

Quantum Recursion and Second Quantisation

Basic Ideas and Examples

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Happy Birthday Prakash!

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I'm very grateful to Prakash for teaching me second quantisation method during his visit to UTS in 2013

Outline

1. Introduction
2. Quantum Case Statement and Quantum Choice
3. Motivating Example: Recursive Quantum Walks
4. Second Quantisation
5. Semantics of Quantum Recursion
7. Conclusion

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Classical Recursion of Quantum Programs

- ▶ Recursive procedure in quantum programming language QPL [Selinger, *Mathematical Structures in Computer Science*'2004].

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- ▶ Termination of quantum while-loops in finite-dimensional state spaces [Ying, Feng, *Acta Informatica*'2010].
- ▶ Selinger's slogan: *Quantum data, classical control* - control flow of the quantum recursions is classical because branchings are determined by the outcomes of certain quantum measurements.

Quantum Control Flow

Quantum programming language QML [Altenkirch and Grattage, *LICS'2005*]:

Two case constructs in the quantum setting:

- ▶ **case**, measure a qubit in the data it analyses - The control flow is determined by the outcome of a measurement and thus is classical.

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- ▶ **case**, measure a qubit in the data it analyses - The control flow is determined by the outcome of a measurement and thus is classical.
- ▶ **case^o**, analyse quantum data without measuring - **if^o** – **then** – **else** statement.

Quantum Control Flow

Hadamard gate:

$$\begin{aligned} \text{had } x = & \text{if}^\circ x \\ & \text{then } \left\{ \frac{1}{\sqrt{2}}(q_{\text{false}} - q_{\text{true}}) \right\} \\ & \text{else } \left\{ \frac{1}{\sqrt{2}}(q_{\text{false}} + q_{\text{true}}) \right\} \end{aligned}$$

Quantum Control Flow

NOT gate:

```
not  $x$  = ifo  $x$   
      then  $q_{\text{false}}$   
      else  $q_{\text{true}}$ 
```

