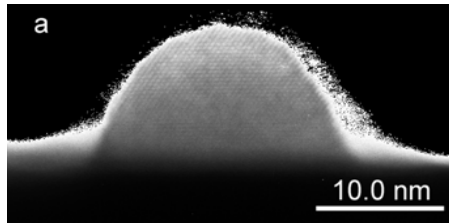




OPTICAL QUANTUM DOTS FOR QUANTUM INFORMATION



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outline

- single spin qubits
- two qubit gates
- entanglement in QD molecules

NRL theory: S. Economou, D. Solenov, T. Reinecke

NRL expt: D. Gammon, A. Bracker, S. Carter

U. Wuerzburg: fabrication, expt

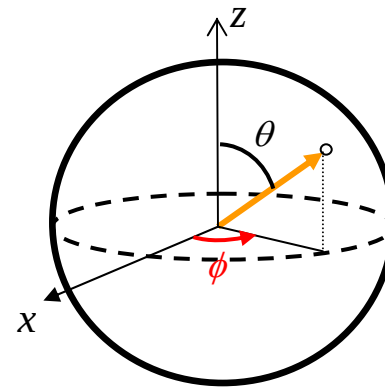
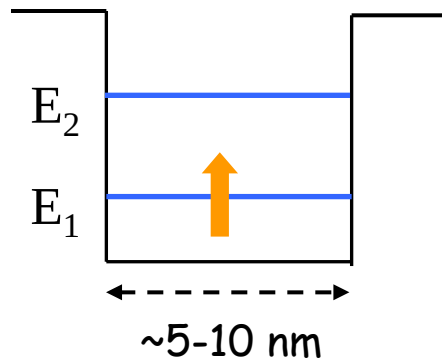
U. Dortmund: expt

support: DARPA, NSA and ONR



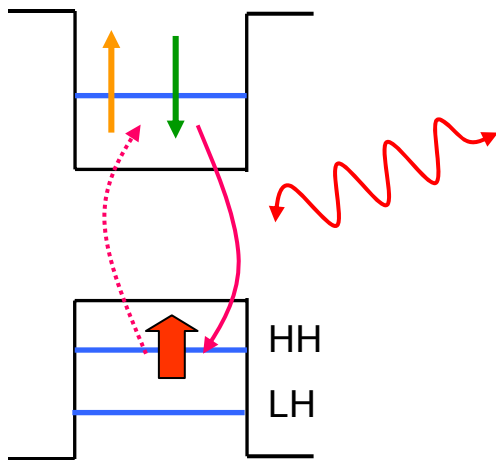
Optical quantum dots

spin in QD



$$\Psi = \alpha|\uparrow\rangle + \beta|\downarrow\rangle$$

optical properties, control (trion)



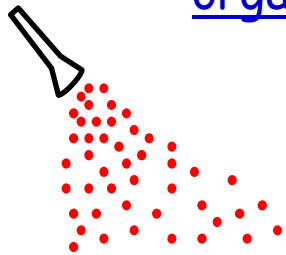
Quantum dots

- spin natural qubit
- 'long' coherence
- fast optical manipulations
- integrate into semiconductor technology
- potentially scalable
- fixed locations

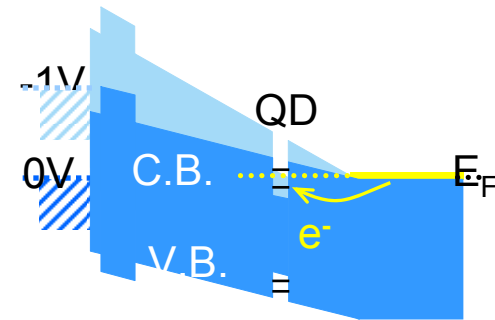
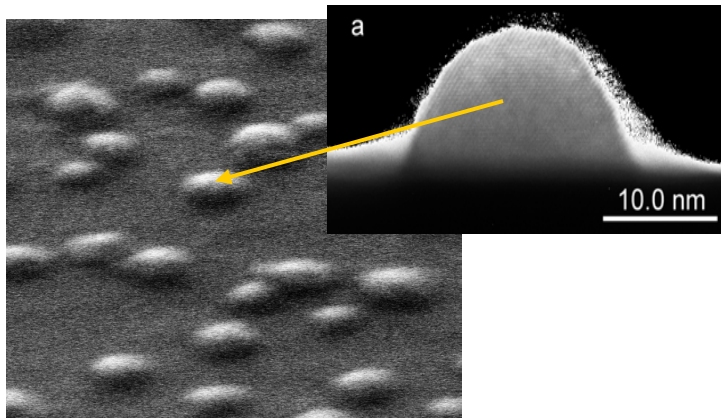
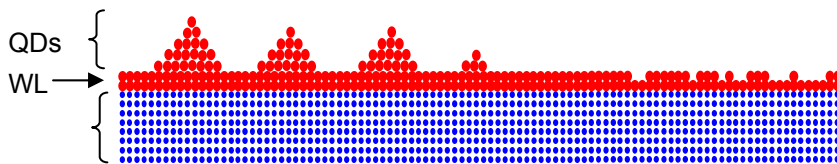


Quantum dots

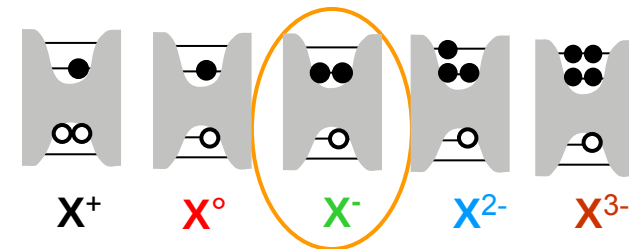
Stranski-Krastanov 'self organized' QDs



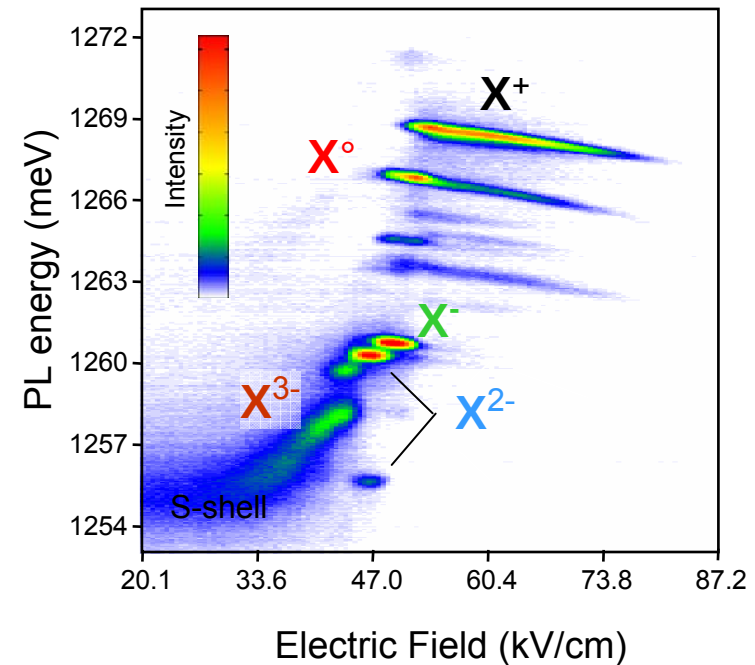
- substrate **GaAs**:
small lattice constant
- deposited Material **InAs**:
large lattice constant



