

Part II: Picturing Even More Quantum Processes

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Spring School on Quantum Structures in Physics and CS

May 29, 2014

1. Review **quantum maps**, **quantum/classical maps**, and **spiders**


Outline

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2. Enrich our language with **multi-coloured spiders** and **phases**


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
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4. Specialise to qubits and define the **ZX-calculus**

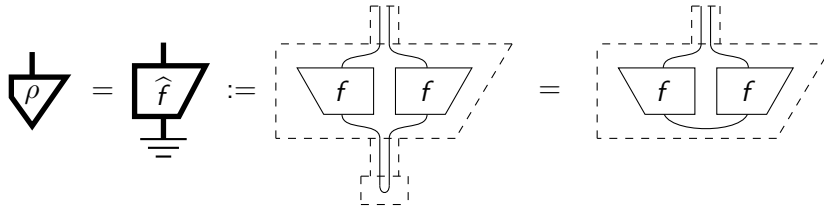
Review – Quantum states

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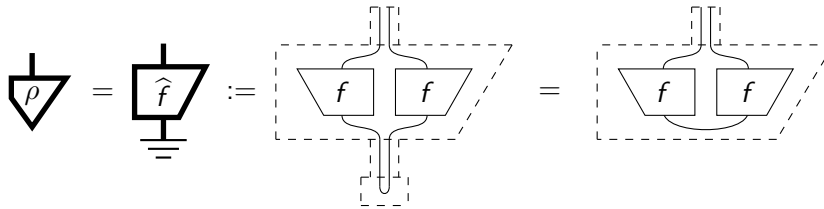
► They can always be written in terms of a **pure state** +  :



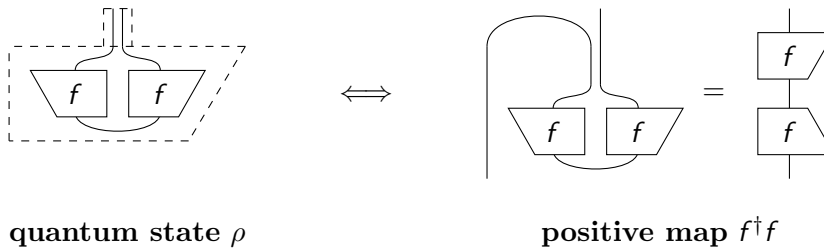
Review – Quantum states

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► So ‘up to bending’, a.k.a. partial transpose:



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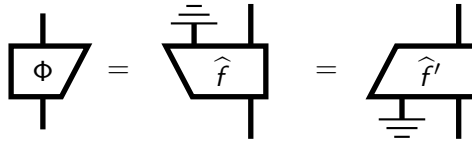


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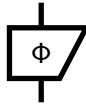


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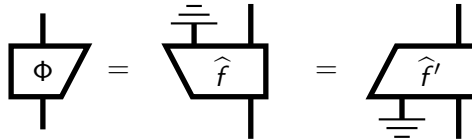


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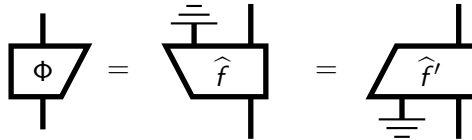


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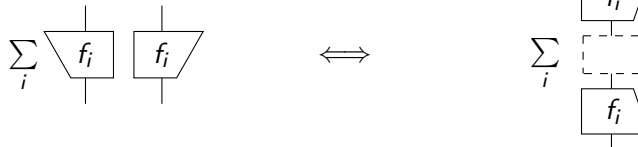
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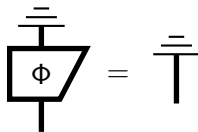


quantum map Φ

CP-map $\sum_i f_i(-) f_i^\dagger$

Review – Discarding and causality

- **Physically realisable** quantum maps satisfy **causality**:



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$$\text{Box } \Phi \text{ with top and bottom wires ending in } \text{---} = \text{---}$$

- **Discarding** a state amounts to taking a **trace**:

$$\text{Triangle } \rho \text{ with top wire ending in } \text{---} = \text{Dashed box with two } f \text{ trapezoids in a loop} = \text{Two } f \text{ trapezoids in a loop} = \text{Tr}(\rho)$$

Review – Discarding and causality

- **Physically realisable** quantum maps satisfy **causality**:

$$\text{Discard} \circ \Phi = \text{Discard}$$

- **Discarding** a state amounts to taking a **trace**:

$$\text{Discard}(\rho) = \text{Tr}(\rho)$$

- **Causal states** \leftrightarrow **positive operators with trace 1**
Causal maps \leftrightarrow **trace-preserving CP-maps (CPTPs)**

Review – Classical states

- **Classical states** look like this:

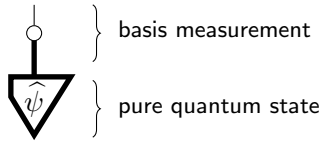


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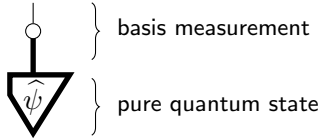


Review – Classical states

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- ...hence the notation. The dot singles out a **preferred basis**, and in that basis, a classical state is a **vector of positive numbers**:

$$\begin{array}{c} \text{dot} \\ | \\ \text{triangle with } \hat{\psi} \end{array} = \sum_i p_i \begin{array}{c} | \\ \text{triangle with } i \end{array} \leftrightarrow \begin{pmatrix} p_1 \\ p_2 \\ \dots \\ p_n \end{pmatrix}$$

Review – Classical states

- ▶ **Classical states** look like this:



- ▶ They can always be written as:



} basis measurement

} pure quantum state

- ▶ ...hence the notation. The dot singles out a **preferred basis**, and in that basis, a classical state is a **vector of positive numbers**:

$$\text{Diagram with dot} = \sum_i p_i \text{Diagram with } i = \leftrightarrow \begin{pmatrix} p_1 \\ p_2 \\ \dots \\ p_n \end{pmatrix}$$

- ▶ Causality forces these numbers to sum to 1:

$$\text{Diagram with dot} = \text{Diagram with double bar} = \text{Empty box} \iff \sum_i p_i = 1$$

Review – Quantum/classical maps

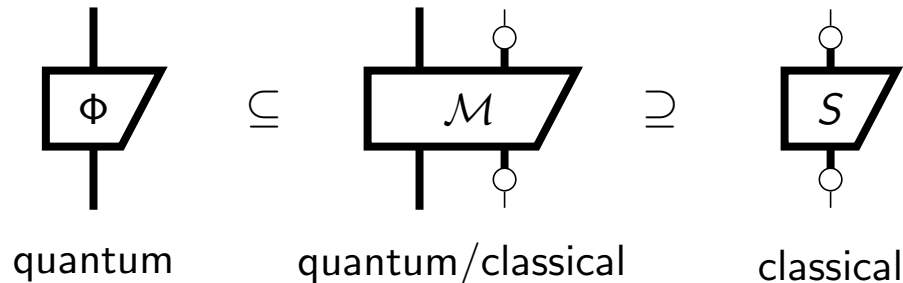
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- ▶ So, **causal classical states** are just plain old **probability distributions**.
- ▶ Similarly, **causal classical maps** are precisely the linear maps that preserve probability distributions, a.k.a. **stochastic maps**.
- ▶ **Quantum/classical maps** generalise both **CP-maps** and **stochastic maps**.



