

Lecture 7

Today: define the process theory of linear maps
 String diagrams + 3 extra ingredients

BASES

NOTE: in most process theories,

$$\text{if } \forall \begin{array}{c} |A \\ \text{f} \\ |B \\ \hline \psi \end{array}, \begin{array}{c} |B \\ f \\ |A \\ \hline \psi \end{array} = \begin{array}{c} |B \\ g \\ |A \\ \hline \psi \end{array} \text{ then } \begin{array}{c} |B \\ f \\ |A \\ \hline \end{array} = \begin{array}{c} |B \\ g \\ |A \\ \hline \end{array}.$$

(extensionality, well-pointed, ...)

DEF A basis for a type A is a minimal set of states:

$$B := \left\{ \begin{array}{c} |A \\ 1 \\ \hline \end{array}, \dots, \begin{array}{c} |A \\ n \\ \hline \end{array} \right\}$$

such that:

$$\text{if for all } \begin{array}{c} |A \\ i \\ \hline \end{array} \in B, \begin{array}{c} |B \\ f \\ |A \\ \hline i \end{array} = \begin{array}{c} |B \\ g \\ |A \\ \hline i \end{array} \text{ then } f = g.$$

INTUITION basis states are "reference points" for a process

DEF The dimension of a system is the size of (the smallest) basis.

DEF Two states are orthogonal if $\langle \psi | \phi \rangle = 0$.

DEF An orthonormal basis $B = \{ \underline{\psi}_i \}_{i=1..n}^{(ONB)}$ is a basis

where

$$\underline{\psi}_i^j = \begin{cases} 1 & i=j \\ 0 & i \neq j \end{cases} \quad \hookrightarrow \delta_i^j \leftarrow \text{Kronecker } \delta.$$

DEF self-conjugate ONB: $\underline{\psi}_i^A = \underline{\psi}_i^1 =: \underline{\psi}_i^A$

MATRICES

Thm For bases:

$$B = \{ \underline{\psi}_i^A \}_{i=1..n} \quad B' = \{ \underline{\psi}_i^B \}_{i=1..n}$$

$$\left(\forall i, j. \quad \underline{\psi}_i^j = \underline{\psi}_j^i \right) \Rightarrow \underline{\psi}_i^B = \underline{\psi}_i^A$$

Pf

$$\forall j \left(\forall i. \quad \underline{\psi}_i^j = \underline{\psi}_j^i \right) \Rightarrow \forall j \left(\underline{\psi}_i^B = \underline{\psi}_i^A \right)$$

$$h_j \left(\begin{array}{c} f^j \\ \hline A \end{array} = \begin{array}{c} g^j \\ \hline A \end{array} \right) \Rightarrow h_j \left(\begin{array}{c} f \\ \hline A \end{array} = \begin{array}{c} g \\ \hline A \end{array} \right)$$

$$\Rightarrow \begin{array}{c} f \\ \hline B \end{array} = \begin{array}{c} g \\ \hline B \end{array} \Rightarrow \begin{array}{c} f \\ \hline A \end{array} = \begin{array}{c} g \\ \hline A \end{array}. \quad \blacksquare$$

NOTATION Let $f_i^j := \begin{array}{c} f^j \\ \hline A \\ \downarrow i \end{array}$ ← number, depends on bases

DEF The set of numbers $\{f_i^j\}_{i=1..m, j=1..n}$ is called the matrix of f .

Think of f_i^j row index
column index (e.g. f_{ji} in linear alg.)

e.g. if $m=3$ and $n=2$, matrix of f

$$\begin{pmatrix} f_1^1 & f_2^1 & f_3^1 \\ f_1^2 & f_2^2 & f_3^2 \end{pmatrix}$$

MATRIX TRANSPOSE

$$\begin{pmatrix} f_1^1 & f_2^1 & f_3^1 \\ f_1^2 & f_2^2 & f_3^2 \end{pmatrix}^T := \begin{pmatrix} f_1^1 & f_1^2 \\ f_2^1 & f_2^2 \\ f_3^1 & f_3^2 \end{pmatrix}$$

(with resp. to a self-conj ONB)

T_{Hm} The matrix of $\begin{pmatrix} A \\ f \\ B \end{pmatrix} := {}^A \begin{pmatrix} f \\ B \end{pmatrix}_B$ is $(f^\top)_i^j = f_j^i$.

