Higher algebra in quantum information theory

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March 9, 2018

• Part I: Shaded tensor networks & biunitaries

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Based on joint work with Jamie Vicary:

Biunitary constructions in quantum information

Shaded tangles for the design and verification of quantum programs

Part 1 Shaded tensor networks & biunitaries

Let's start with a very concrete problem.

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Hadamard matrices H

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 U_i unitary $\operatorname{Tr}(U_i^{\dagger}U_j) = n\delta_{i,j}$ $\begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix}, \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}, \begin{pmatrix} 0 & -i \\ 1 & 0 \end{pmatrix}, \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}$

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What is higher algebra?

• Ordinary algebra lets us compose along a line:

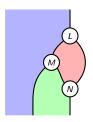
$$xy^2zyx^3$$

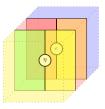
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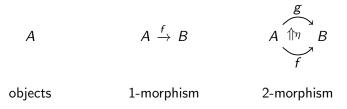
• Higher algebra lets us compose in higher dimensions:





Planar algebra = 2-category theory

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$$A \hspace{1cm} A \xrightarrow{f} B \hspace{1cm} A \xrightarrow{f} B$$
 objects 1-morphism 2-morphism

We can compose 2-morphisms like this:

$$A \xrightarrow{\uparrow \epsilon} B \qquad A \xrightarrow{\uparrow \eta} B \xrightarrow{\uparrow \epsilon} C$$

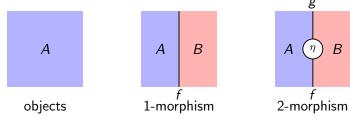
vertical composition

horizontal composition

These are pasting diagrams.

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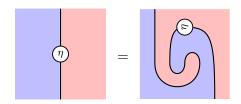
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The dual diagrams are the graphical calculus.

We use monoidal dagger pivotal 2-categories:

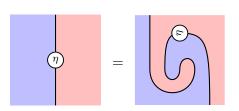
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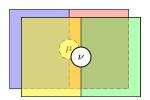
• Dagger pivotal 2-categories have a very flexible graphical calculus.



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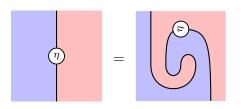
- Dagger pivotal 2-categories have a very flexible graphical calculus.
- In a monoidal 2-category, we can layer surfaces on top of each other.

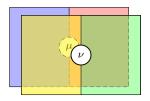




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⇒ surfaces, wires and vertices in three-dimensional space

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$$\left(\begin{array}{ccc}
H_{11} & \cdots & H_{1n} \\
\vdots & \ddots & \vdots \\
H_{m1} & \cdots & H_{mn}
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- 2-morphisms $H \stackrel{\phi}{\Rightarrow} H'$ are matrices of linear maps

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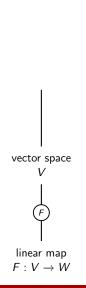
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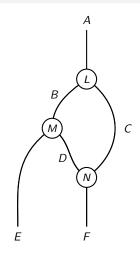
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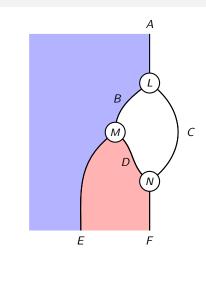
This well-studied structure plays a key role in higher representation theory.

A direct perspective: tensor networks





A (composed) linear map $E \otimes F \rightarrow A$



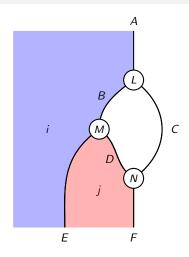




indexing set $i \in S$







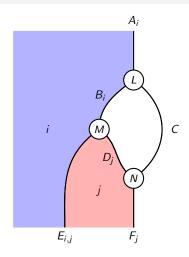


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family of vector spaces $V_{i,j}$