Review – Spiders

- ▶ Linear/quantum maps can be defined in terms of **basis states** (and numbers) using **sums**.

Review – Spiders

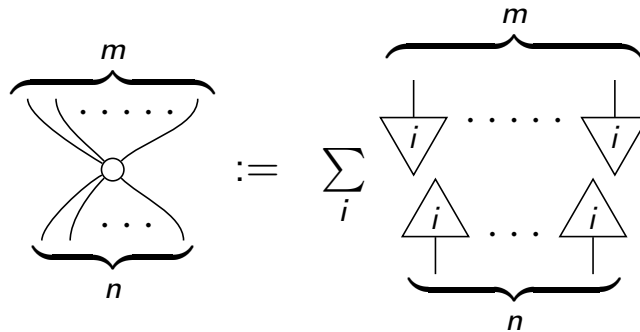
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Spiders!

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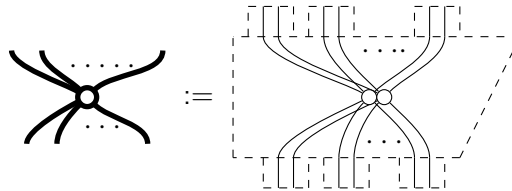


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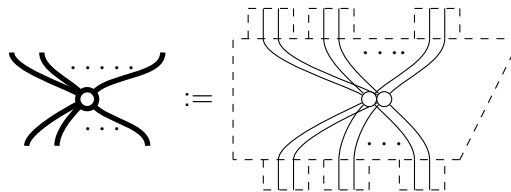


Review – Spiders

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- ▶ We have seen classical spiders (single wires):



- ▶ ...quantum spiders (double wires):



- ▶ ...and classical/quantum (a.k.a. bastard) spiders:

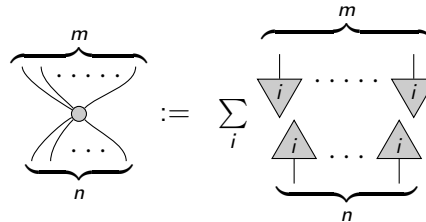
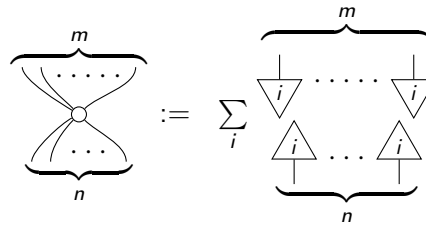


Multi-coloured spiders

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Multi-coloured spiders

- ▶ Most interesting quantum features appear only when we ditch **preferred bases** for systems and instead study **interaction of multiple bases**.
- ▶ Different bases \rightarrow different coloured spiders



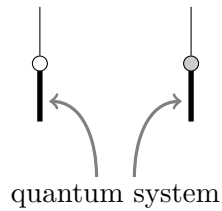
Two kinds of measurement

- Each spider induces a basis **measurement**:



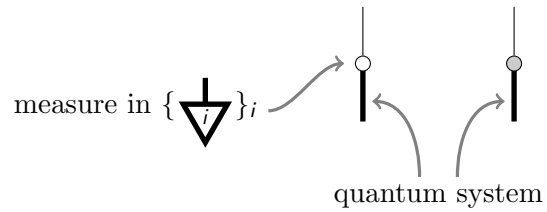
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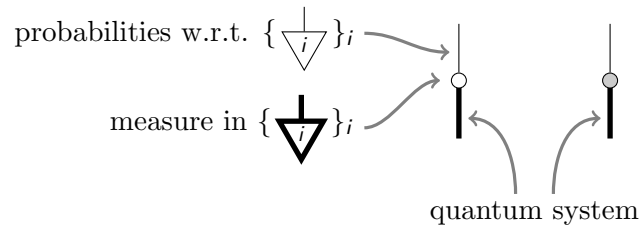
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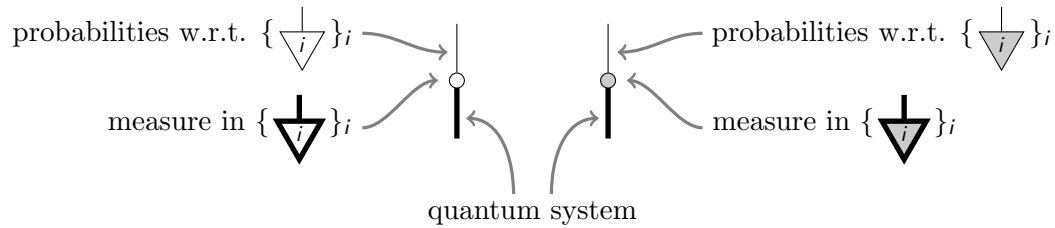
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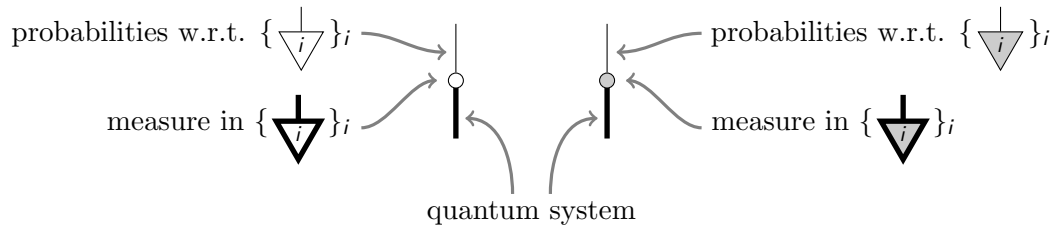
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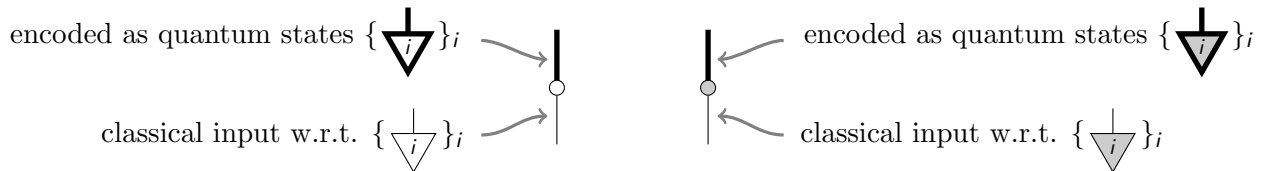