spins in low states in QDs

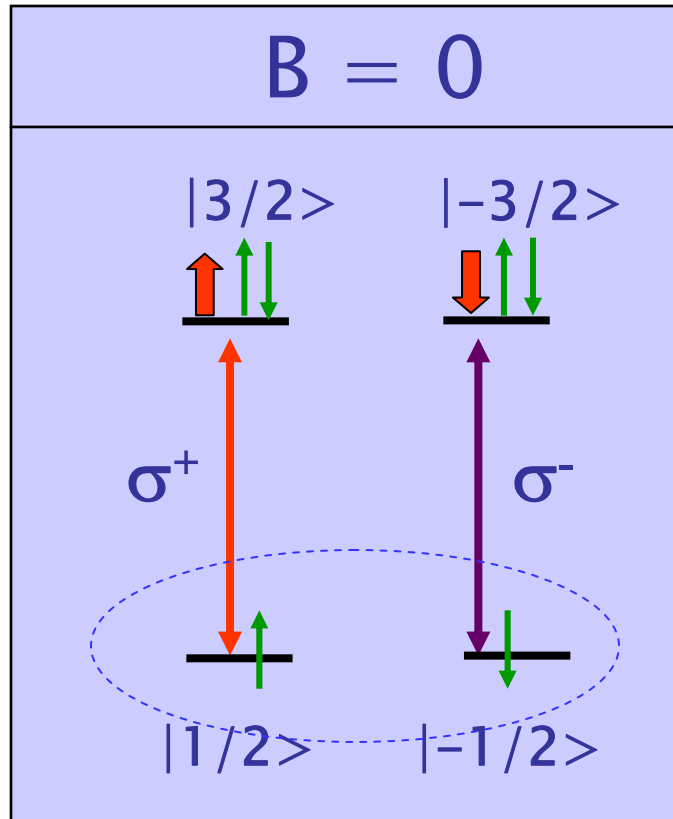


trion

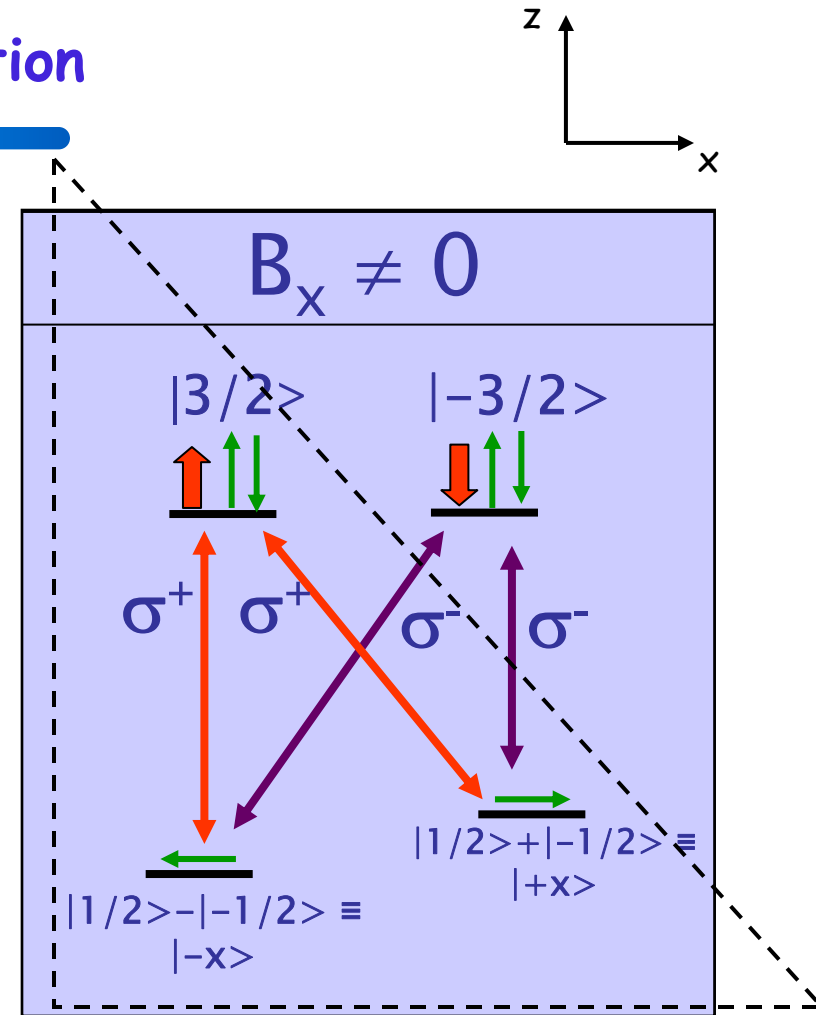




Spin initialization, manipulation



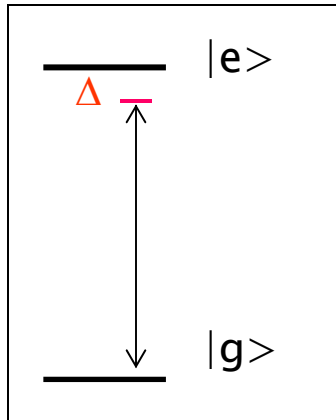
- No Raman transitions $|1/2\rangle \rightarrow |-1/2\rangle$
 - Cannot make arbitrary spin rotation
- ('z basis' - growth direction/optical axis)



- Voigt field mixes spin states & enables Raman transitions
 - optical pumping gives initialization
- ('x basis')



Useful feature of 2-level systems



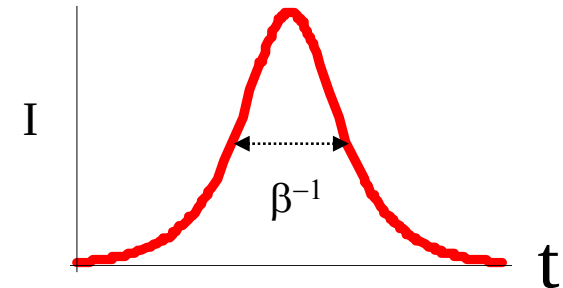
$$H(t) \sim \begin{pmatrix} 0 & V_{ge}(t) \\ V_{ge}(t)^* & (E_e - E_0) \end{pmatrix}$$

analytical solution

$$V_{ge} = \Omega \operatorname{sech}(\beta t) e^{i\Delta t}$$

Δ = detuning

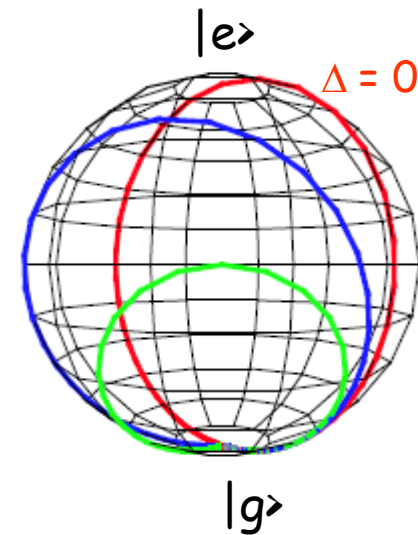
β = pulse bandwidth



properties:

- (i) $\Omega = \sigma \rightarrow 2\pi$, no excitation left in $|e\rangle$
- (ii) pulse area indep of detuning
- (iii) $|g\rangle$ acquires phase

$$\tan(\phi) = \frac{2\beta\Delta}{\Delta^2 - \beta^2}$$



(Rosen & Zener, Phys. Rev., 1932)

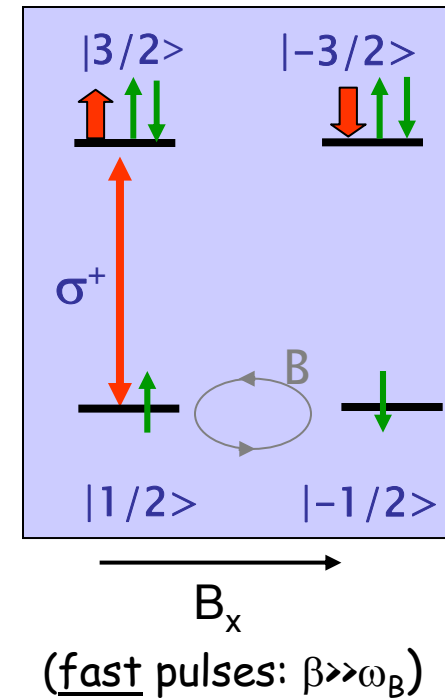


Spin rotations around one axis (z)

$$H(t) \sim \begin{pmatrix} 0 & \omega_B & 0 \\ \omega_B & 0 & \Omega(t)e^{i\omega_0 t} \\ 0 & \Omega^*(t)e^{-i\omega_0 t} & \varepsilon_T \end{pmatrix}$$

'z-basis' $\{|\bar{z}\rangle, |z\rangle, T\}$

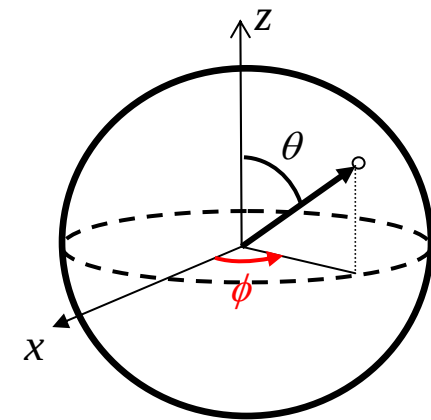
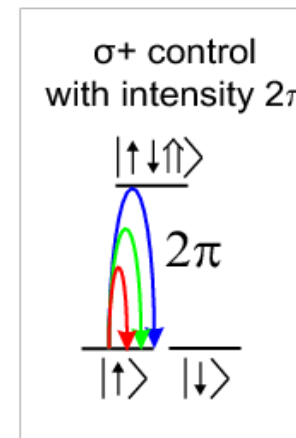
Pulses: $V_{ge} = \Omega \sec h(\beta t) e^{i\Delta t}$



Properties:

- (i) 2π pulses return to spin subspace with no trion excitation
- (ii) 2π pulses give phase ϕ of $|z\rangle$ wrt $|z'\rangle$ equivalent to a z rotation

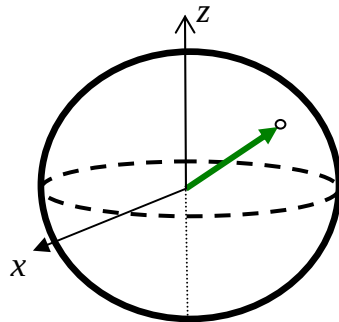
$$\phi = 2 \arctan(\beta / \Delta)$$



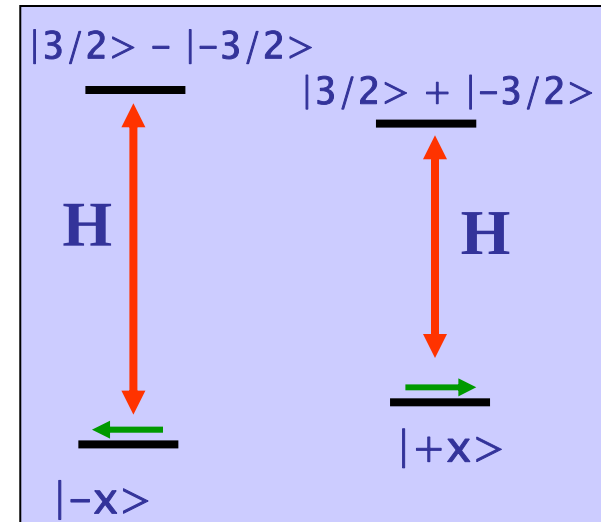


Rotations around another axis (x)

- single, broad band, linearly polarized pulse couples both transitions
- different phase $|x\rangle$ and $|x'\rangle$
- pulse area independent of detuning
- gives rotation $(\phi_1 - \phi_2)$



general rotations from
those about z and x



π_x polarizations

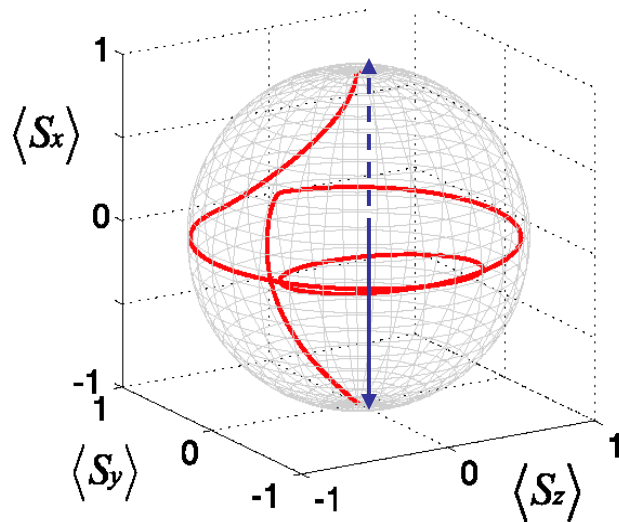


general rotations, full solutions

- density matrix calculations including mixing, losses
- arbitrary rotation from rotations around two axes

