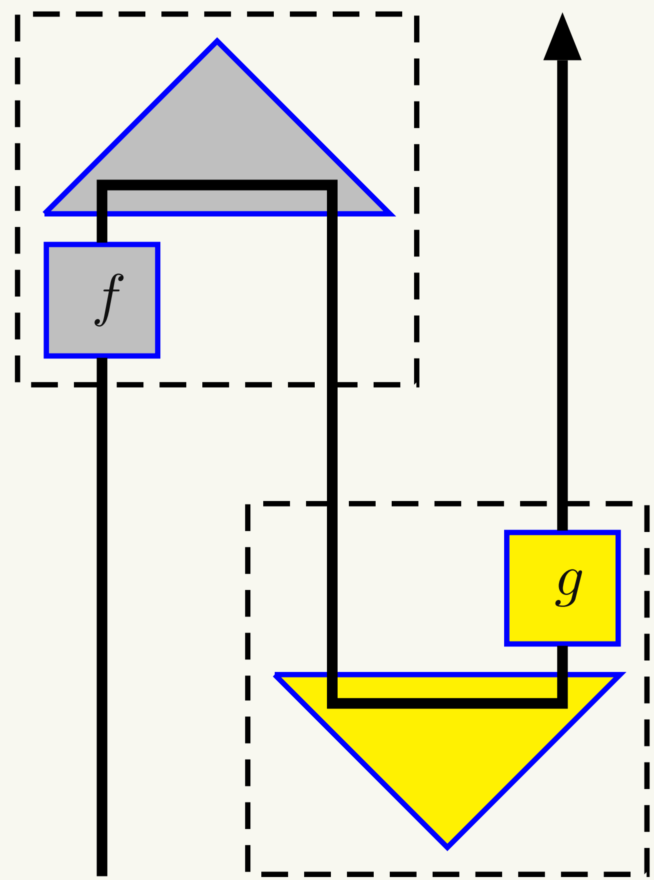


Projectors Decomposed



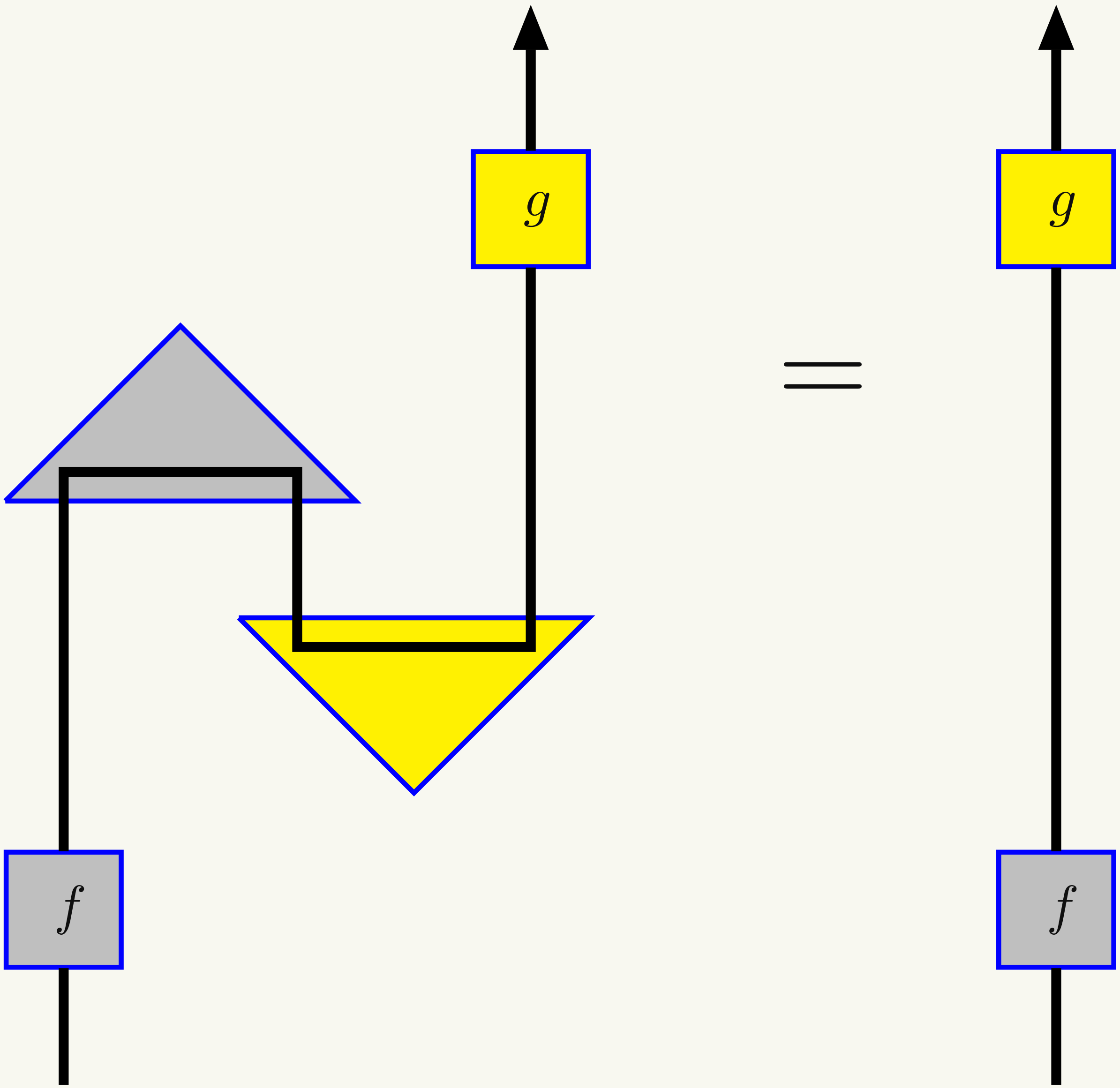
The key algebraic fact from which teleportation (and many other protocols) can be derived.