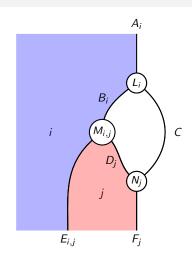
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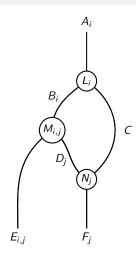
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A family of linear maps, indexed by i and j $E_{i,j} \otimes F_j \to A_i$



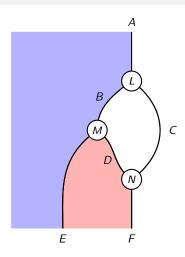
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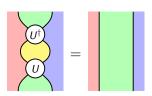


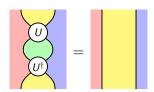
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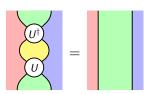
• (vertically) unitary:

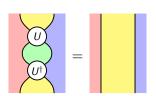




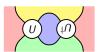
A biunitary is a 2-morphism that is

• (vertically) unitary:





• horizontally unitary:



 $=\lambda$



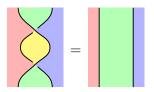


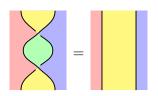




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These look just like the second Reidemeister move.

Quantum structures are biunitaries in 2Hilb

Result 1: Hadamards and UEBs are biunitaries of the following type:

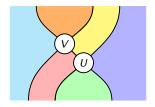


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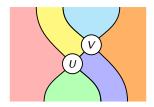


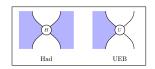
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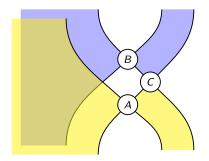
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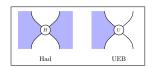


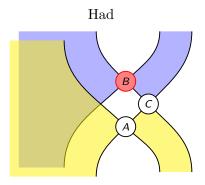
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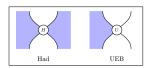


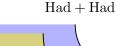


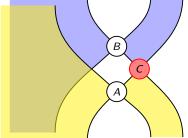


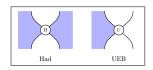




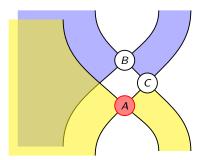


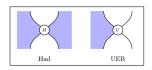




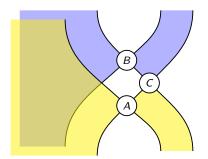


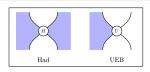
Had + Had + Had



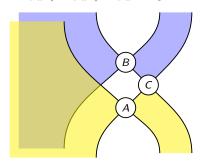


 $Had + Had + Had \rightsquigarrow UEB$

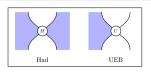




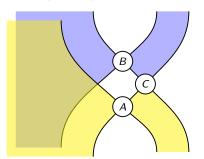
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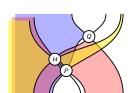
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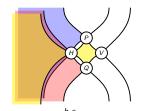
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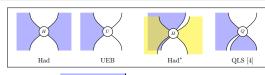


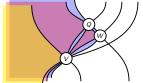
Composing biunitaries



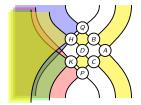
 $U_{abc,de,fg} = H_{a,eg}^{b,c} P_{e,b,f}^{c,g} Q_{c,g,d}$





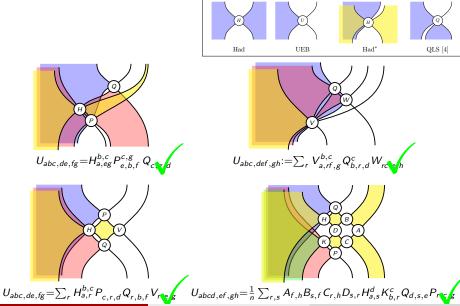


 $U_{abc,def,gh} := \sum_{r} V_{a,rf,g}^{b,c} Q_{b,r,d}^{c} W_{rc,e,h}$



 $U_{abc,de,fg} = \sum_{r} H_{a,r}^{b,c} P_{c,r,d} Q_{r,b,f} V_{r,e,g} \quad U_{abcd,ef,gh} = \frac{1}{n} \sum_{r,s} A_{f,h} B_{s,f} C_{r,h} D_{s,r} H_{a,s}^{d} K_{b,r}^{c} Q_{d,s,e} P_{r,c,g}$