Pf

$$(f^\top)_i^j := \begin{pmatrix} j \\ f \\ i \end{pmatrix} = \begin{pmatrix} j \\ f \\ i \end{pmatrix} \xrightarrow{\text{self-conj}} \begin{pmatrix} i \\ f \\ j \end{pmatrix} = f_j^i. \quad \blacksquare$$

on " " " "

Similarly, the conjugate and conjugate-transpose:

$$(\bar{f})_i^j = \overline{f_i^j} \quad \begin{pmatrix} f \\ B \end{pmatrix} \mapsto \begin{pmatrix} \bar{f} \\ B \end{pmatrix} \quad \left(\begin{array}{ccc} \bar{\cdot} & \bar{\cdot} & \bar{\cdot} \\ \vdots & \vdots & \vdots \end{array} \right)$$

$$(f^\dagger)_i^j = \overline{f_j^i} \quad \begin{pmatrix} f \\ B \end{pmatrix} \mapsto \begin{pmatrix} f \\ \bar{B} \end{pmatrix} \quad \left(\begin{array}{ccc} \bar{\cdot} & \bar{\cdot} & \bar{\cdot} \\ \vdots & \vdots & \vdots \end{array} \right)$$

Process

$$\begin{pmatrix} B \\ f \\ A \end{pmatrix}$$



MATRIX

$$\left(\begin{array}{ccc} \cdot & \cdot & \cdot \\ \vdots & \vdots & \vdots \end{array} \right)$$

MATRIX

$$\left(\begin{array}{ccc} \cdot & \cdot & \cdot \\ \vdots & \vdots & \vdots \end{array} \right)$$

Process

?

SUPPOSE I HAVE $\left(\begin{array}{ccc} \cdot & \cdot & \cdot \\ \cdot & g_1^1 & \cdot \\ \cdot & \cdot & \cdot \end{array} \right)$, then I can make the proc:

$$\begin{pmatrix} f \\ g \\ B \end{pmatrix} := \begin{pmatrix} f \\ g \\ B \end{pmatrix} \xrightarrow{i} \begin{pmatrix} f \\ g \\ i \end{pmatrix}$$

then $\begin{pmatrix} f \\ g \\ B \end{pmatrix}$ has matrix:

$$\left(\begin{array}{cccc} 0 & 0 & 0 & \cdots \\ \cdots & 0 & g_1^1 & 0 \\ \cdots & 0 & 0 & 0 \end{array} \right) \xleftarrow{\text{j}^{\text{th}} \text{ row}} \xrightarrow{\text{i}^{\text{th}} \text{ column}}$$

DEF A process theory has sums if, for any two processes of the same type

$$\begin{array}{c} \boxed{f} \\ \text{in} \end{array}, \begin{array}{c} \boxed{g} \\ \text{in} \end{array} \rightsquigarrow \begin{array}{c} \boxed{f+g} \\ \text{in} \end{array}$$

1. associative, commutative, and has unit \emptyset .

$$f + (g + h) = f + (g + h) \quad f + g = g + f \quad \begin{array}{c} \boxed{f} \\ \text{in} \end{array}, \begin{array}{c} \boxed{\emptyset} \\ \text{in} \end{array} \rightsquigarrow f + \emptyset = f.$$

Let $\sum_i \begin{array}{c} \boxed{f_i} \\ \text{in} \end{array} := f_1 + \dots + f_n$

2. Sums distribute over diagrams

$$\left(\sum_i \begin{array}{c} \boxed{h_i} \\ \text{in} \end{array} \right) \begin{array}{c} \boxed{g} \\ \text{out} \end{array} = \sum_i \begin{array}{c} \boxed{g} \\ \text{out} \end{array} \begin{array}{c} \boxed{h_i} \\ \text{in} \end{array} \begin{array}{c} \boxed{f} \\ \text{out} \end{array}$$