Two kinds of measurement

- Each spider induces a basis **measurement**:



- Their adjoints are **preparations**:



Measuring \Rightarrow preparing

- What happens when we **measure** then **prepare**? Decoherence.

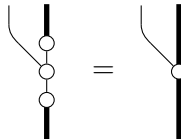
$$\left(\begin{array}{c} \text{---} \\ | \\ \text{---} \\ \rho \\ \text{---} \end{array} \right) = \sum_{ij} \rho_{ij} \begin{array}{c} | \\ \text{---} \\ i \\ \text{---} \end{array} \begin{array}{c} | \\ \text{---} \\ j \\ \text{---} \end{array} \right) \mapsto \left(\begin{array}{c} \text{---} \\ | \\ \circ \\ | \\ \circ \\ | \\ \text{---} \\ \rho \\ \text{---} \end{array} \right) = \sum_i \rho_{ii} \begin{array}{c} | \\ \text{---} \\ i \\ \text{---} \end{array} \begin{array}{c} | \\ \text{---} \\ i \\ \text{---} \end{array} \right)$$

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$$\left(\text{triangle with } \rho \text{ and a line} = \sum_{ij} \rho_{ij} \text{triangle } i \text{ triangle } j \right) \mapsto \left(\text{triangle with } \rho \text{ and a line with two circles} = \sum_i \rho_{ii} \text{triangle } i \text{ triangle } i \right)$$

- Decoherence models the situation where we **forget** the classical in the middle. However, we may have access to this classical data, i.e. if the detector clicks. So, we could just as well **keep a copy**.

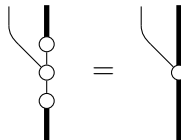


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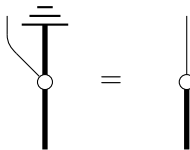
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- This lets us model **non-demolition** measurement devices. The demolition measurement can be recovered just by discarding the (quantum) output:

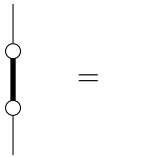


Preparing \Rightarrow measuring

- What happens when we **prepare** then **measure**? It depends on the choice of bases.

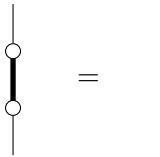
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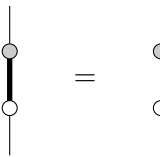


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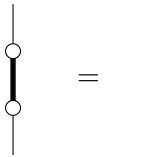


- ▶ The other extreme is:

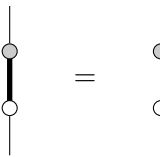


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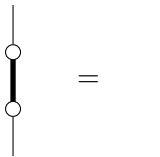
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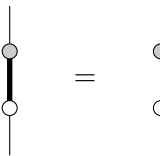
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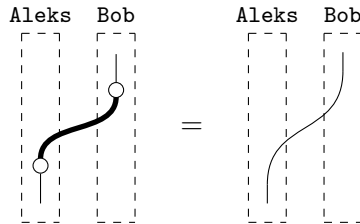
- ▶ In other words: (encode in ○) + (measure in ●) = (no data transfer)
- ▶ This is precisely what it means for two bases to be **complementary**

Complementarity – QKD

- ▶ This is at the heart of quantum key distribution.

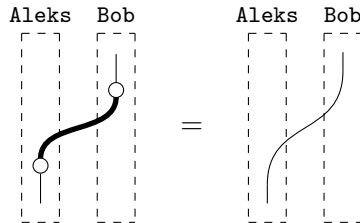
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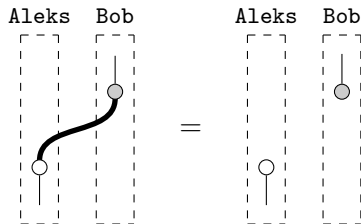


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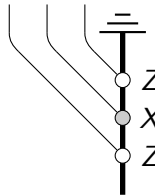


- ▶ When Bob measures in the **incorrect** basis, he gets noise:



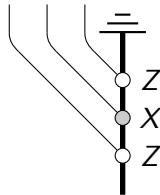
Complementarity – Stern-Gerlach

- Suppose \circ is a spin- Z measurement and \bullet is a spin- X measurement, then we could imagine a Stern-Gerlach type setup:

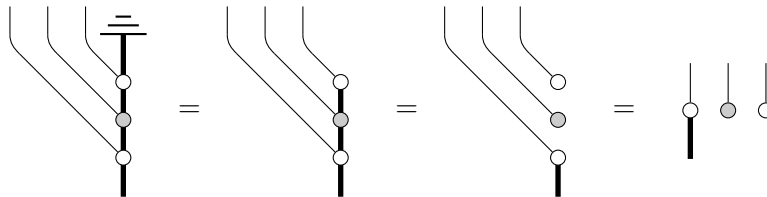


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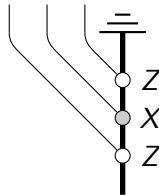


- Since Z and X are complementary, this simplifies as:

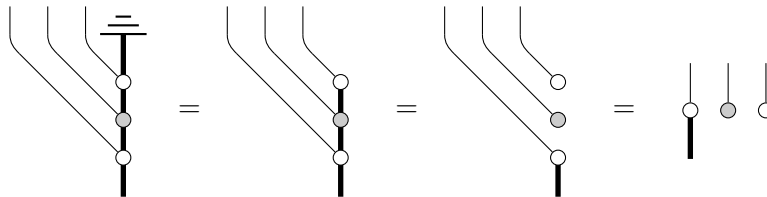


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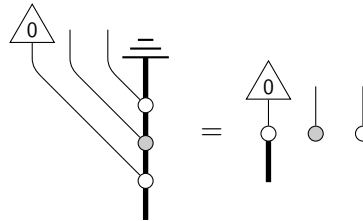
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- Thus the outcome of final measurement is **uniformly random**.
(recall $\circ = \text{flat probability distribution w.r.t. } \{\downarrow_j\}_j$).