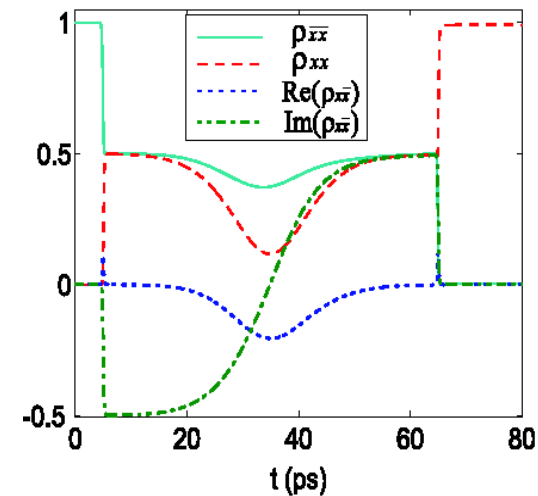
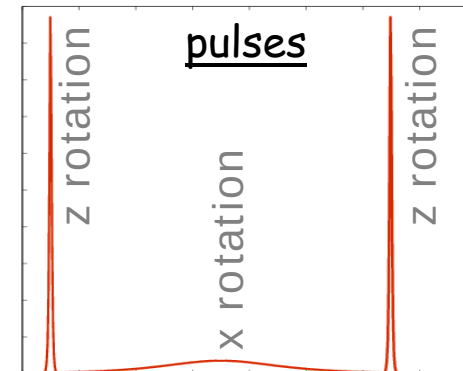
$$R_y(\phi) = R_z^+(\pi/2) R_x(\phi) R_z(\pi/2)$$

Fidelity 99.28%



parameters for InAs QDs

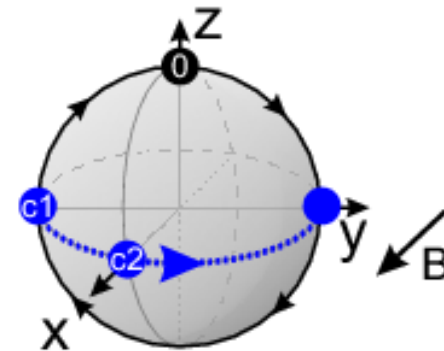
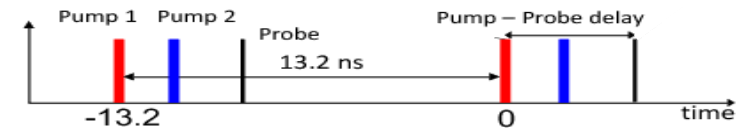
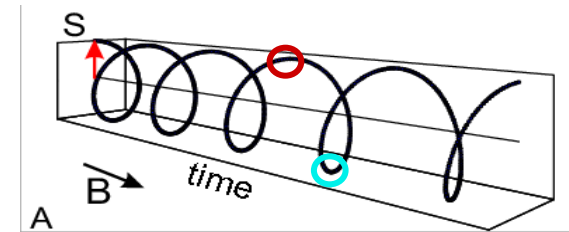
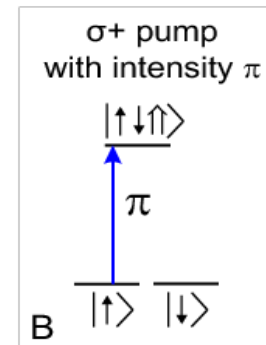
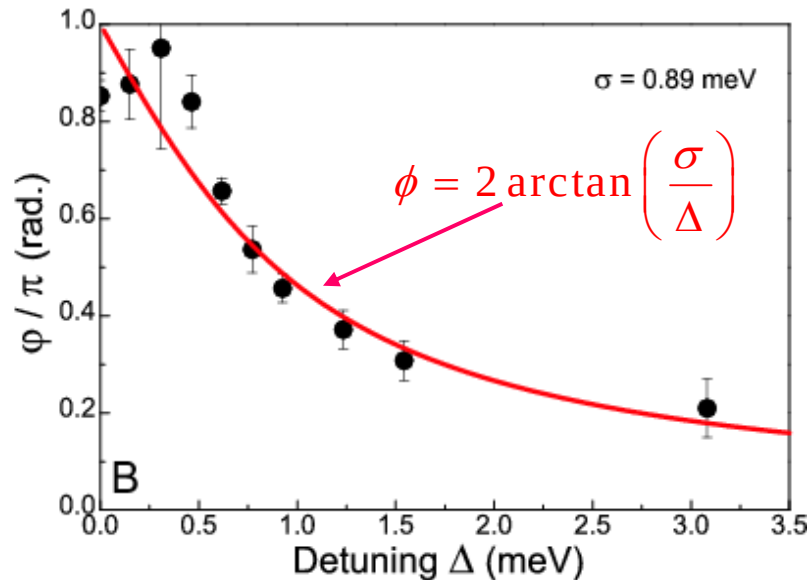
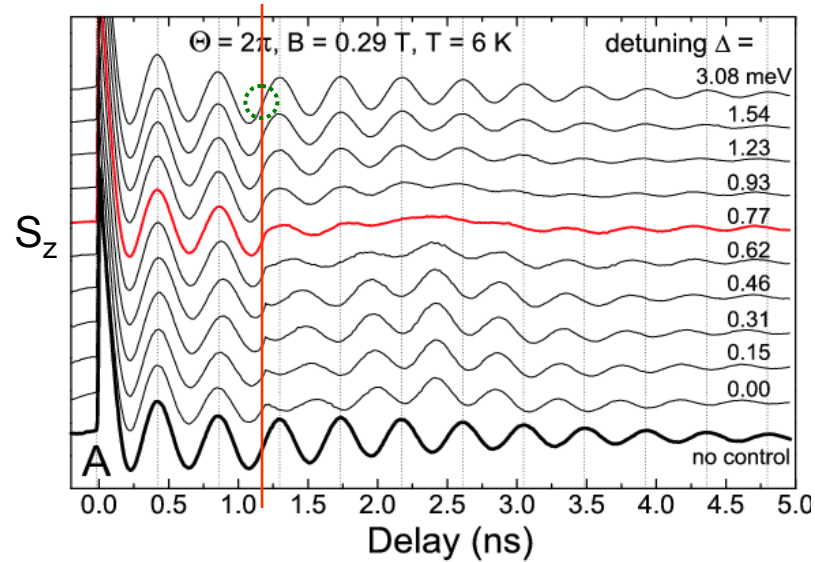
Economou & Reinecke, PRL (2007)



fast gates ~ 60 psec



Spin rotations - ϕ vs Δ



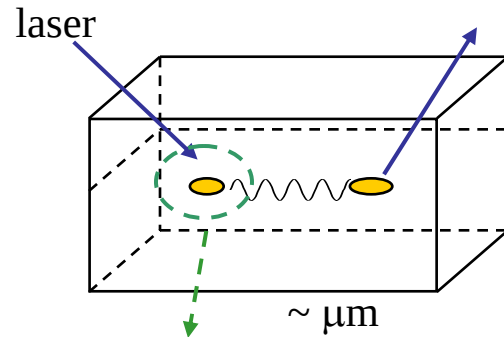
general rotations demonstrated
by varying precession

Greilich *et al*, Nature Physics 5, 262 (2009)



Cavity coupling

fast, distant, frequency selective
coupling and architectures

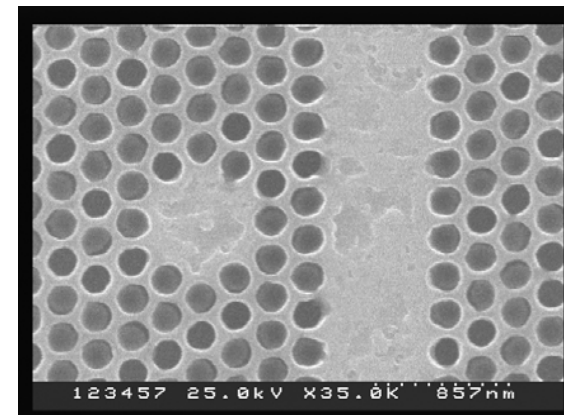
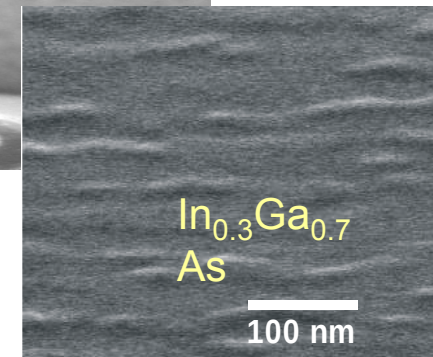
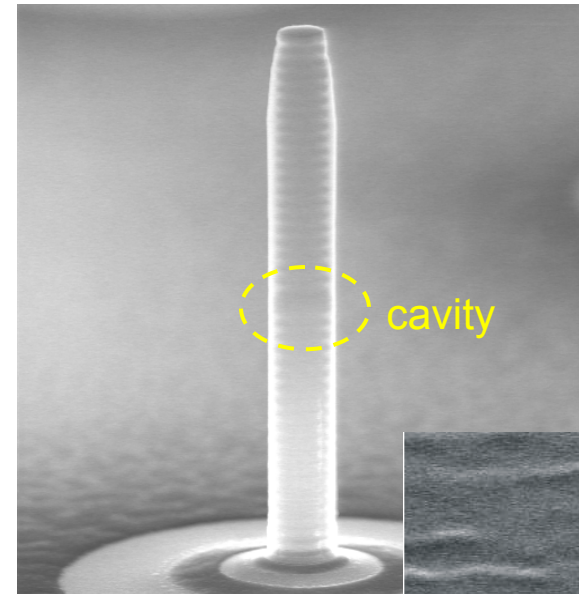


strong coupling (coherent)
weak coupling (Purcell)

