

Solving simple PDEs using the finite element method

FEM for simple PDEs: elliptic and parabolic linear PDEs

Second-order PDEs commonly arise in physical models. There are three archetypal second-order PDEs

- 1 Elliptic PDEs, for example, Poisson's equation $\nabla^2 u + f = 0$
- 2 Parabolic PDEs, for example, the heat equation $u_t = \nabla^2 u + f$
- 3 Hyperbolic PDEs, for example, the wave equation $u_{tt} = \nabla^2 u$

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- 2 **Parabolic PDEs, for example, the heat equation** $u_t = \nabla^2 u + f$
- 3 **Hyperbolic PDEs, for example, the wave equation** $u_{tt} = \nabla^2 u$

Defining PDEs in an object oriented manner

First, an abstract class defining a general linear elliptic PDE $\nabla \cdot D \nabla u = f$, where D is a matrix-valued function of position (the diffusion tensor):

AbstractLinearEllipticPde:

Abs. method: GetDiffusionTensor(x)

Abs. method: GetForceTerm(x)

MyEllipticPde: *inherits from AbstractLinearEllipticPde*

Implemented method: GetDiffusionTensor(x)

Implemented method: GetForceTerm(x)

For example $\nabla^2 u = 0$

LaplacesEquation: *inherits from AbstractLinearEllipticPde*

Implemented method: GetDiffusionTensor(x)

▷ *return identity matrix*

Implemented method: GetForceTerm(x)

▷ *return zero*

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Defining PDEs in an object oriented manner

Next, an abstract class defining a general linear parabolic PDE

$$\alpha u_t = \nabla \cdot D \nabla u + f$$

where α , D and f are functions of space and time.

AbstractLinearParabolicPde:

Abs. method: `GetDuDtCoefficientTerm(t,x)`

Abs. method: `GetDiffusionTensor(t,x)`

Abs. method: `GetForceTerm(t,x)`

For example $u_t = \nabla^2 u$

HeatEquation: inherits from `AbstractLinearParabolicPde`

Implemented method: `GetDuDtCoefficientTerm(t,x)`

▷ return 1

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▷ return identity matrix

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FEM for simple PDEs: introduction to FEM

The finite element method

Stages

- 1 Convert equation from **strong form** to **weak form**
- 2 Convert infinite-dimensional problem into a finite dimensional one
- 3 Set up the finite element linear system to be solved

Weak form of Poisson's equation

Consider Poisson's equation:

$$\nabla^2 u + f = 0$$

subject to boundary conditions

$$\begin{aligned} u &= 0 && \text{on } \Gamma_1 \\ \nabla u \cdot \mathbf{n} &= g && \text{on } \Gamma_2 \end{aligned}$$

Weak form

Multiply by a test function v satisfying $v = 0$ on Γ_1 , and integrate:

$$\begin{aligned} v(\nabla^2 u) &= -fv \\ \int_{\Omega} v(\nabla^2 u) \, dV &= -\int_{\Omega} fv \, dV \\ \int_{\partial\Omega} v(\nabla u \cdot \mathbf{n}) \, dS - \int_{\Omega} \nabla u \cdot \nabla v \, dV &= -\int_{\Omega} fv \, dV \\ \int_{\Omega} \nabla u \cdot \nabla v \, dV &= \int_{\Omega} fv \, dV + \int_{\Gamma_2} gv \, dS \end{aligned}$$

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Weak form of Poisson's equation

Let \mathcal{V} be the space of all differentiable functions on Ω (more precisely, \mathcal{V} is the Sobolev space $H^1(\Omega)$). Let

$$\mathcal{V}_0 = \{v \in \mathcal{V} : v = 0 \text{ on } \Gamma_1\}$$

Weak form

Find $u \in \mathcal{V}_0$ satisfying

$$\int_{\Omega} \nabla u \cdot \nabla v \, dV = \int_{\Omega} f v \, dV + \int_{\Gamma_2} g v \, dS \quad \forall v \in \mathcal{V}_0$$