Composing biunitaries



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 see structural properties hidden in conventional matrix notation

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~→



Part 2 Untangling quantum circuits

Basic states and gates





$$|+\rangle = |0\rangle + |1\rangle$$

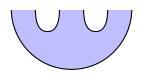
$$|\mathrm{Bell}\rangle = |00\rangle + |11\rangle$$

$$|+\rangle = |0\rangle + |1\rangle \qquad |\mathrm{Bell}\rangle = |00\rangle + |11\rangle \qquad |\mathrm{GHZ}\rangle = |000\rangle + |111\rangle$$

Basic states and gates







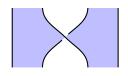
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$$|i\rangle\mapsto\sum_{j}H_{ij}|j\rangle$$



$$|i\rangle \otimes |j\rangle \mapsto H_{ij} |i\rangle \otimes |j\rangle$$

Basic states and gates





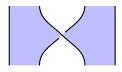
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Hadamard gate



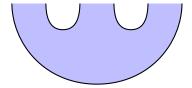
CZ gate

Creating GHZ states

How to create a GHZ state from $|+\rangle$ states?

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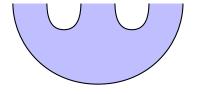




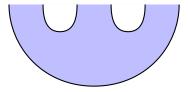


Creating GHZ states

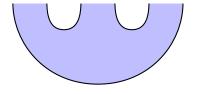
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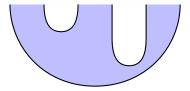
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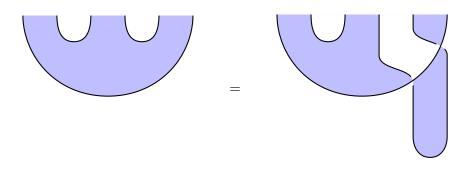


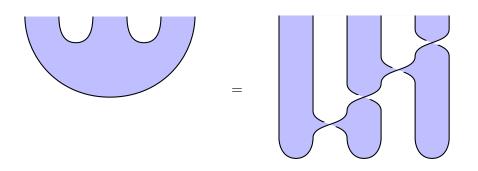
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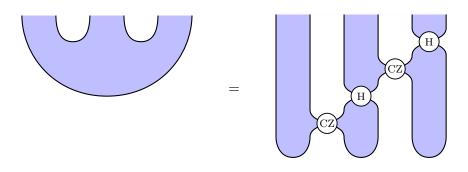


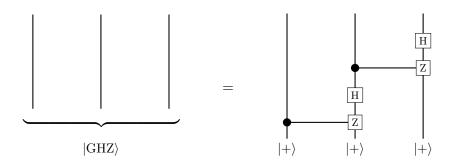
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Quantum error correction

A k-local quantum code is an isometry $H \stackrel{\text{enc}}{\longrightarrow} H^{\otimes n}$, s.t.

$$H \stackrel{\mathsf{enc}}{\longrightarrow} H^{\otimes n} \stackrel{E}{\longrightarrow} H^{\otimes n} \stackrel{\mathsf{enc}^\dagger}{\longrightarrow} H$$

is proportional to the identity for every k-local error $E: H^{\otimes n} \to H^{\otimes n}$.