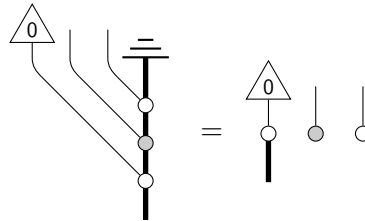
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- ▶ Since it disconnects, the output **stays random**, even when we post-select the first measurement to be spin-up (i.e. ‘block off the spin-down output’):



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- ▶ Since it disconnects, the output **stays random**, even when we post-select the first measurement to be spin-up (i.e. ‘block off the spin-down output’):



- ▶ We conclude from above that the X measurement (maximally) disturbs the system, w.r.t. the final Z measurement.

Complementarity \leftrightarrow Mutually unbiased bases

Definition

Two bases $\{\downarrow_j\}_j$ and $\{\downarrow_j\}_j$ are called *mutually unbiased* if:

$$\forall i, j. \quad \left| \frac{\downarrow_j}{\downarrow_i} \right| = \frac{1}{D}$$

or equivalently,

$$\forall i, j. \quad \left| \frac{\downarrow_j}{\downarrow_i} \right| = \frac{1}{\sqrt{D}}$$

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Theorem

Two bases are mutually unbiased iff they satisfy the *complementarity equation*:

$$\begin{array}{c} \bullet \\ | \\ \bullet \end{array} = \frac{1}{D} \begin{array}{c} \bullet \\ | \\ \bullet \end{array} \quad \text{or equivalently,} \quad \begin{array}{c} \bullet \\ \bigcirc \\ \bullet \end{array} = \frac{1}{D} \begin{array}{c} \bullet \\ | \\ \bullet \end{array}$$

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$$\begin{array}{c} \bullet \\ \text{thick line} \\ \circ \end{array} = \frac{1}{D} \begin{array}{c} \bullet \\ \text{thin line} \\ \circ \end{array} \quad \text{or equivalently,} \quad \begin{array}{c} \bullet \\ \text{loop} \\ \circ \end{array} = \frac{1}{D} \begin{array}{c} \bullet \\ \text{thin line} \\ \circ \end{array}$$

Proof.

(Compl. \Rightarrow MUB)

$$\begin{array}{c} \triangleup_j \\ \text{thick line} \\ \triangleup_i \end{array} = \begin{array}{c} \triangleup_j \\ \bullet \\ \text{thick line} \\ \circ \\ \triangleup_i \end{array} = \frac{1}{D} \begin{array}{c} \triangleup_j \\ \bullet \\ \text{thin line} \\ \circ \\ \triangleup_i \end{array} = \frac{1}{D}$$

(MUB \Rightarrow Compl.) follows similarly by comparing matrix entries.

□

General unbiased points

- Any pure state $\hat{\psi}$ is called *unbiased* w.r.t. to a basis if

$$\forall i. \quad \begin{array}{c} \triangleup \\ i \\ \hline \hat{\psi} \\ \triangleleft \end{array} = \lambda$$

where λ doesn't depend on i (and $= \frac{1}{D}$ when $\hat{\psi}$ is normalised).

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where λ doesn't depend on i (and $= \frac{1}{D}$ when $\hat{\psi}$ is normalised).

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- We could just as easily use this definition of unbiasedness for MUBs. Then, the complementarity equation follows just by evaluating on basis elements:

$$\begin{array}{c} \bullet \\ \text{---} \\ \circ \\ \text{---} \\ \triangleup i \end{array} = \begin{array}{c} \bullet \\ \text{---} \\ \triangleup i \end{array} = \frac{1}{D} \quad \bullet = \frac{1}{D} \quad \begin{array}{c} \bullet \\ \text{---} \\ \circ \\ \text{---} \\ \triangleup i \end{array}$$

Phase-states

- Killing the global phase, unbiased states can be parametrised by $D - 1$ complex phase factors:

$$\textcircled{\alpha} := \text{double} \left(\begin{array}{c} | \\ \nabla \\ 0 \end{array} + \sum_j e^{i\alpha_j} \begin{array}{c} | \\ \nabla \\ j \end{array} \right)$$

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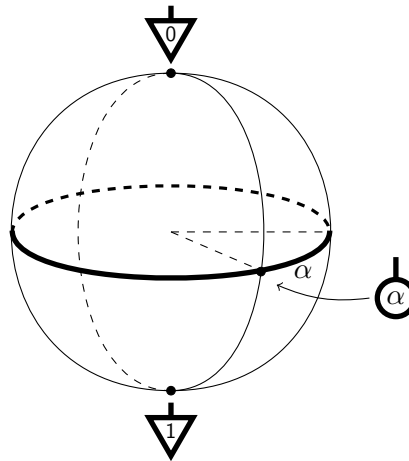
- Thus, unbiased states are also called *phase states*
- Specialising to the 2D case:

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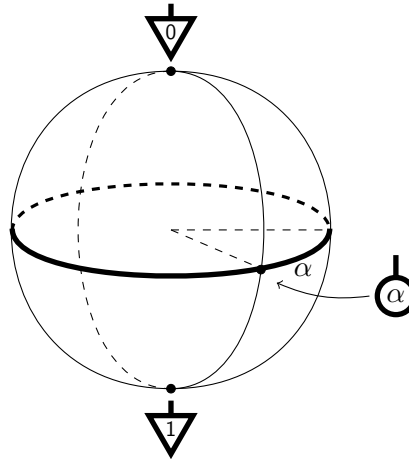
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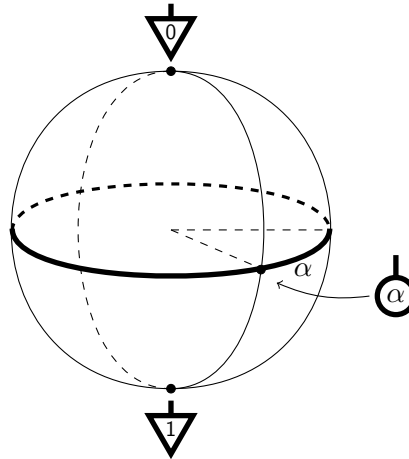


- Since decoherence projects to the axis of the Bloch ball, in particular:

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- ▶ So, phases get clobbered in the quantum/classical passage

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- How do we define **phase rotations**?