Example

Solve $\frac{d^2 u}{dx^2} = 1$, $u(0) = u(1) = 0$ vs

Find differentiable u satisfying

$u(0) = u(1) = 0$ and:

$\int_0^1 \frac{du}{dx} \frac{dv}{dx} \, dx = - \int_0^1 v \, dx$ for all
 v s.t. $v(0) = v(1) = 0$

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FEM discretisation

Find $u \in \mathcal{V}_0$ satisfying

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Take

$$\mathcal{V}_0^h = \text{span}\{\phi_1, \phi_2\}$$

(where ϕ_1, ϕ_2 satisfy the Dirichlet boundary conditions), so

$$u_h = \alpha \phi_1 + \beta \phi_2$$

Linear system:

$$\begin{bmatrix} \int_{\Omega} \nabla \phi_1 \cdot \nabla \phi_1 \, dV & \int_{\Omega} \nabla \phi_1 \cdot \nabla \phi_2 \, dV \\ \int_{\Omega} \nabla \phi_2 \cdot \nabla \phi_1 \, dV & \int_{\Omega} \nabla \phi_2 \cdot \nabla \phi_2 \, dV \end{bmatrix} \begin{bmatrix} \alpha \\ \beta \end{bmatrix} = \begin{bmatrix} \int_{\Omega} f \phi_1 \, dV + \int_{\Gamma_2} g \phi_1 \, dS \\ \int_{\Omega} f \phi_2 \, dV + \int_{\Gamma_2} g \phi_2 \, dS \end{bmatrix}$$

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FEM discretisation

Find $u_h \in \mathcal{V}_0^h$ satisfying

$$\int_{\Omega} \nabla u_h \cdot \nabla \phi_j \, dV = \int_{\Omega} f \phi_j \, dV + \int_{\Gamma_2} g \phi_j \, dS \quad \text{for } j = 1, 2$$

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FEM discretisations

Take

$$V_h = \text{span}\{\phi_1, \phi_2, \dots, \phi_N\}$$

(satisfying $\phi_j = 0$ on Γ_1) so

$$u_h = \alpha_1 \phi_1 + \dots + \alpha_N \phi_N$$

Let the stiffness matrix and RHS vector be given by

$$K_{jk} = \int_{\Omega} \nabla \phi_j \cdot \nabla \phi_k \, dV$$
$$b_j = \int_{\Omega} f \phi_j \, dV + \int_{\Gamma_2} g \phi_j \, dS$$

and solve

$$K \begin{bmatrix} \alpha_1 \\ \vdots \\ \alpha_N \end{bmatrix} = \mathbf{b}$$

Basis functions

FEM discretisations

Let

$$K_{jk} = \int_{\Omega} \nabla \phi_j \cdot \nabla \phi_k \, dV \quad \text{stiffness matrix}$$

$$M_{jk} = \int_{\Omega} \phi_j \phi_k \, dV \quad \text{mass matrix}$$

$$b_j = \int_{\Omega} f \phi_j \, dV + \int_{\Gamma_2} g \phi_j \, dS$$

FEM discretisations

$$\text{Laplace's equation: } \nabla^2 u + f = 0 \quad \rightarrow \quad K\mathbf{U} = \mathbf{b}$$

Heat equation:

$$\frac{\partial u}{\partial t} = \nabla^2 u + f \quad \rightarrow \quad M \frac{d\mathbf{U}}{dt} + K\mathbf{U} = \mathbf{b}$$

Time-discretised heat equation:

$$\frac{u^{n+1} - u^n}{\Delta t} = \nabla^2 u^{n+1} + f^{n+1} \quad \rightarrow \quad M\mathbf{U}^{n+1} + \Delta t K\mathbf{U}^{n+1} = M\mathbf{U}^n + \Delta t \mathbf{b}^{n+1}$$