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phase error

Quantum error correction

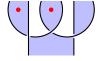
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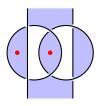
$$H \stackrel{\mathsf{enc}}{\longrightarrow} H^{\otimes n} \stackrel{E}{\longrightarrow} H^{\otimes n} \stackrel{\mathsf{enc}^\dagger}{\longrightarrow} H$$

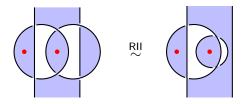
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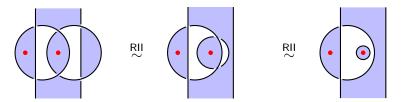


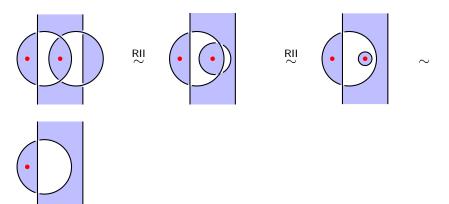


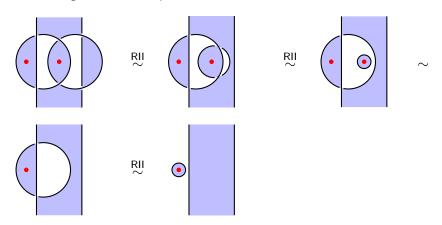


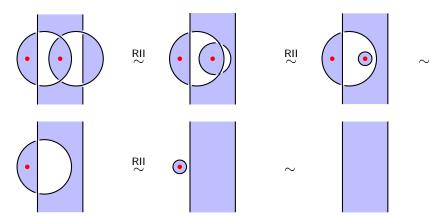




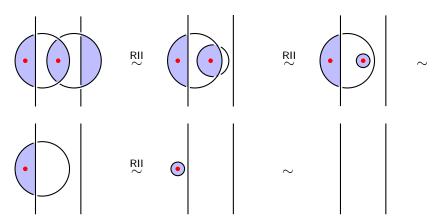






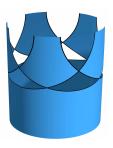


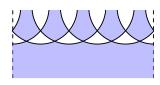
The following is a 2-local phase error code $H \to H^{\otimes 3}$:

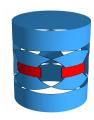


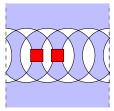
New construction of a phase code from unitary error bases.

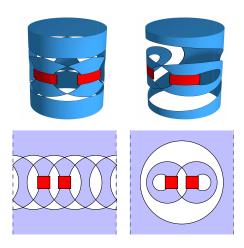
A 2-local full error correcting code $H \to H^{\otimes 5}$:

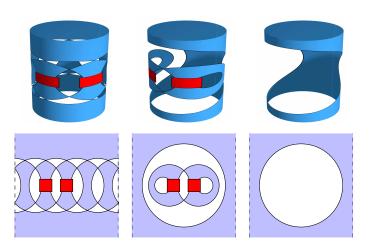


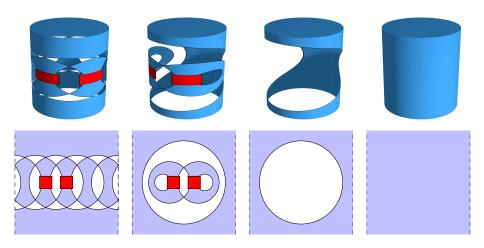


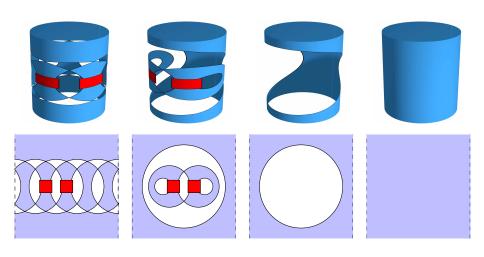




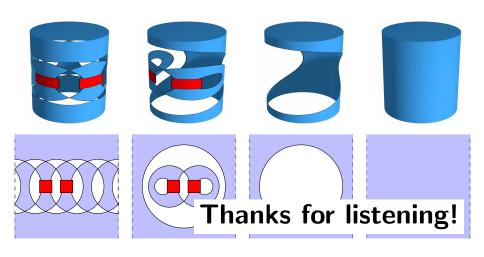








Caveat: We cannot yet handle two non-adjacent errors.



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