The phase group

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- ▶ A clue comes from the the **phase group** structure of spiders

$$\text{Spider}(\alpha, \beta) = \text{Circle}(\alpha + \beta)$$

$$\overline{(\text{Circle}(\alpha))} = \text{Circle}(-\alpha)$$

$$\text{Circle}(0) = \text{Leg}$$

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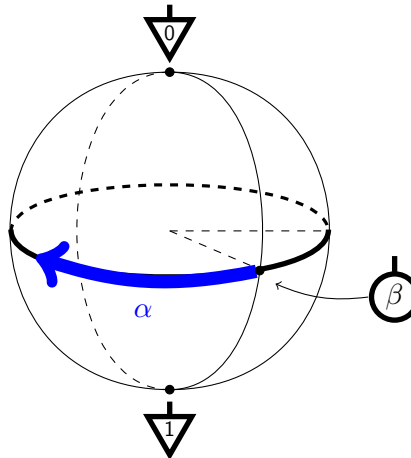
$$\text{Circle with } 0 = \text{Wire}$$

- ▶ If we multiply on the left (or the right) with a phase-state α , it performs an α rotation:

$$\text{Circle with } \alpha := \text{Spider with inputs } \alpha, \text{ wire}$$

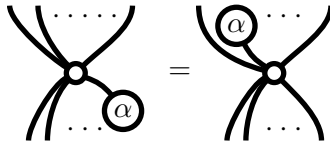
$$\vdots$$

$$\text{Circle with } \beta \mapsto \text{Circle with } \alpha + \beta$$



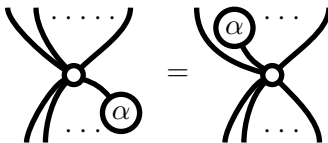
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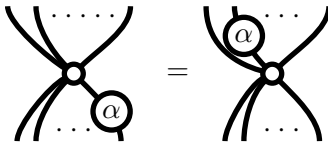


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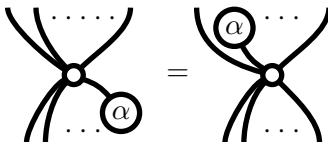


- A consequence is that **phase maps** commute through spiders:

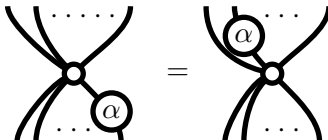


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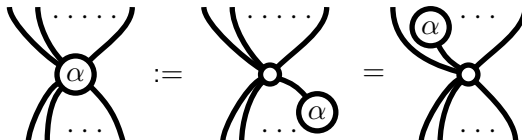
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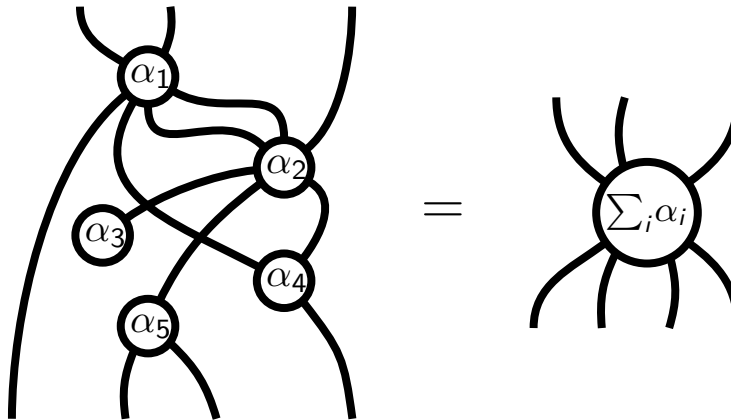


- We simplify our notation by letting spiders **eat connected phases**:



Generalised spider law

(phase group) + (spider fusion) = **(phase-spider fusion)**



Basis elements as phase states

- For a complementary pair \circ/\bullet the **basis states** of \circ are unbiased w.r.t. \bullet , so we could also write them as **phase states**. For $\circ := Z$ and $\bullet := X$,

$$\begin{array}{c} \downarrow \\ \nabla \\ 0 \end{array} = \begin{array}{c} \downarrow \\ \circ \\ 0 \end{array} \qquad \begin{array}{c} \downarrow \\ \nabla \\ 1 \end{array} = \begin{array}{c} \downarrow \\ \bullet \\ \pi \end{array}$$

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- So, since \circ gives us a way multiply phases, we can multiply ∇ -basis elements.

$$\nabla_i \nabla_j = \circ_{\alpha_i} \circ_{\alpha_j} = \circ_{\alpha_i + \alpha_j}$$

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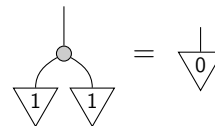
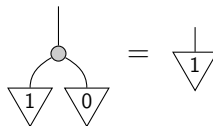
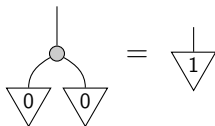
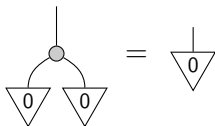
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- While in general, $\alpha_i + \alpha_j$ won't be another basis element, this *is* the case for Z/X :

$$\begin{array}{c} | \\ \bigcirc \\ \swarrow \searrow \\ \bigcirc \quad \bigcirc \\ 0 \quad 0 \end{array} = \begin{array}{c} | \\ \bigcirc \\ 0 \end{array} \qquad \begin{array}{c} | \\ \bigcirc \\ \swarrow \searrow \\ \bigcirc \quad \bigcirc \\ 0 \quad \pi \end{array} = \begin{array}{c} | \\ \bigcirc \\ \pi \end{array} \qquad \begin{array}{c} | \\ \bigcirc \\ \swarrow \searrow \\ \bigcirc \quad \bigcirc \\ \pi \quad 0 \end{array} = \begin{array}{c} | \\ \bigcirc \\ \pi \end{array} \qquad \begin{array}{c} | \\ \bigcirc \\ \swarrow \searrow \\ \bigcirc \quad \bigcirc \\ \pi \quad \pi \end{array} = \begin{array}{c} | \\ \bigcirc \\ 0 \end{array}$$


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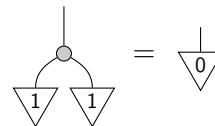
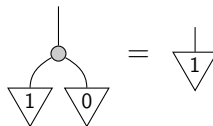
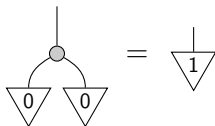
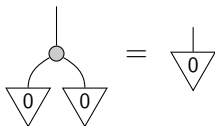
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namely, \mathbb{Z}_2 -multiplication.