Anisotropic diffusion

Suppose we have an anisotropic diffusion tensor D (symmetric, positive definite), for example, in Poisson's equation:

$$\nabla \cdot (D \nabla u) + f = 0$$

subject to boundary conditions

$$\begin{aligned} u &= 0 && \text{on } \Gamma_1 \\ (D \nabla u) \cdot \mathbf{n} &= g && \text{on } \Gamma_2 \end{aligned}$$

The weak form is: find $u \in \mathcal{V}_0$ satisfying

$$\int_{\Omega} (D \nabla u) \cdot \nabla v \, dV = \int_{\Omega} f v \, dV + \int_{\Gamma_2} g v \, dS \quad \forall v \in \mathcal{V}_0$$

and the only change in the FEM discretisation is that the stiffness matrix becomes

$$K_{jk} = \int_{\Omega} \nabla \phi_j \cdot (D \nabla \phi_k) \, dV$$

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Implementing Dirichlet boundary conditions

In practice, rather using the basis functions in \mathcal{V}_0^h (i.e. bases satisfying $\phi_i = 0$ on Γ_1), we use \mathcal{V}^h , i.e. all the basis functions corresponding to all nodes in the mesh.

We then impose (any) Dirichlet boundary conditions by altering the appropriate rows of the linear system, for example, for $K\mathbf{U} = \mathbf{b}$, if we want to impose $U_1 = c$

$$\begin{bmatrix} K_{11} & K_{12} & \dots & K_{1N} \\ K_{21} & K_{22} & \dots & K_{2N} \\ \vdots & \vdots & \ddots & \vdots \\ K_{N1} & K_{N2} & \dots & K_{NN} \end{bmatrix} \begin{bmatrix} U_1 \\ U_2 \\ \vdots \\ U_N \end{bmatrix} = \begin{bmatrix} b_1 \\ b_2 \\ \vdots \\ b_N \end{bmatrix}$$

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FEM stages

Solve:

$$\nabla \cdot (D\nabla u) + f = 0$$

subject to boundary conditions

$$\begin{aligned} u &= u^* && \text{on } \Gamma_1 \\ (D\nabla u) \cdot \mathbf{n} &= g && \text{on } \Gamma_2 \end{aligned}$$

- 1 Set up the computational mesh and choose basis functions
- 2 Compute the matrix K and vector \mathbf{b} :

$$\begin{aligned} K_{jk} &= \int_{\Omega} \nabla \phi_j \cdot (D\nabla \phi_k) \, dV \\ b_j &= \int_{\Omega} f \phi_j \, dV + \int_{\Gamma_2} g \phi_j \, dS \end{aligned}$$

- 3 Alter linear system $K\mathbf{U} = \mathbf{b}$ to impose Dirichlet BCs
- 4 Solve linear system

FEM for simple PDEs: FEM details

Computing a finite element matrix/vector by *assembly*

Consider computing the mass matrix $M_{jk} = \int_{\Omega} \phi_j \phi_k \, dV$, an N by N matrix say, and let's suppose (for clarity only) that we are in 2D. Also, assume we are using linear basis functions.

We **do not** write out the full basis functions explicitly in computing this integral. Instead: firstly, we break the integral down into an integral over elements:

$$M_{jk} = \sum_{\mathcal{K}} \int_{\mathcal{K}} \phi_j \phi_k \, dV$$

Consider $\int_{\mathcal{K}} \phi_j \phi_k \, dV$. **Key point:** The only basis functions with are non-zero in the triangle are the 3 basis functions corresponding to the 3 nodes of the element.

Therefore: compute the **elemental contribution to the mass matrix**, a **3 by 3** matrix of the form $\int_{\mathcal{K}} \phi_j \phi_k \, dV$ for **3 choices of j and k only**.

Then add elemental contribution to full N by N mass matrix.

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Computing an elemental contribution

We have reduced the problem to computing small matrices/vectors, for example the 3 by 3 matrix

$$\int_{\mathcal{K}} \phi_j \phi_k \, dV$$

where ϕ_j, ϕ_k are the 3 basis functions corresponding to the 3 nodes of the mesh.