3. Sums preserve adjoints

$$\left(\sum_i \begin{array}{c} \boxed{f_i} \\ \text{in} \end{array} \right)^* = \sum_i \begin{array}{c} \boxed{f_i} \\ \text{out} \end{array}$$

DEF The process theory of linear maps is the theory where:

1. for each $D \in \mathbb{N}$, there is a system \cong ONB of size D .
2. processes admit sums
3. the numbers are complex numbers.
(ie. every process \boxed{x} corresponds to a complex number)

\rightarrow (we call a system in linear maps a Hilbert space)

Lecture 8

From the past

EXAMPLES OF SUMS DIST. OVER DiAGS.

* linearity.

$$\frac{1_B}{\boxed{f}} \circ \left(\frac{1_A}{\boxed{\psi}} \right) = \frac{1_A}{\boxed{f} \circ \boxed{\psi}} = \lambda \cdot \left(\frac{1_B}{\boxed{f}} \circ \frac{1_A}{\boxed{\psi}} \right)$$

$$f(\lambda \cdot v) = \lambda \cdot f(v)$$

$$f(v_1 + v_2) = f(v_1) + f(v_2)$$

$$f(\sum_i v_i) = \sum_i f(v_i)$$

$$\frac{1_B}{\boxed{f}} \circ \left(\sum_i \frac{1_A}{\boxed{v_i}} \right) = \frac{1_B}{\boxed{f} \circ \boxed{w}}$$

$$\text{where } \boxed{w} = \sum_i \boxed{v_i}$$

$$= \sum_i \frac{1_f}{\boxed{v_i}} = \sum_i \left[\frac{1_f}{\boxed{v_i}} \circ \frac{1_A}{\boxed{v_i}} \right]$$

* linearity of $\langle - | - \rangle$

$$\langle w | v_1 + v_2 \rangle = \langle w | v_1 \rangle + \langle w | v_2 \rangle$$

$$\langle w | \sum_i v_i \rangle = \sum_i \langle w | v_i \rangle$$

* conjugate-linearity of $\langle - | - \rangle$

$$\langle \sum_i v_i | w \rangle = \sum_i \langle v_i | w \rangle$$

$$\langle \sum_i \lambda_i v_i | w \rangle = \sum_i \bar{\lambda}_i \langle v_i | w \rangle$$

* bilinearity of \otimes . $(\sum_i \frac{1_f}{\boxed{v_i}}) \otimes \frac{1_g}{\boxed{w}} = \sum_i \left[\frac{1_f}{\boxed{v_i}} \otimes \frac{1_g}{\boxed{w}} \right]$

$$\frac{1_f}{\boxed{v}} \otimes (\sum_i \frac{1_g}{\boxed{w}}) = \sum_i \left[\frac{1_f}{\boxed{v}} \otimes \frac{1_g}{\boxed{w}} \right]$$

$$\textcircled{2} \quad \frac{1_f}{\boxed{f}} = \left(\sum_{i=0}^1 \dots \right) \otimes \frac{1_f}{\boxed{f}} = \sum_{i=0}^1 \frac{1_f}{\boxed{f}} = \frac{1_f}{\boxed{f}} + \frac{1_f}{\boxed{f}}$$

Matrix M \rightsquigarrow process f

Theorem For any matrix M with entries m_{ij}^j , the process:

$$\boxed{f} := \sum_{ij} \begin{array}{c} \nearrow \\ m_{ij}^j \\ \searrow \end{array}$$

has matrix M .