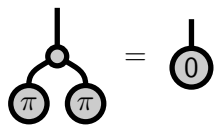
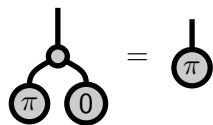
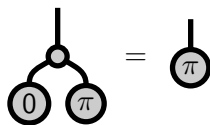
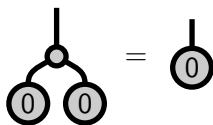
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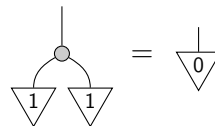
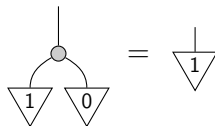
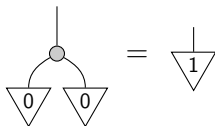
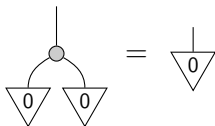
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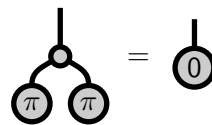
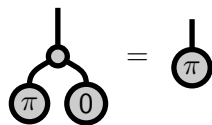
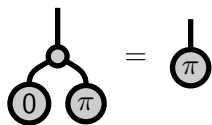
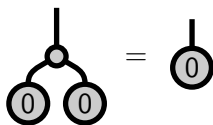
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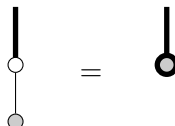
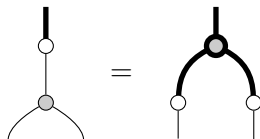


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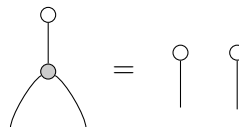
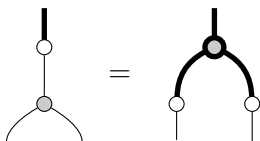
- ...and since $\{\triangle_j\}_j$ **encodes** the phase-states (via \circ preparation):



Strong complementarity

Definition

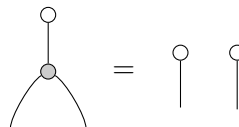
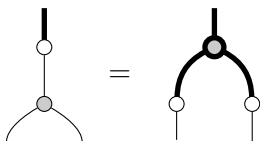
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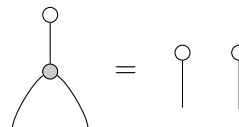
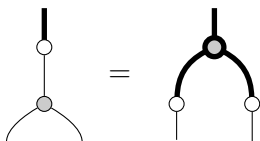
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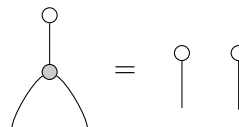
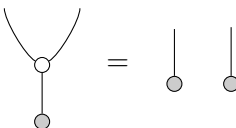
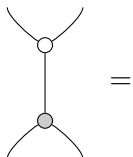
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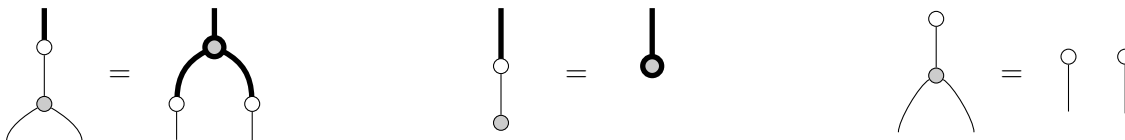
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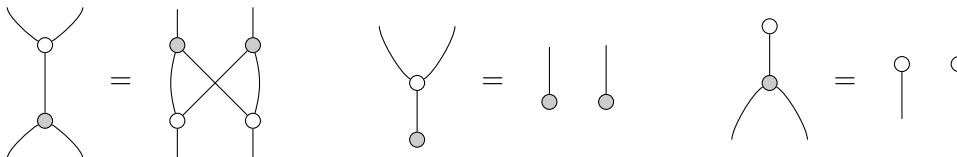
Strong complementarity

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- Strongly complementary pairs of spiders form **bi-algebras**!

Strong complementarity \Rightarrow complementarity

Theorem

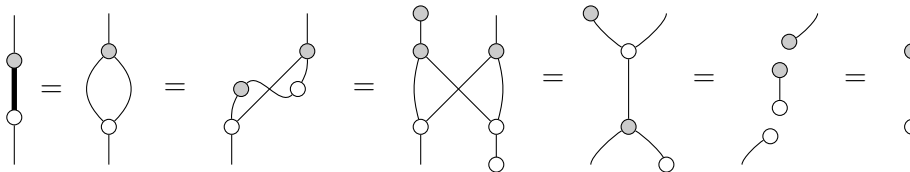
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□

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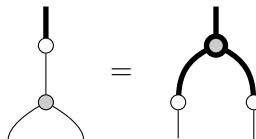
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Proof.

(sketch)  acts as a group operation on $\{\triangle_j\}_j$. Fixing *which* group operation totally characterises , and hence $\{\triangle_j\}_j$. □

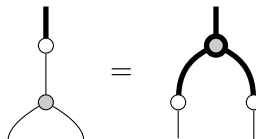
Making sense of phase-multiply


- We tried to give some (pseudo-)operational interpretation of this equation:

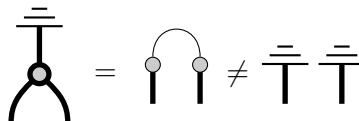


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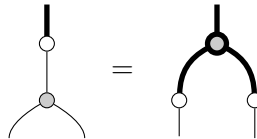
- ▶ But it falls down because, while  is a good **quantum map**, it isn't **causal**:




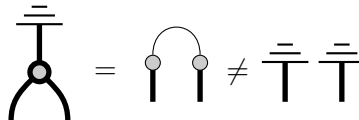
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So it isn't *physical*.


- ▶ This is because, it is both **pure**, and **it throws stuff away**. E.g. for the Z/X example before, it is \mathbb{Z}_2 -multiply, a.k.a. XOR.