Next, map to the **reference triangle** (also known as the canonical triangle), \mathcal{K}_{ref} , the triangle with nodes $(0, 0)$, $(0, 1)$, $(1, 0)$.

The basis functions on the reference triangle are easy to write down

$$N_1(\xi, \eta) = 1 - \xi - \eta$$

$$N_2(\xi, \eta) = \xi$$

$$N_3(\xi, \eta) = \eta$$

Computing an elemental contribution

We now need to be able to compute

$$\int_{\mathcal{K}} \phi_j \phi_k \, dx dy = \int_{\mathcal{K}_{\text{ref}}} N_j N_k \det J \, d\xi d\eta$$

where J is the Jacobian of the mapping from the true element to the canonical element.

J is also required if $\nabla \phi_i$ is needed (for example, in computing the stiffness matrix), since $\nabla \phi_i = J \nabla_{\xi} N_i$.

Consider the mapping from an element with nodes $\mathbf{x}_1, \mathbf{x}_2, \mathbf{x}_3$, to the canonical element. The inverse mapping can in fact be easily written down using the basis functions.

$$\mathbf{x}(\xi, \eta) = \sum_{j=1}^3 \mathbf{x}_j N_j(\xi, \eta)$$

from which it is easy to show that J is the following function of nodal positions

$$J = \text{inv} \begin{bmatrix} x_2 - x_1 & x_3 - x_1 \\ y_2 - y_1 & y_3 - y_1 \end{bmatrix}$$

Computing an elemental contribution

We now need to be able to compute

$$\int_{\mathcal{K}} \phi_j \phi_k \, dx dy = \int_{\mathcal{K}_{\text{ref}}} N_j N_k \det J \, d\xi d\eta$$

where J is the Jacobian of the mapping from the true element to the canonical element.

J is also required if $\nabla \phi_i$ is needed (for example, in computing the stiffness matrix), since $\nabla \phi_i = J \nabla_{\xi} N_i$.

Consider the mapping from an element with nodes x_1, x_2, x_3 , to the canonical element. The inverse mapping can in fact be easily written down using the basis functions.

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Computing an elemental contribution - the general case

Suppose we want to compute

$$\int_{\mathcal{K}} \mathcal{F}(x, y, u, \phi_1, \phi_2, \phi_3, \nabla \phi_1, \nabla \phi_2, \nabla \phi_3) dx dy$$

We map to the reference element:

$$\int_{\mathcal{K}_{\text{ref}}} \mathcal{F}(x, y, u, \phi_1, \phi_2, \phi_3, \nabla \phi_1, \nabla \phi_2, \nabla \phi_3) \det J d\xi d\eta$$

and then use **numerical quadrature**, which means f just has to be evaluated at the quadrature points.

FEM stages - full algorithm

Solve:

$$\nabla \cdot (D\nabla u) + f = 0$$

subject to boundary conditions

$$\begin{aligned} u &= u^* && \text{on } \Gamma_1 \\ (D\nabla u) \cdot \mathbf{n} &= g && \text{on } \Gamma_2 \end{aligned}$$

- 1 Set up the computational mesh and choose basis functions
- 2 Compute the matrix K and vector \mathbf{b} :

$$\begin{aligned} K_{jk} &= \int_{\Omega} \nabla \phi_j \cdot (D\nabla \phi_k) \, dV \\ b_j &= \int_{\Omega} f \phi_j \, dV + \int_{\Gamma_2} g \phi_j \, dS \end{aligned}$$

- 3 Alter linear system $K\mathbf{U} = \mathbf{b}$ to impose Dirichlet BCs
- 4 Solve linear system

FEM stages - full algorithm

Write

$$b_j = \int_{\Omega} f \phi_j dV + \int_{\Gamma_2} g \phi_j dS$$

as $\mathbf{b} = \mathbf{b}^{\text{vol}} + \mathbf{b}^{\text{surf}}$

- 1 Set up the computational mesh and choose basis functions
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