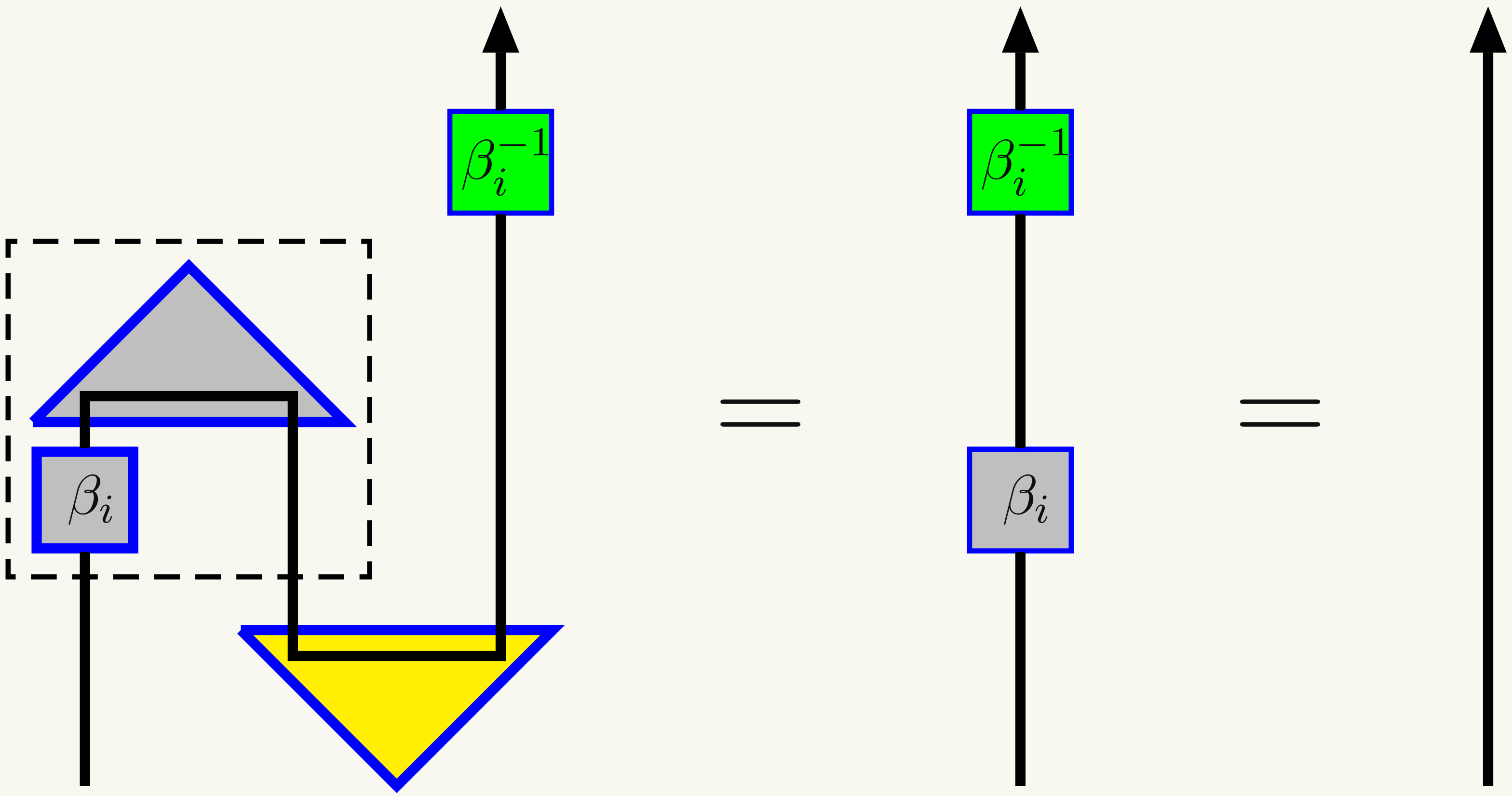



f
g

$=$

Compositionality ctd




β_i^{-1}
β_i
=
β_i^{-1}
β_i
=

Measurements and Classical Information

Suppose we assume biproducts (need only assume products: Robin Houston) in a dagger compact category.

These can be used to represent **branching on measurement outcomes**:

$$A \xrightarrow{\quad M \quad} \bigoplus_{i=1}^{n} A_i \xrightarrow{\bigoplus_{i=1}^{n} U_i} \bigoplus_{i=1}^{n} B_i$$

Propagation of the outcome of a measurement performed on one part of a compound system to other parts — "classical communication" — can be expressed using **distributivity**:

$$(A_1 \oplus A_2) \otimes B \;\cong\; (A_1 \otimes B) \oplus (A_2 \otimes B).$$

$$
\begin{tikzcd}
Q \arrow[r, "(1\otimes\eta)", "(1)"'] \arrow[d, equal] & Q \otimes Q \otimes Q \arrow[r, "\langle \llcorner U_i \lrcorner \rangle_{i=1}^{4} \otimes 1", "(2)"'] & \bigoplus_{i=1}^{4} I \otimes Q \arrow[r, "\text{dist}", "(3)"'] & \bigoplus_{i=1}^{4} Q \arrow[r, "\bigoplus_{i=1}^{4} U_i^{-1}", "(4)"'] & \bigoplus_{i=1}^{4} Q \arrow[d, equal] \\
Q \arrow[rrrr, "\langle 1 \rangle_{i=1}^{i=4}"'] & & & & \bigoplus_{i=1}^{4} Q
\end{tikzcd}
$$

(1) **Produce EPR pair**

(2) **Perform measurement in Bell-basis**

(3) **Propagate classical information**

(4) **Perform unitary correction**.

Teleportation categorically

$$\begin{tikzcd}[column sep=large]
Q \arrow[r, "(1\otimes\eta)", "(1)"'] \arrow[d, equal] & Q \otimes Q \otimes Q \arrow[r, "\langle \llcorner U_i \lrcorner \rangle_{i=1}^{4} \otimes 1", "(2)"'] & \bigoplus_{i=1}^{4} I \otimes Q \arrow[r, "\text{dist}", "(3)"'] & \bigoplus_{i=1}^{4} Q \arrow[r, "\bigoplus_{i=1}^{4} U_i^{-1}", "(4)"'] & \bigoplus_{i=1}^{4} Q \arrow[d, equal] \\
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N.B. Alternative approach (which brings important new structure to light): **classical structures** (Coecke and Pavlovic), *i.e.* Frobenius dagger-algebras.

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- Distributed QM, “discrete QFT”??