$$\omega_{\pm} = \omega_c \pm \sqrt{g_{e-p}^2 - (\gamma_c + \gamma_e)^2 / 4}$$

$$g = \sqrt{\frac{e^2 f}{4\pi\kappa\epsilon_0 V_m}}$$

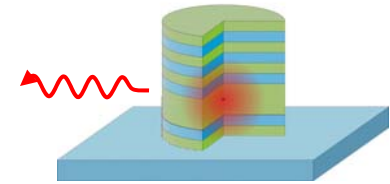
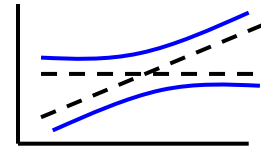
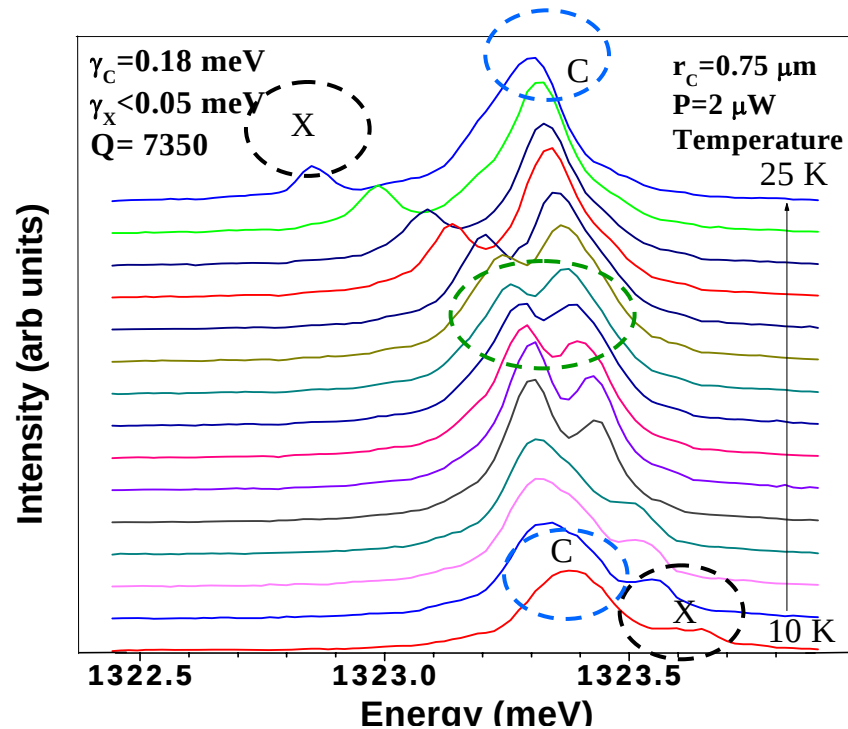
Cavity development



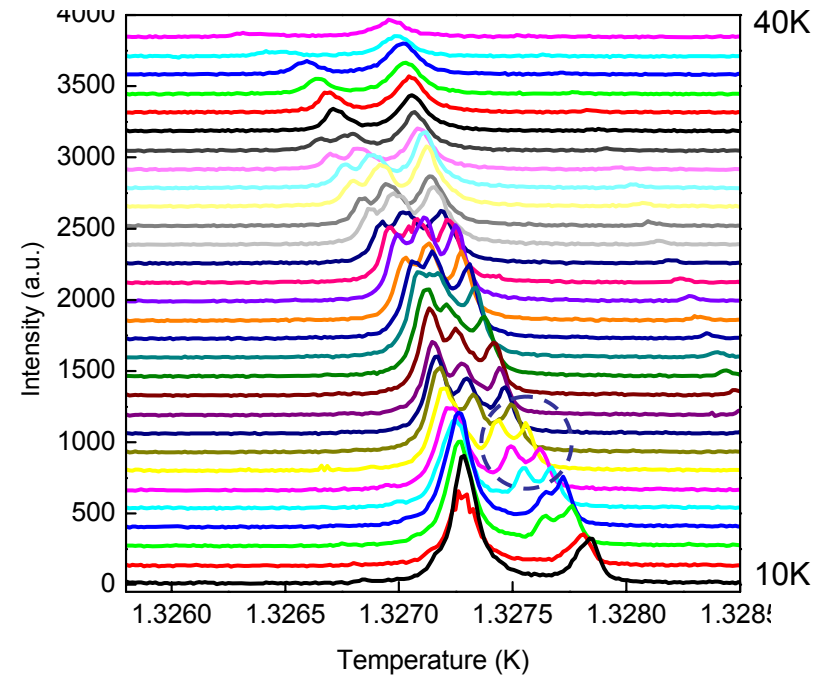


Strong coupling

one dot spectra



one dot spectra

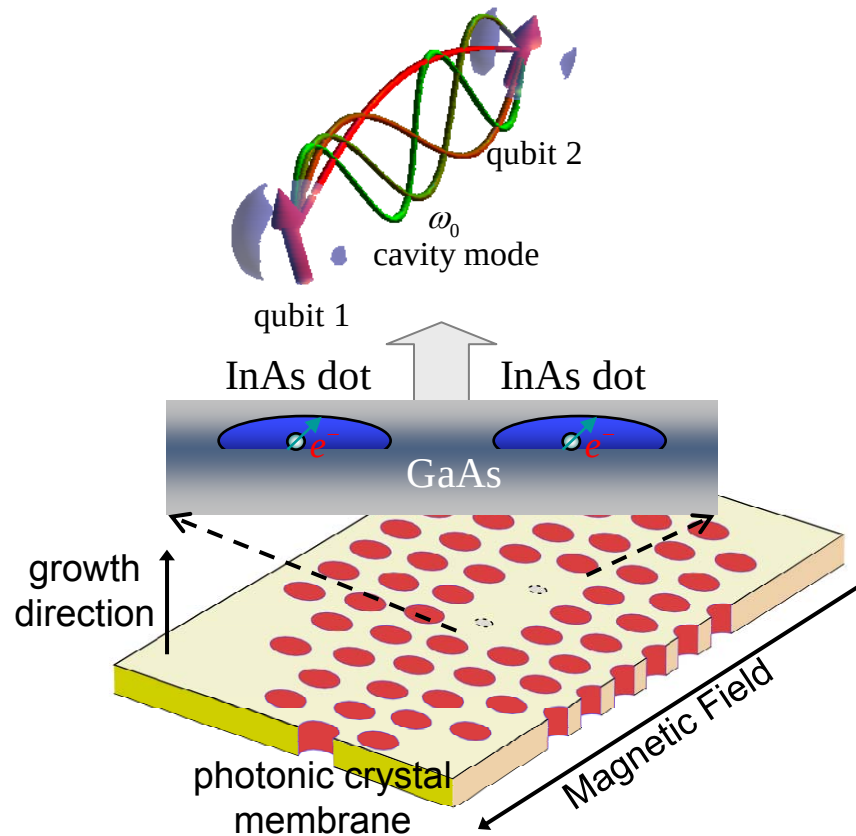


- $g_{\text{ex-p}} \sim 100\text{-}200 \text{ meV}$
- 40 pillars from several wafers with $g_{\text{ex-p}} > 50 \text{ meV}$
- confirmed by intensity and linewidths

Reithmaier et al, Nature **432**, 197 (2005)

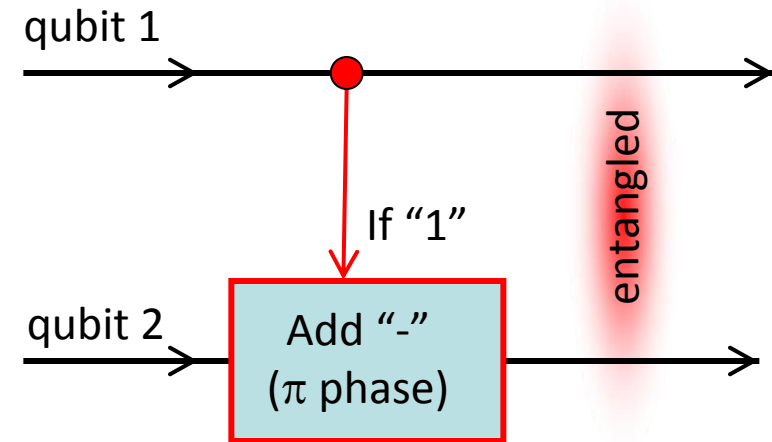


Two qubit entangling gate (theory)



Requirements for two qubit (entangling) gate:

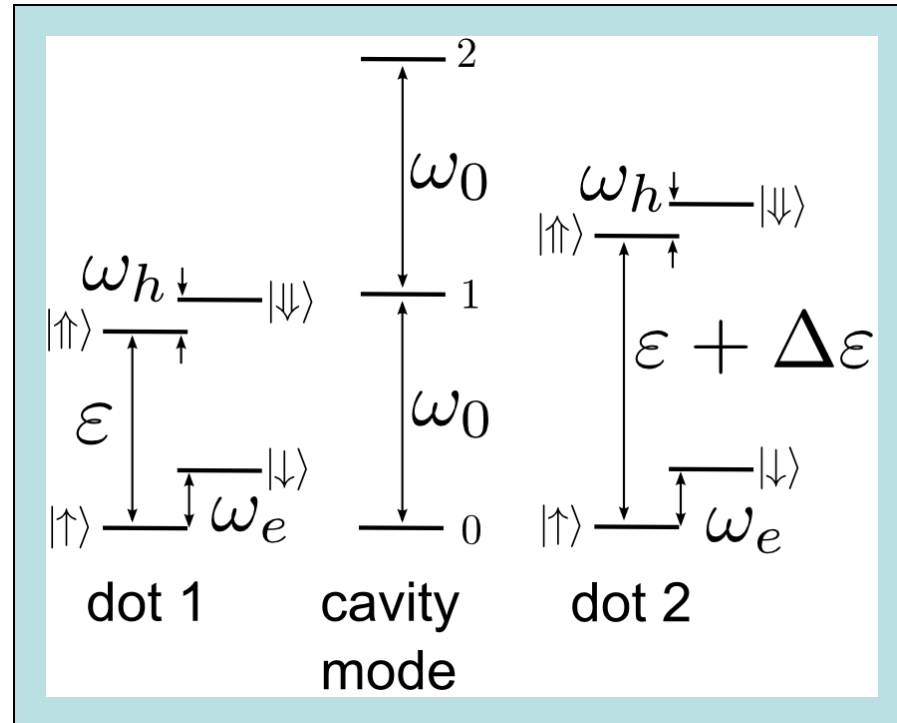
- (Long-range) physical interaction
- Dynamical control of interaction
- Conditionality (control-U operation)
- Scalability



compatible with one qubit operations



system



System ($n=1,2$)

$$H_{e-t,n} = -\omega_e c_{n\uparrow}^\dagger c_{n\uparrow} + \omega_e c_{n\downarrow}^\dagger c_{n\downarrow} + (\varepsilon_n - \omega_h) t_{n\uparrow}^\dagger t_{n\uparrow} + (\varepsilon_n + \omega_h) t_{n\downarrow}^\dagger t_{n\downarrow}$$