Making sense of phase-multiply

- ▶ However,  is *part* of a physical map, if we play a standard trick from quantum computing. We simply **copy** (some of) the input:

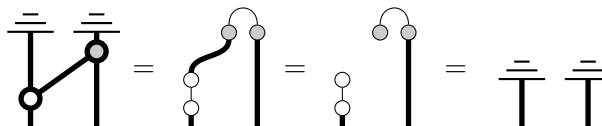


Making sense of phase-multiply


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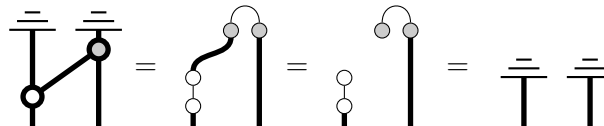


Making sense of phase-multiply

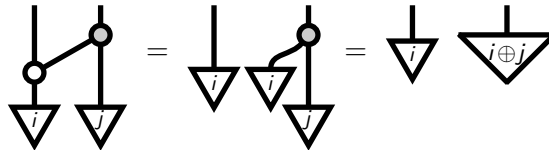
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- ▶ Returning to the Z/X example, this in fact gives us a CNOT gate:

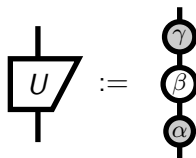


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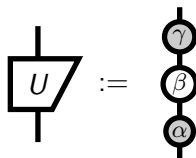
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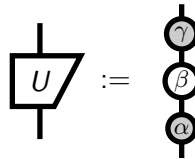
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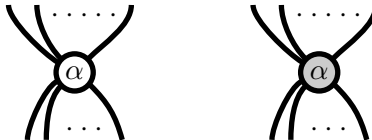
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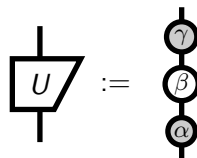
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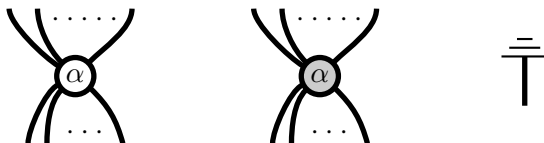
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Corollary

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Completeness?

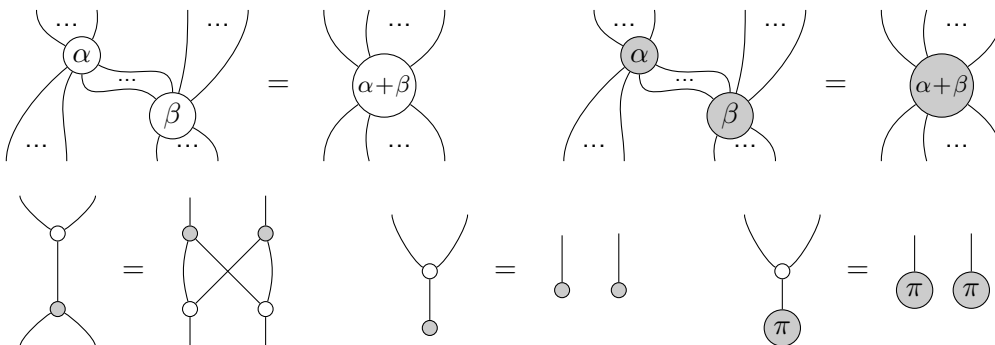
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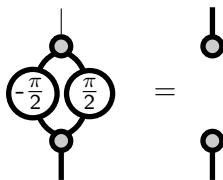
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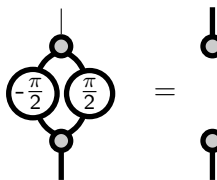
Clifford maps

- But there are still some equations that can't be proven, e.g.



Clifford maps

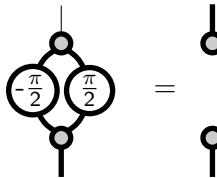
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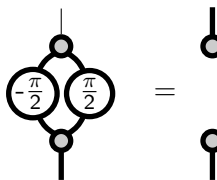
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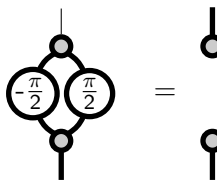
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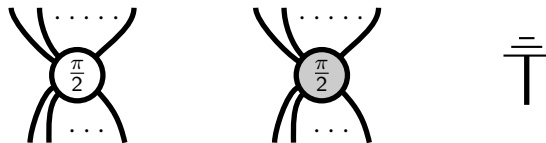
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Definition

Let the family of *Clifford maps* consist of any map generated by:



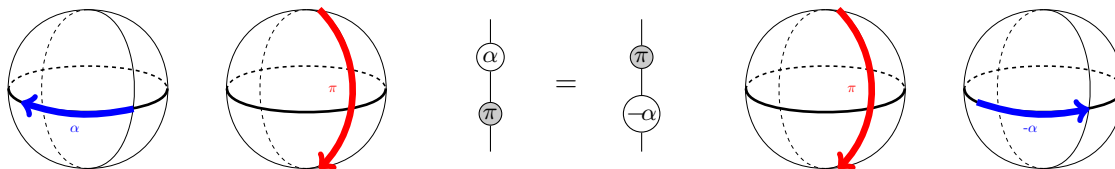
(*Clifford circuit* := unitary Clifford map)

Geometry

- ▶ We nearly have a complete set of equations for the Clifford maps, but we're missing some info about the **geometry of the Bloch sphere**

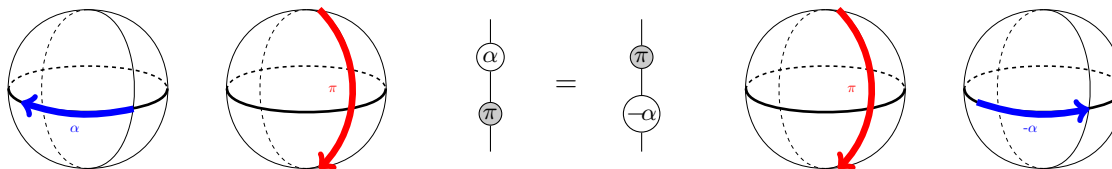
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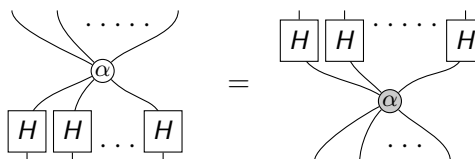


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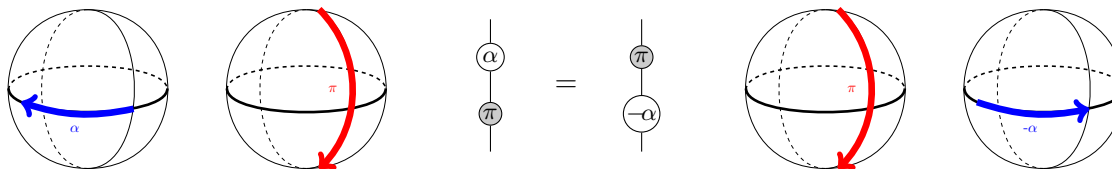