Pf matrix of f has entries $f_i^j := \begin{array}{c} \nearrow \\ f \\ \searrow \end{array}$.

$$f_i^j = \begin{array}{c} \nearrow \\ f \\ \searrow \end{array} = \left(\sum_{k,l} m_{k,l}^l \begin{array}{c} \nearrow \\ k \\ \searrow \end{array} \right) = \sum_{k,l} m_{k,l}^l \begin{array}{c} \nearrow \\ k \\ \searrow \\ i \end{array} = \sum_{k,l} m_{k,l}^l \delta_{i,k} \delta_{l,j}$$

$$= \underbrace{\sum_{k \neq i, l} m_{k,l}^l \delta_{i,k} \delta_{l,j}}_0 + \sum_l m_{i,l}^l \delta_{i,i} \delta_{l,j} = \sum_l m_{i,l}^l \delta_{l,j} = m_i^j \quad \blacksquare$$

NOTE $\rightsquigarrow \int_A$ has matrix $\begin{pmatrix} 1 & 0 & \dots \\ 0 & 1 & \dots \\ \vdots & \vdots & \ddots \end{pmatrix}$, so $\int_A = \sum_j \delta_j^i \begin{array}{c} \nearrow \\ j \\ \searrow \end{array} = \sum_i \begin{array}{c} \nearrow \\ i \\ \searrow \end{array}$

COROLLARY $\int_A = \sum_i \begin{array}{c} \nearrow \\ i \\ \searrow \end{array}$

Thm If f has matrix M_f and g has matrix M_g , then

the matrix δ_f^g 

is $M_g \cdot M_f$.

 \leftarrow matrix multiplication

Pf

$$M_g M_f = \begin{pmatrix} g_1 & \cdots & g_n \end{pmatrix} \begin{pmatrix} f_1 \\ \vdots \\ f_n \end{pmatrix} = \begin{pmatrix} \cdots & \cdots & \cdots \\ \cdots & \sum_k g_k^j f_i^k & \cdots \\ \cdots & \cdots & \cdots \end{pmatrix} \leftarrow \text{if } j$$

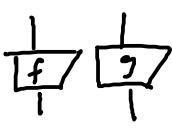
need to show

the matrix δ_f^g 

has entries $\sum_k g_k^j f_i^k$.

 $\left\{ \begin{array}{c} \uparrow \\ \downarrow \end{array} \right\}_{k=1..n}$ basis for B

$$(M_{gof})_i^j = \begin{pmatrix} \overset{j}{\uparrow} \\ \overset{i}{\downarrow} \\ \begin{matrix} g \\ f \\ A \end{matrix} \end{pmatrix} = \sum_{k=1}^n \begin{pmatrix} \overset{j}{\uparrow} \\ \overset{k}{\downarrow} \\ \begin{matrix} g \\ f \\ A \end{matrix} \end{pmatrix} = \sum_k g_k^j f_i^k. \quad \square$$

Thm The matrix δ_f^g 

 \leftarrow tensor product

$$\begin{pmatrix} f_1' & f_2' \\ f_1^2 & f_2^2 \end{pmatrix}_{2 \times 2} \otimes \begin{pmatrix} g_1' & g_2' \\ g_1^2 & g_2^2 \end{pmatrix}_{2 \times 2}^G = \begin{pmatrix} f_1' G & f_2' G \\ f_1^2 G & f_2^2 G \end{pmatrix}$$

$$= \begin{pmatrix} f_1' g_1' & f_1' g_2' & f_2' g_1' & f_2' g_2' \\ f_1^2 g_1' & f_1^2 g_2' & f_2^2 g_1' & f_2^2 g_2' \\ f_1' g_1^2 & f_1' g_2^2 & f_2' g_1^2 & f_2' g_2^2 \\ f_1^2 g_1^2 & f_1^2 g_2^2 & f_2^2 g_1^2 & f_2^2 g_2^2 \end{pmatrix}_{4 \times 4}$$