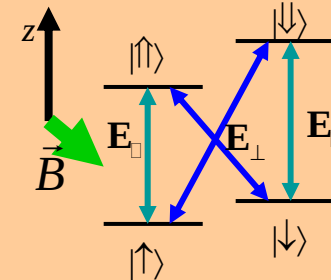
$$H_c = \omega_0 a^\dagger a$$

$$H_{e-c,n} = g(t_{n\uparrow}^\dagger c_{n\uparrow} + t_{n\downarrow}^\dagger c_{n\downarrow} + h.c.)(a^\dagger + a)$$

Pulses

$$V(t) = \sum_{\substack{p \\ \xi\xi'=\uparrow,\downarrow \\ n=1,2}} \Omega_{p(\xi\xi')}(t-t_p) \cos \omega_p t (t_{n\xi}^\dagger c_{n\xi'} + h.c.)$$

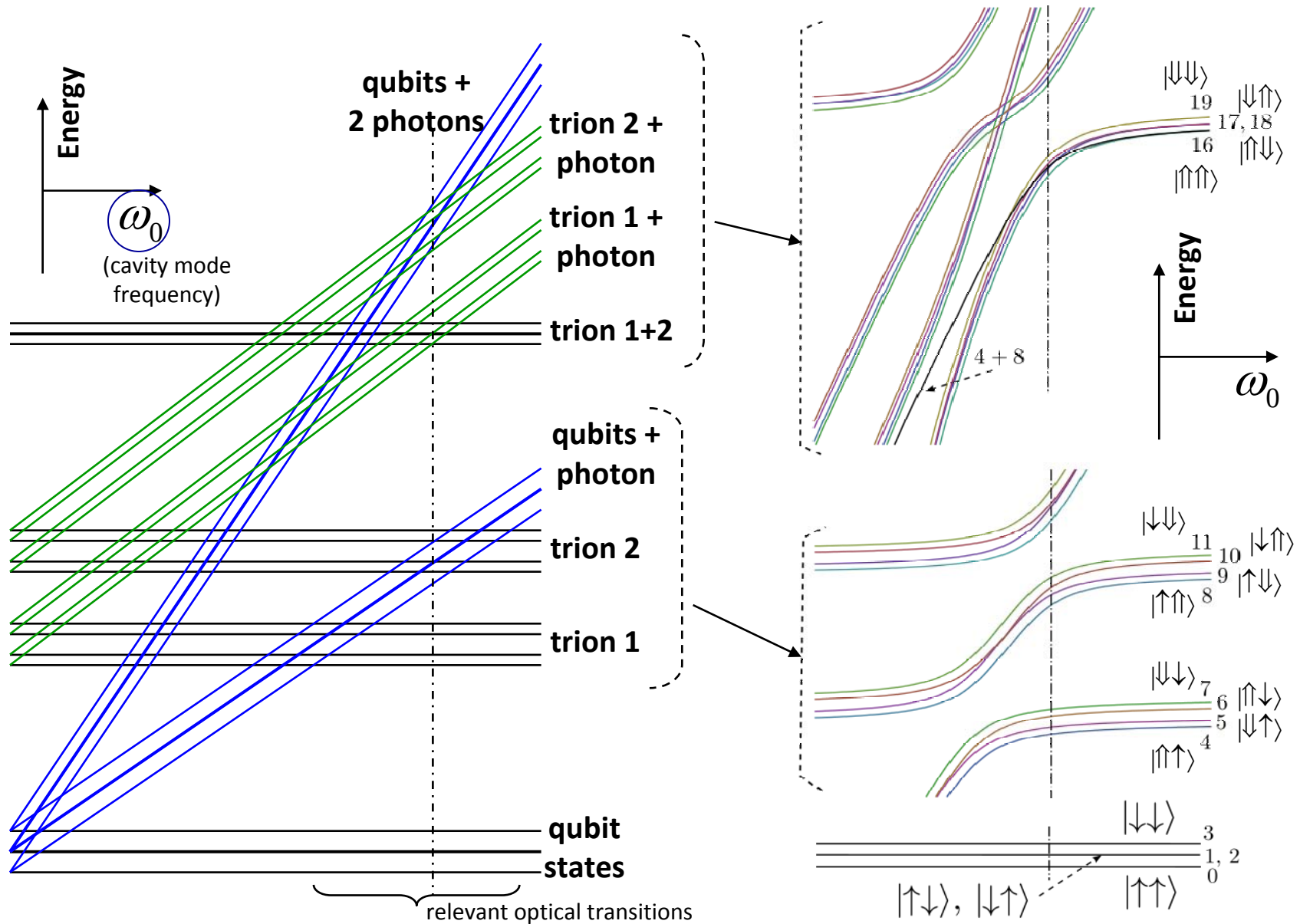
Magnetic field perpendicular to the growth axis (z)



$$H_{qubits} = \frac{\omega_e}{2} \sigma_z^{(1)} + \frac{\omega_e}{2} \sigma_z^{(2)}$$



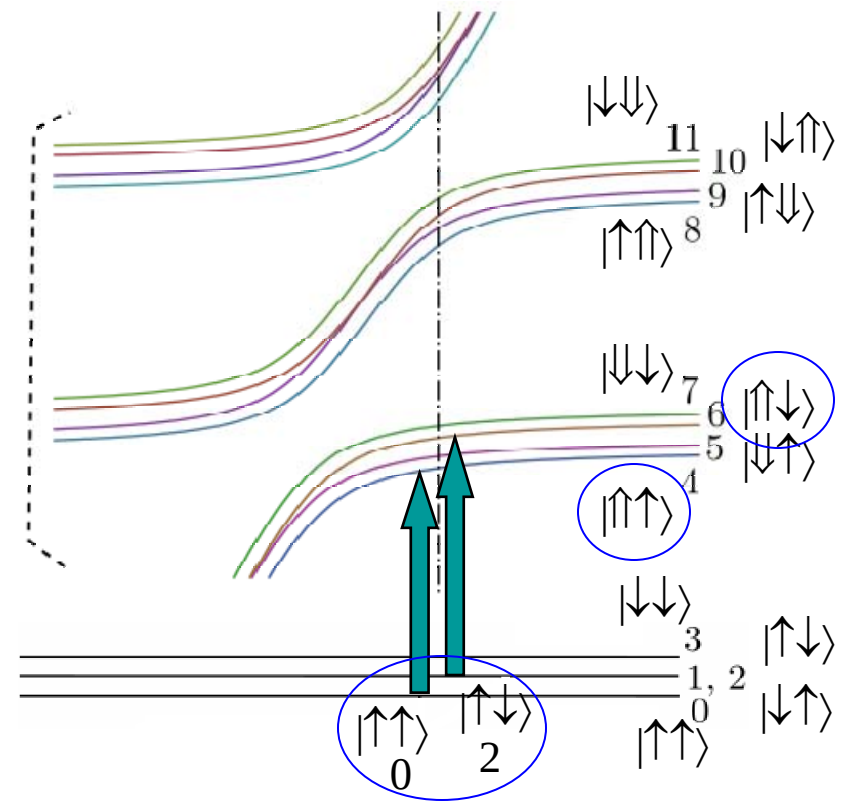
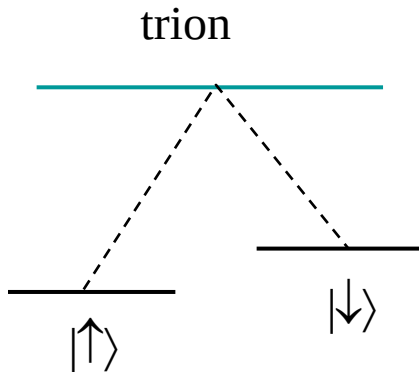
spectra





Single qubit operations: compatibility

- effects of cavity-induced splitting in one-particle sector negligible for short pulses
- single qubit operations same as previously
- need single and two-qubit operations for complete set of operations

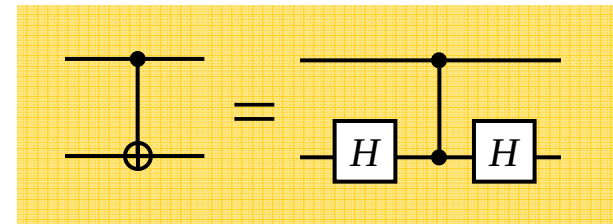
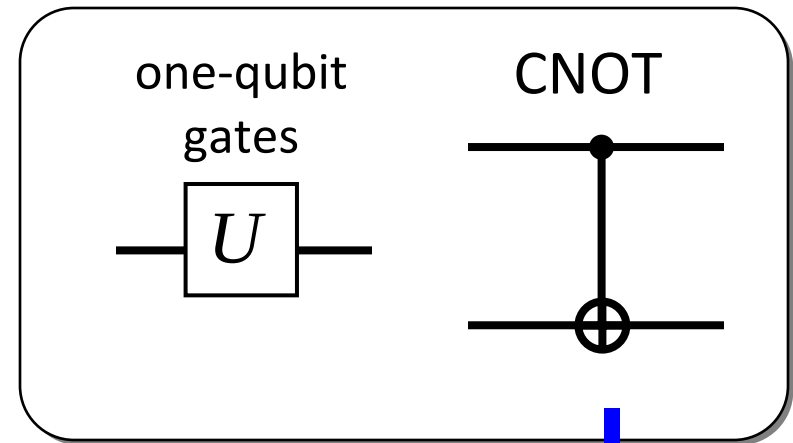