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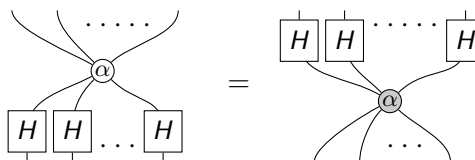


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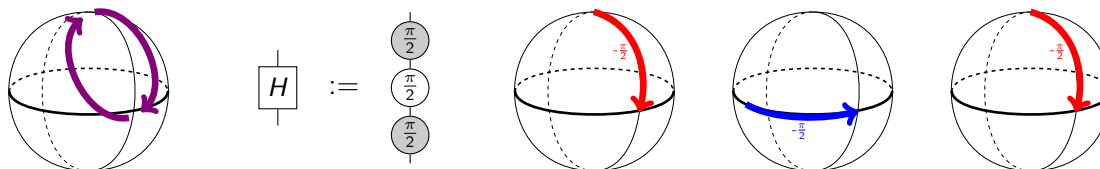
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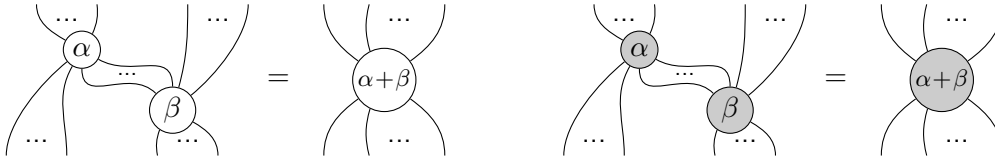


The ZX-Calculus

Definition

The *ZX-calculus* consists of:

- Two **spider-fusion** rules:

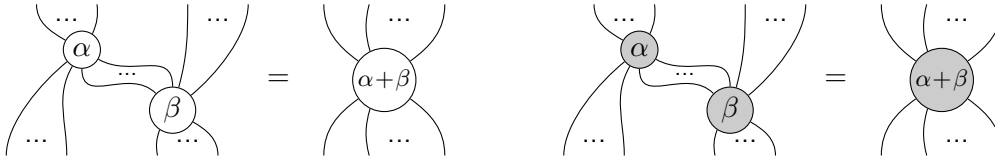


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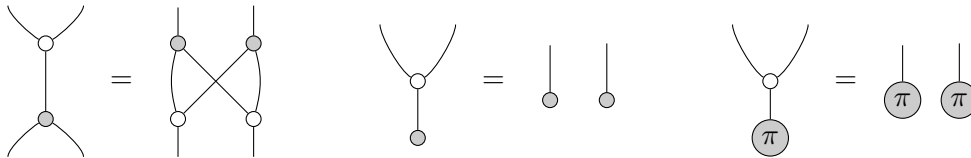
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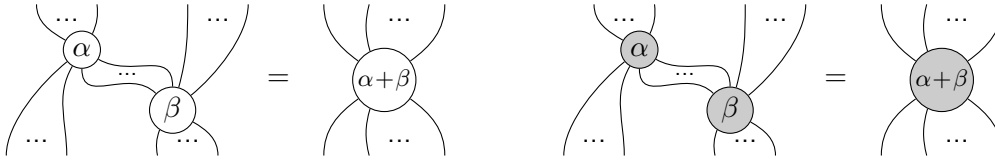


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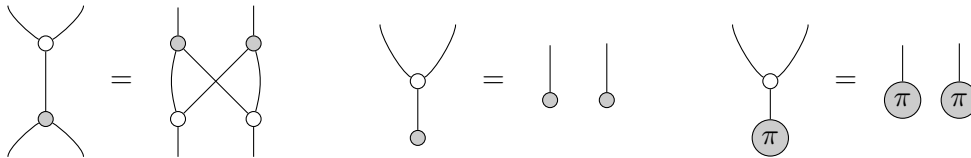
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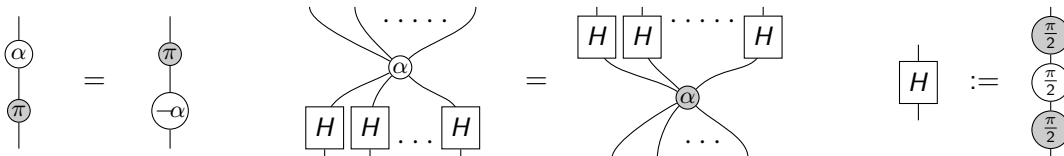
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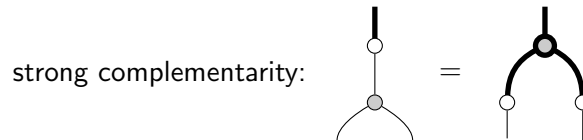
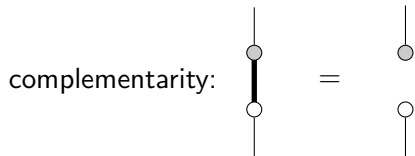
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- ▶ ...but it is complete for at least one other fragment: **single-qubit unitaries** with $\frac{\pi}{4}$ **phase maps** (a.k.a. Clifford + T).

Summary

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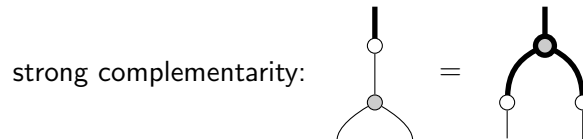
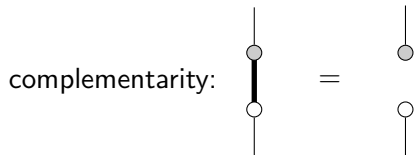
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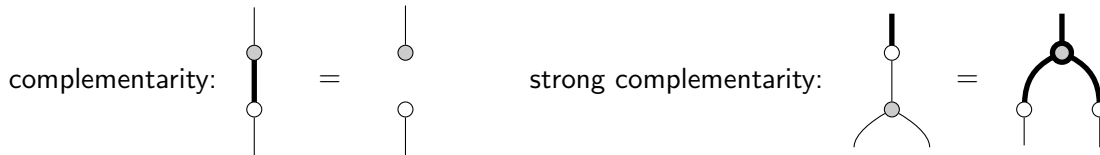
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- ▶ ...and demonstrate a tool for automating calculation in ZX: **QuantoDerive**