CNOT gate:

```
cnot  $c$   $x$  = ifo  $c$   
      then ( $q_{\text{true}}, \textit{not } x$ )  
      else ( $q_{\text{false}}, x$ )
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The state Hilbert space $\mathcal{H}_c = \text{span}\{|0\rangle, |1\rangle\}$
- ▶ U_0 and U_1 two unitary transformations on a quantum system q - the state Hilbert space $\mathcal{H}_q$.
- ▶ A *quantum case statement* employing “quantum coin” c :

qif $[c] \ |0\rangle \rightarrow U_0[q]$

$\square \ |1\rangle \rightarrow U_1[q]$

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“coined” Quantum Case Statement

- ▶ The semantics is an unitary operator U in $\mathcal{H}_c \otimes \mathcal{H}_q$ - the state Hilbert space of the composed system of “coin” c and principal system q :

$$U|0, \psi\rangle = |0\rangle U_0|\psi\rangle, \quad U|1, \psi\rangle = |1\rangle U_1|\psi\rangle$$

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- ▶ Matrix representation:

$$U = |0\rangle\langle 0| \otimes U_0 + |1\rangle\langle 1| \otimes U_1 = \begin{pmatrix} U_0 & 0 \\ 0 & U_1 \end{pmatrix}.$$

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- ▶ The quantum choice of $U_0[q]$ and $U_1[q]$ with “coin-tossing” $V[c]$:

$$U_0[q] \oplus_{V[c]} U_1[q]$$

$\stackrel{\text{def}}{=}$

$$V[c]; \text{ \textbf{qif} } [c] \begin{array}{l} |0\rangle \rightarrow U_0[q] \\ \square \quad |1\rangle \rightarrow U_1[q] \end{array}$$

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- ▶ Compare with probabilistic choice [McIver and Morgan, *Abstraction, Refinement and Proof for Probabilistic Systems*, 2005]

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- ▶ Quantum walks [Ambainis, Bach, Nayak, Vishwanath, Watrous, *STOC'2001*; Aharonov, Ambainis, Kempe, Vazirani, *STOC'2001*].
- ▶ Unitary transformations $U_0[q], U_1[q]$ are replaced by general quantum programs that may contain quantum measurements? [Ying, Yu, Feng, arXiv:1209.4379]

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Example - One-dimensional quantum walk

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 - ▶ $\mathcal{H}_p = \text{span}\{|n\rangle : n \in Z\}$, n indicates the position marked by integer n .

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▶

$$H = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 & 1 \\ 1 & -1 \end{pmatrix}$$

is Hadamard transform in the direction space $\mathcal{H}_d$

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- ▶ Define the left and right translation operators T_L and T_R in the position space $\mathcal{H}_p$:

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- Hadamard walk — *repeated applications* of operator W .

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- ▶ Recursive Hadamard walk — program X declared by the recursive equation:

$$X \Leftarrow T_L[p] \oplus_{H[d]} (T_R[p]; X)$$

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- ▶ A variant of the bidirectional recursive Hadamard walk is declared by the following system of recursive equations:

$$\begin{cases} X \Leftarrow T_L[p] \oplus_{H[d]} (T_R[p]; Y), \\ Y \Leftarrow (T_L[p]; X) \oplus_{H[d]} T_R[p] \end{cases}$$

A More Interesting Recursive Quantum Walk

- ▶ Let $n \geq 2$. Another variant of unidirectional recursive quantum walk is defined as the program declared by the following recursive equation:

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How to solve these quantum recursive equations?

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- ▶ Semantics $\llbracket X \rrbracket$ of X is the limit

$$\llbracket X \rrbracket = \lim_{n \rightarrow \infty} \llbracket X^{(n)} \rrbracket$$

Example - (Unidirectional) Recursive Hadamard Walk

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$$X^{(1)} = T_L[p] \oplus_{H[d]} (T_R[p]; \mathbf{abort}),$$

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Observation

- ▶ Continuously introduce new “coin” to avoid variable conflict.
- ▶ Variables $d, d_1, d_2, \dots$ denote identical particles.

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- ▶ Variables $d, d_1, d_2, \dots$ denote identical particles.
- ▶ The number of the “coin” particles that are needed in running the recursive walk is unknown beforehand.
- ▶ We need to deal with *quantum systems where the number of particles of the same type may vary*.

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Fock Spaces

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- ▶ For each permutation π of $1, \dots, n$, define the permutation operator P_π in $\mathcal{H}^{\otimes n}$ by

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- ▶ Define the symmetrisation and antisymmetrisation operators in $\mathcal{H}^{\otimes n}$:

$$S_+ = \frac{1}{n!} \sum_{\pi} P_\pi, \quad S_- = \frac{1}{n!} \sum_{\pi} (-1)^\pi P_\pi$$

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- ▶ The free Fock space:

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$$\mathbf{U}|\psi_1, \dots, \psi_n\rangle_v = |U\psi_1, \dots, U\psi_n\rangle_v.$$

Evolution in the Fock Spaces

- ▶ Let the (discrete-time) evolution of one particle be unitary operator U .
- ▶ The evolution of n particles without mutual interactions is operator $\mathbf{U}$ in $\mathcal{H}^{\otimes n}$:

$$\mathbf{U}|\psi_1 \otimes \dots \otimes \psi_n\rangle = |U\psi_1 \otimes \dots \otimes U\psi_n\rangle$$

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$$\mathbf{U}|\psi_1, \dots, \psi_n\rangle_v = |U\psi_1, \dots, U\psi_n\rangle_v.$$

- ▶ Extend to the Fock spaces $\mathcal{F}_v(\mathcal{H})$ and $\mathcal{F}(\mathcal{H})$:

$$\mathbf{U} \left(\sum_{n=0}^{\infty} |\Psi(n)\rangle \right) = \sum_{n=0}^{\infty} \mathbf{U} |\Psi(n)\rangle$$

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- ▶ Annihilation operator $a(\psi)$ — the Hermitian conjugate of $a^*(\psi)$:

$$a(\psi)|\mathbf{0}\rangle = 0,$$

$$a(\psi)|\psi_1, \dots, \psi_n\rangle_v = \frac{1}{\sqrt{n}} \sum_{i=1}^n (v)^{i-1} \langle \psi | \psi_i \rangle |\psi_1, \dots, \psi_{i-1}, \psi_{i+1}, \dots, \psi_n\rangle_v$$

Decrease the number of particles by one unit, while preserving the symmetry of the state.

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Example - (Unidirectional) Recursive Hadamard Walk

Semantics of the recursive Hadamard walk:

$$\llbracket X \rrbracket = \left[\sum_{i=0}^{\infty} \left(\bigotimes_{j=0}^{i-1} |R\rangle_{d_j} \langle R| \otimes |L\rangle_{d_i} \langle L| \right) \otimes T_L T_R^i \right] (\mathbf{H} \otimes I)$$

- ▶ An operator in

$$\mathcal{F}_v(\mathcal{H}_d) \otimes \mathcal{H}_p \rightarrow \mathcal{F}(\mathcal{H}_d) \otimes \mathcal{H}_p.$$

- ▶ The sign v is $+$ or $-$, depending on using bosons or fermions to implement the “direction coins” $d, d_1, d_2, \dots$

Principal System Semantics

- ▶ Each state $|\Psi\rangle$ in Fock space $\mathcal{F}_v(\mathcal{H}_d)$ induces mapping:

$\llbracket X, \Psi \rrbracket_p : \text{pure states} \rightarrow \text{partial density operators in } \mathcal{H}_p$

$$\llbracket X, \Psi \rrbracket_p(|\psi\rangle) = \text{tr}_{\mathcal{F}(\mathcal{H}_d)}(|\Phi\rangle\langle\Phi|)$$

where $|\Phi\rangle = \llbracket X \rrbracket(|\Psi\rangle \otimes |\psi\rangle)$

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where $|\Phi\rangle = \llbracket X \rrbracket(|\Psi\rangle \otimes |\psi\rangle)$

- ▶ Mapping $\llbracket X, \Psi \rrbracket_p$ is called the principal system semantics of X with “coin” initialisation $|\Psi\rangle$.

Bidirectional Recursive Quantum Walk

$$\begin{cases} X \Leftarrow T_L[p] \oplus_{H[d]} (T_R[p]; Y), \\ Y \Leftarrow (T_L[p]; X) \oplus_{H[d]} T_R[p] \end{cases}$$

- Coherent state of bosons in the symmetric Fock space $\mathcal{F}_+(\mathcal{H})$ over $\mathcal{H}$:

$$|\psi\rangle_{\text{coh}} = \exp\left(-\frac{1}{2}\langle\psi|\psi\rangle\right) \sum_{n=0}^{\infty} \frac{[a^*(\psi)]^n}{n!} |\mathbf{0}\rangle$$

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- ▶ The walk starts from position 0 and the coins are initialised in the coherent states of bosons corresponding to $|L\rangle$:

$$\begin{aligned} \llbracket X, L_{\text{coh}} \rrbracket_p(|0\rangle) &= \frac{1}{\sqrt{e}} \left(\sum_{k=0}^{\infty} \frac{1}{2^{2k+1}} |-1\rangle\langle -1| + \sum_{k=0}^{\infty} \frac{1}{2^{2k+2}} |2\rangle\langle 2| \right) \\ &= \frac{1}{\sqrt{e}} \left(\frac{2}{3} |-1\rangle\langle -1| + \frac{1}{3} |2\rangle\langle 2| \right). \end{aligned}$$

Quantum while-loop

- ▶ Program X declared by the recursive equation

$$\begin{aligned} X \Leftarrow & W[c, q]; \mathbf{qif}[c] \, |0\rangle \rightarrow \mathbf{skip} \\ & \square \, |1\rangle \rightarrow U[q]; X \\ & \mathbf{fiq} \end{aligned}$$

where W a unitary operator in $\mathcal{H}_c \otimes \mathcal{H}_q$ — the interaction between the “coin” c and the principal system q .

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- Semantics of X :

$$\llbracket X \rrbracket = \sum_{k=1}^{\infty} \prod_{j=0}^{k-1} W[c_j, q] \left(\bigotimes_{j=0}^{k-2} |1\rangle_{c_j} \langle 1| \otimes |0\rangle_{c_{k-1}} \langle 0| \otimes U^{k-1}[q] \right)$$

from the space $\mathcal{F}_v(\mathcal{H}_2) \otimes \mathcal{H}_q$ into $\mathcal{F}(\mathcal{H}_2) \otimes \mathcal{H}_q$.

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 - ▶ Combine linear logic with Hoare logic for quantum programs [Ying, *TOPLAS*'2011]?
- ▶ What kind of physical systems can be used to implement quantum recursion where new “coins” must be continuously created?

Thank You!