fast pulses

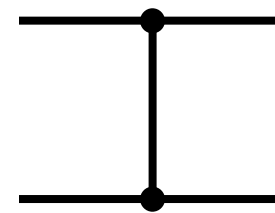


Two-qubit gates: CZ gate

universal



CZ



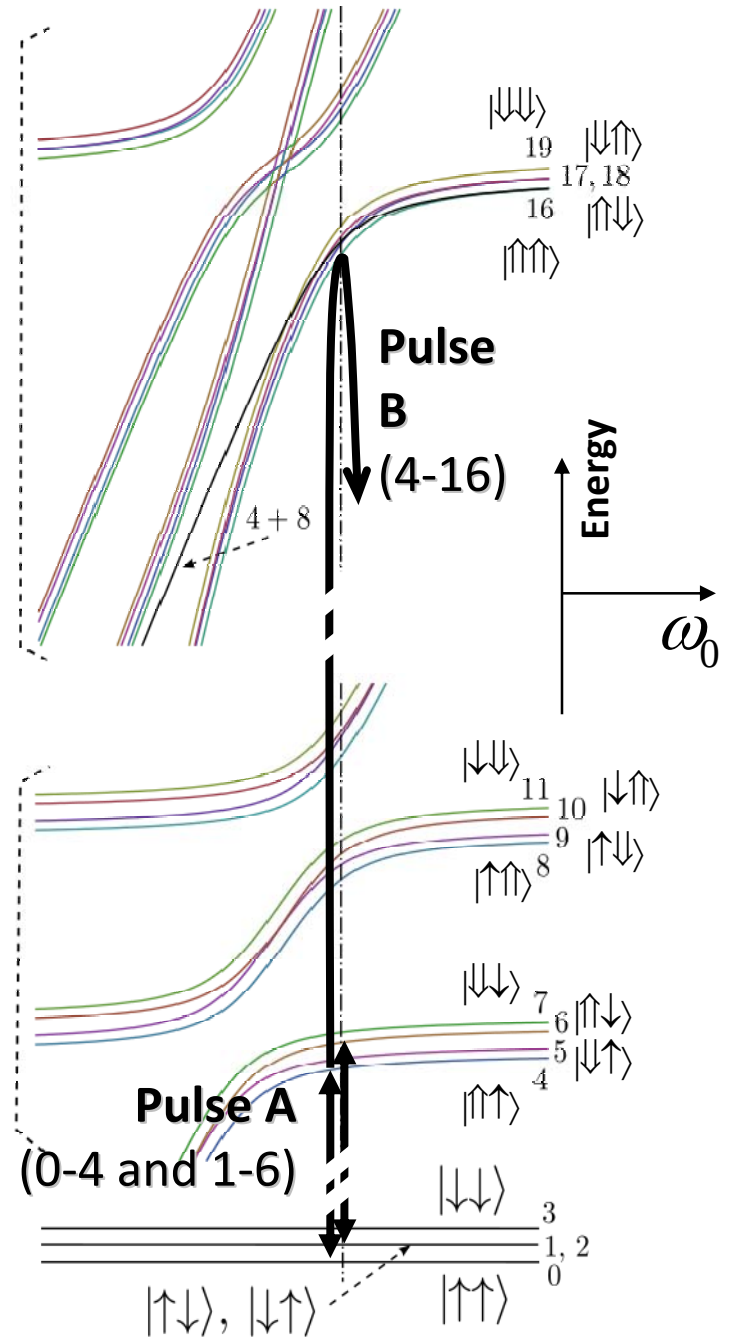
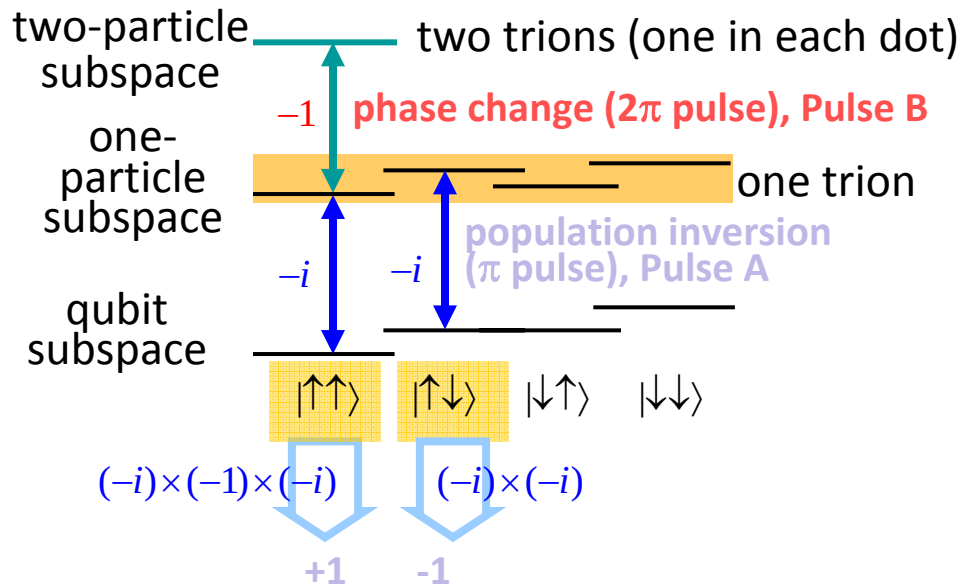
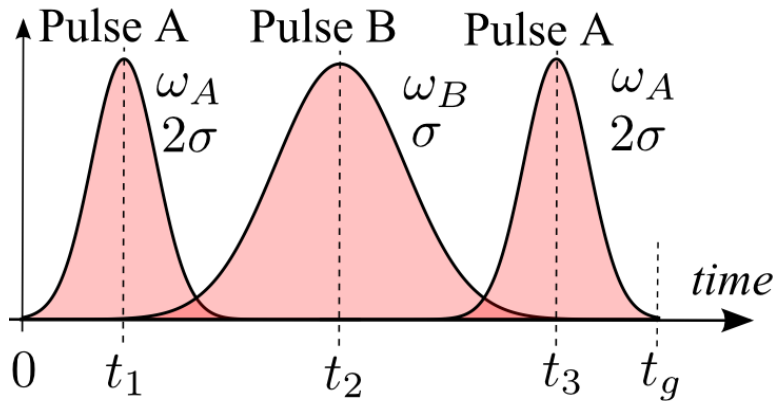
$$|\psi\rangle = \alpha |\uparrow\uparrow\rangle + \beta |\uparrow\downarrow\rangle + \delta |\downarrow\uparrow\rangle + \gamma |\downarrow\downarrow\rangle$$

$$|\psi\rangle = \alpha |\uparrow\uparrow\rangle - \beta |\uparrow\downarrow\rangle + \delta |\downarrow\uparrow\rangle + \gamma |\downarrow\downarrow\rangle$$

$$U_g = \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & -1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \end{pmatrix}$$



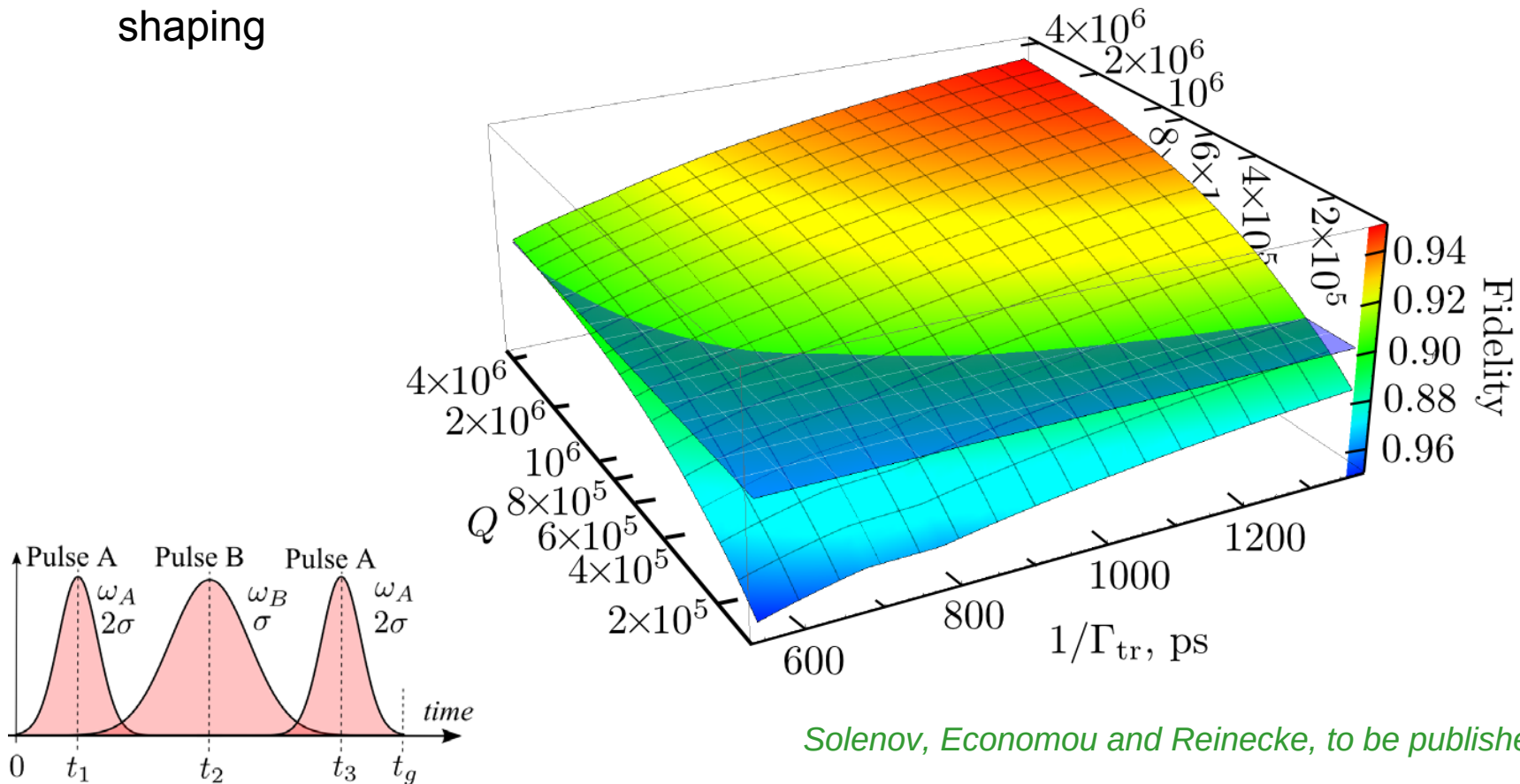
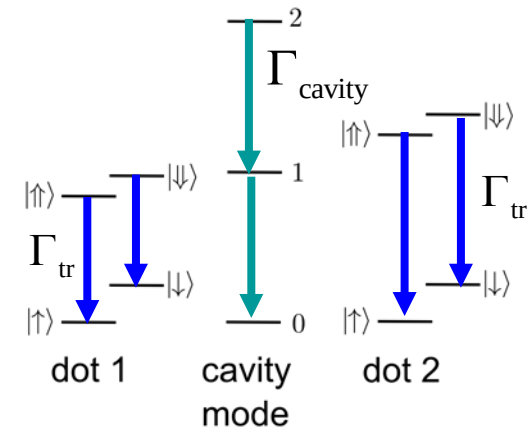
CZ gate





Fidelity

1. losses due to unwanted coherent dynamics of off-resonant transitions
 2. losses due to carrier recombination and cavity losses
- can be improved through e.g., pulse shaping

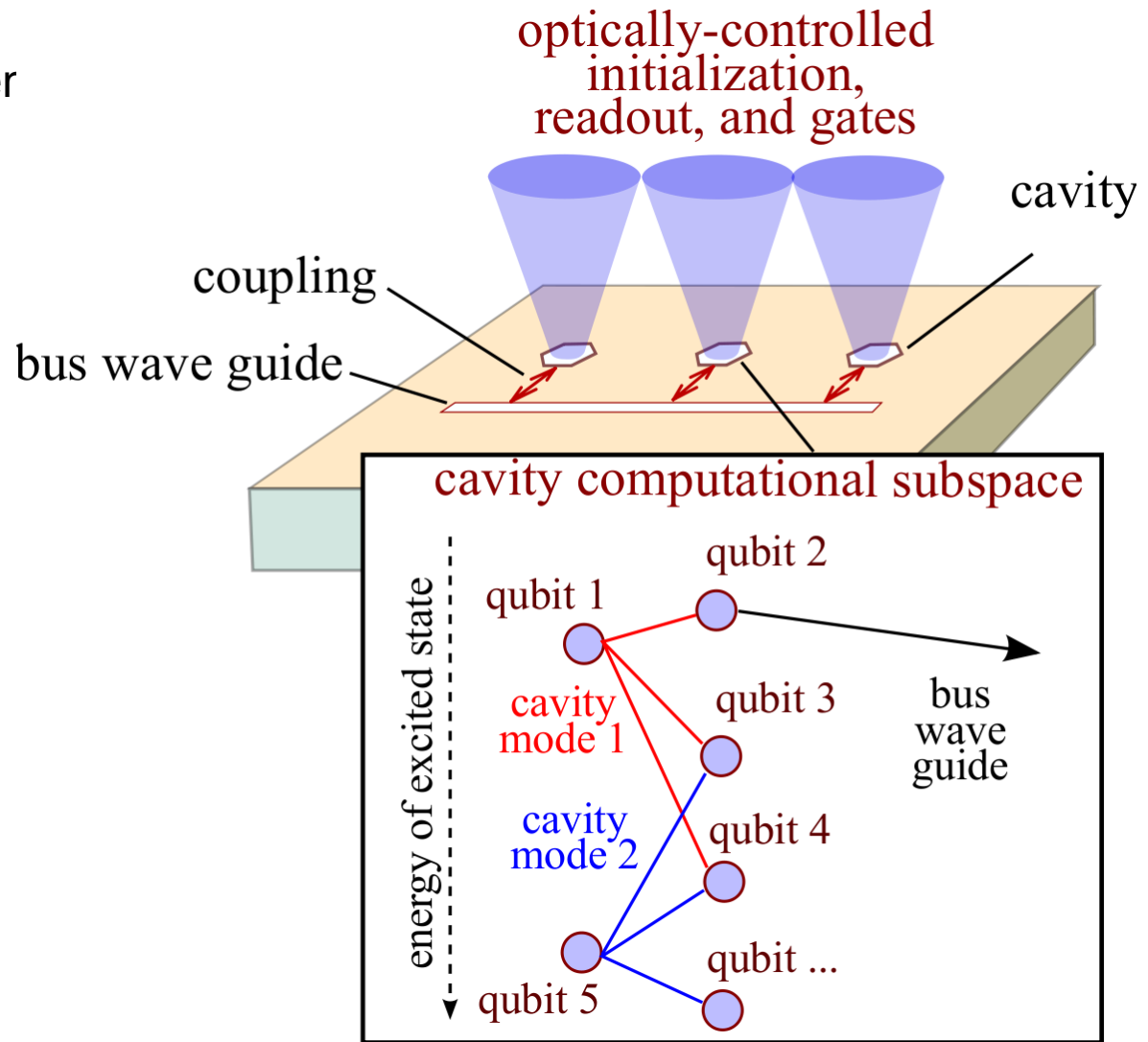


Solenov, Economou and Reinecke, to be published



Scalability

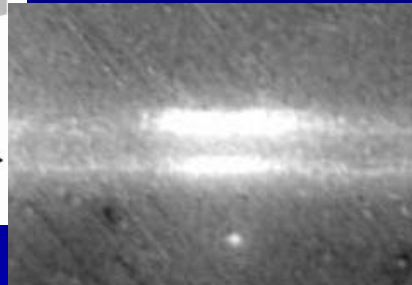
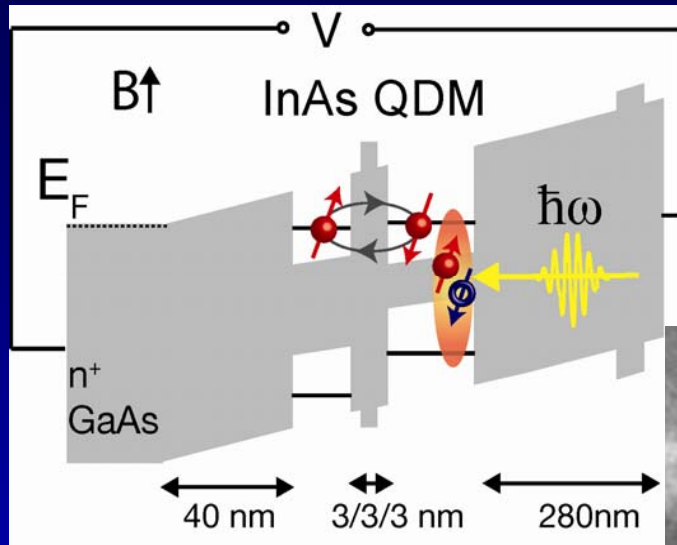
- gates for unequal dot excitation energies allow coupling between multiple dots
- multiple confined photon exist modes in cavities



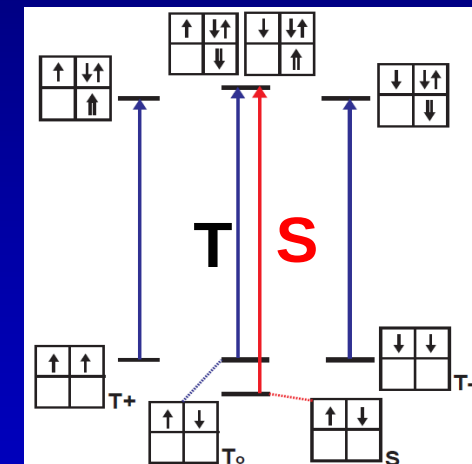


2 electron spins in two dots

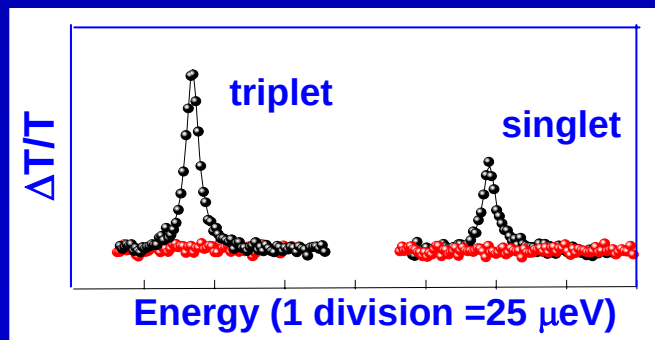
- Tunnel coupled QDs
- Exchange interaction always on
- Can access optically single qubit and two-qubit states



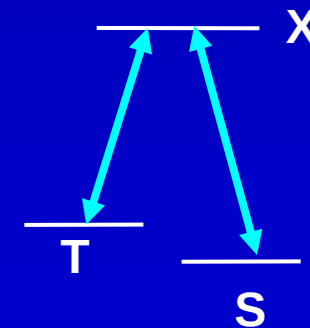
Energy diagram



Initialization



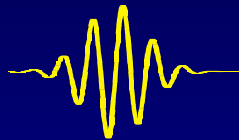
kinetic exchange Δ_{ee}





Short vs. Long optical pulses

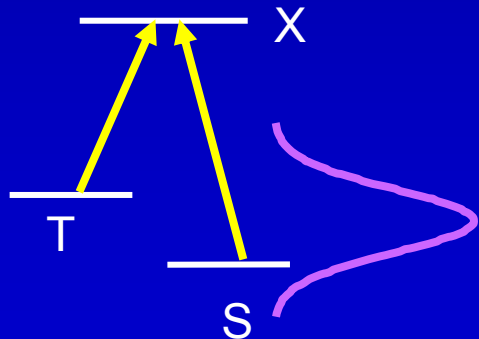
Short pulses



Time domain:

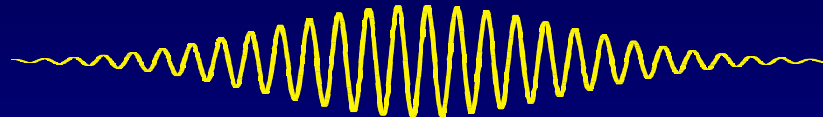
- Acts on *single* spin state because *faster* than exchange interaction

Frequency domain:



Acts on S + T

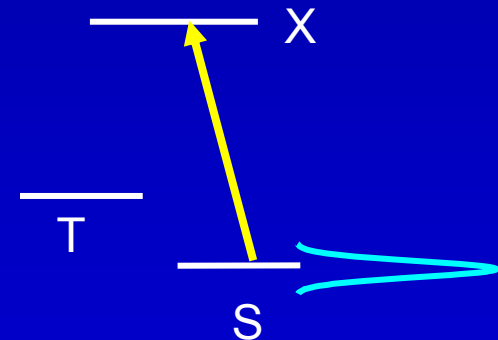
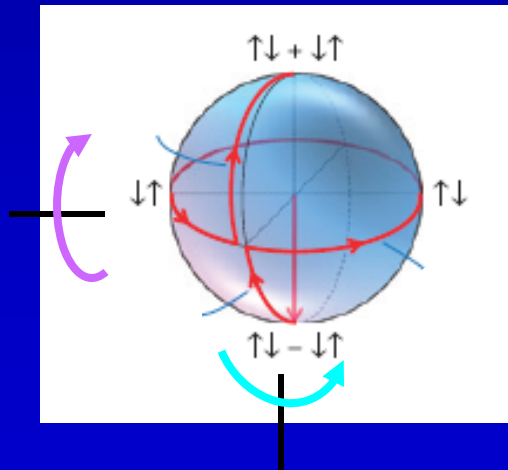
Long pulses



Time domain:

- Acts on *joint* spin state because *slower* than exchange interaction

Frequency domain:



Acts just on S



Optical 2-qubit phase gate

$$\begin{aligned} |\uparrow\downarrow\rangle - |\downarrow\uparrow\rangle &\Rightarrow e^{i\phi} (|\uparrow\downarrow\rangle - |\downarrow\uparrow\rangle) \\ |\uparrow\downarrow\rangle &\Rightarrow |\uparrow\downarrow\rangle - i|\downarrow\uparrow\rangle, \text{ for } \phi = \pi/2 \end{aligned}$$

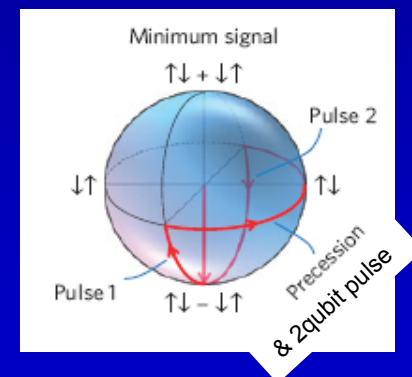
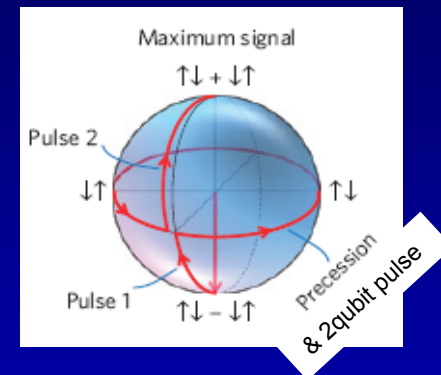
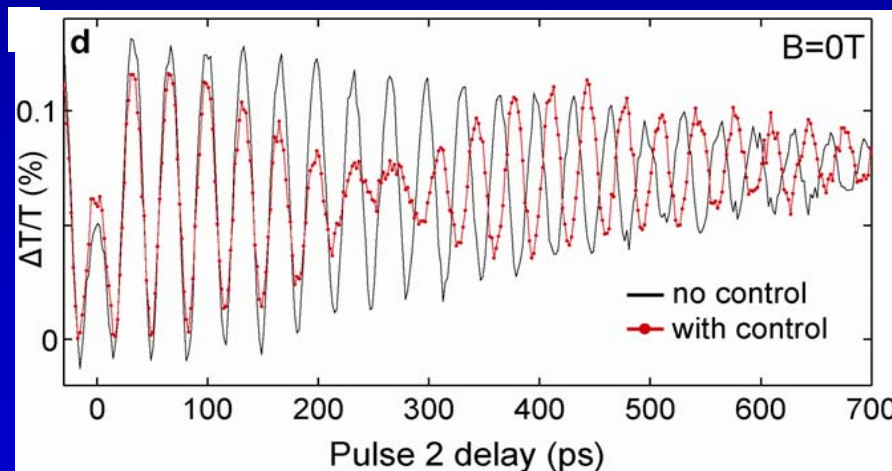


$\sqrt{\text{SWAP}}$ gate

Pulse sequence



Ramsey fringes with control pulse



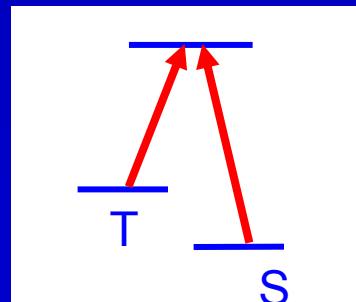
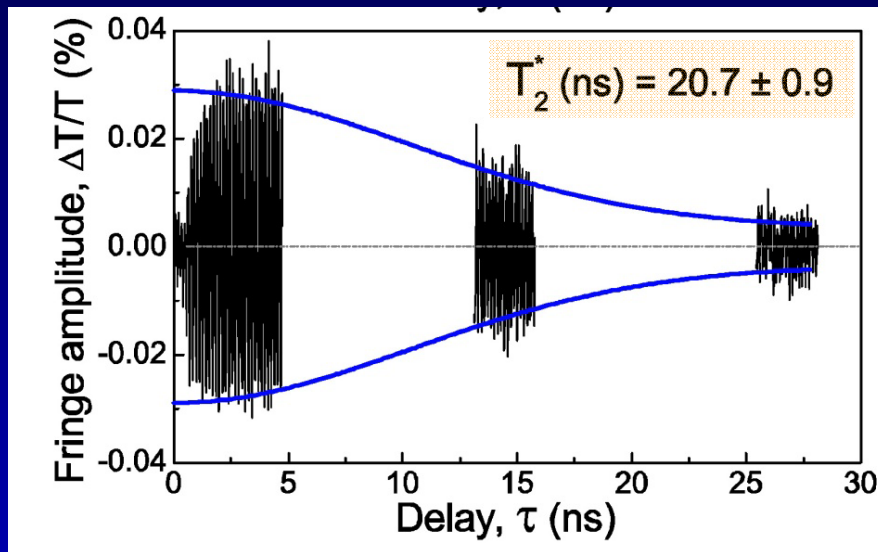
Can vary phase change from -180° to $+180^\circ$ with pulse detuning

Kim et al, Nat. Physics, 7, 223 (2011)

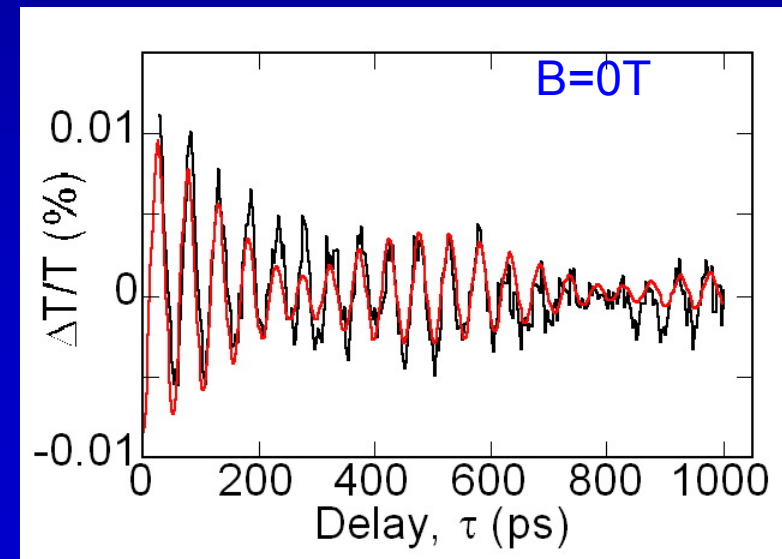


Alternative qubit: hole spin

Ramsey fringes of single hole qubit



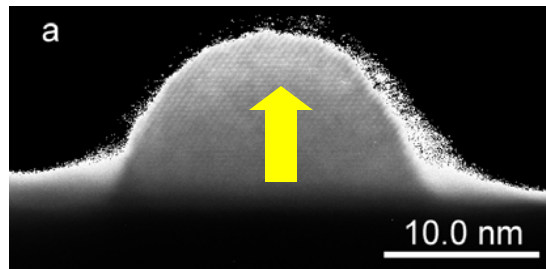
Coupled hole qubits



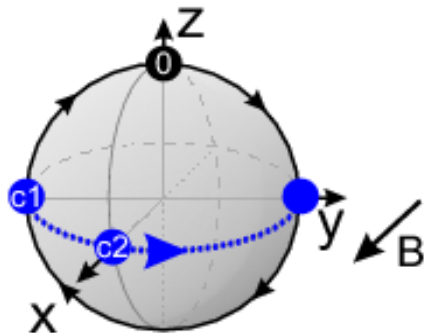


SUMMARY

- Spins in quantum dots



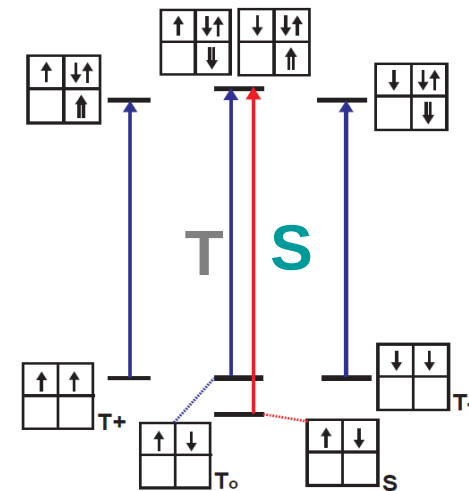
- single qubit operations



- Two qubit phase gate

$$U_g = \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & -1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \end{pmatrix}$$

- entanglement in QD molecules





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