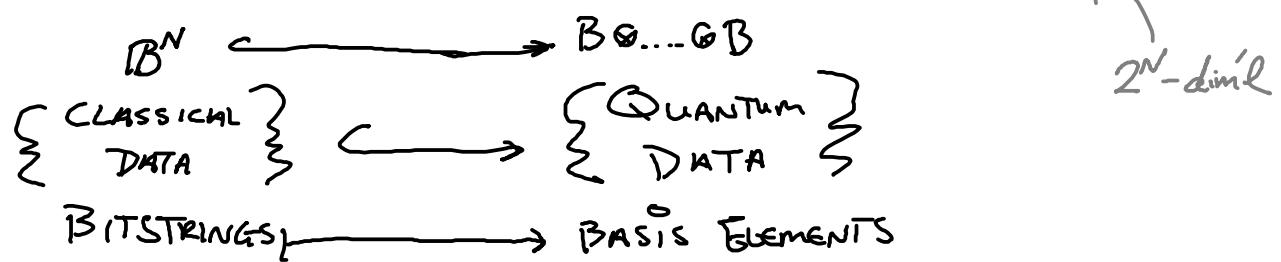
$$\begin{pmatrix} \cdot & \cdot & \cdot \\ \cdot & \cdot & \cdot \\ \cdot & \cdot & \cdot \end{pmatrix}_{m \times n} \otimes \begin{pmatrix} \cdot & \cdot \\ \cdot & \cdot \\ \cdot & \cdot \end{pmatrix}_{m' \times n'} = \begin{pmatrix} \cdot & \cdot & \cdot & \cdot \\ \cdot & \cdot & \cdot & \cdot \\ \cdot & \cdot & \cdot & \cdot \\ \cdot & \cdot & \cdot & \cdot \end{pmatrix}_{mm' \times nn'}$$

The reason: if $\left\{ \begin{smallmatrix} |A \\ \downarrow i \end{smallmatrix} \right\}_{i=1..m}$ is a basis for A and if $\left\{ \begin{smallmatrix} |B \\ \downarrow j \end{smallmatrix} \right\}_{j=1..m'}$ is a basis for B, then $\left\{ \begin{smallmatrix} |A & |B \\ \downarrow i & \downarrow j \end{smallmatrix} \right\}_{i=1..m, j=1..m'}$ is a basis for $A \otimes B$.

$$\left(\forall_{ij} \begin{smallmatrix} |f \\ \downarrow_i \downarrow_j \end{smallmatrix} = \begin{smallmatrix} |g \\ \downarrow_i \downarrow_j \end{smallmatrix} \right) \Rightarrow \begin{smallmatrix} |c \\ \downarrow_f \downarrow_B \end{smallmatrix} = \begin{smallmatrix} |g \\ \downarrow_h \downarrow_o \end{smallmatrix}$$

Ex assume B has basis $\left\{ \begin{smallmatrix} |1^B \\ \downarrow_0 \end{smallmatrix}, \begin{smallmatrix} |1^B \\ \downarrow_1 \end{smallmatrix} \right\}$ \leftarrow 2-dim'l
 Then $B \otimes B$ has basis $\left\{ \begin{smallmatrix} |11^B \\ \downarrow_0 \downarrow_0 \end{smallmatrix}, \begin{smallmatrix} |11^B \\ \downarrow_0 \downarrow_1 \end{smallmatrix}, \begin{smallmatrix} |11^B \\ \downarrow_1 \downarrow_0 \end{smallmatrix}, \begin{smallmatrix} |11^B \\ \downarrow_1 \downarrow_1 \end{smallmatrix} \right\}$ \leftarrow 4-dim'l

and $\underbrace{B \otimes \dots \otimes B}_N$ has basis $\left\{ \begin{smallmatrix} |1\dots 1 \\ \downarrow_0 \dots \downarrow_0 \end{smallmatrix}, \begin{smallmatrix} |1\dots 1 \\ \downarrow_0 \dots \downarrow_1 \end{smallmatrix}, \dots, \begin{smallmatrix} |1\dots 1 \\ \downarrow_1 \dots \downarrow_1 \end{smallmatrix} \right\}$



Lecture 9

linear maps \rightsquigarrow quantum maps

1. Systems \mathbb{C}^n for every $n \in \mathbb{N}$ with an ONB $\{\downarrow, \dots, \downarrow\}$
2. Sums
3. numbers are complex numbers \mathbb{C}

* in particular, the trivial system $I = \mathbb{C} (= \mathbb{C}^1)$ has an ONB $\{\downarrow, \dots, \downarrow\}$

$$\underbrace{\mathbb{C}^m \otimes \mathbb{C}^n}_{= \mathbb{C}^{m \cdot n}}$$

$$(\downarrow \psi = \downarrow \phi) \Rightarrow \downarrow \psi = \downarrow \phi$$

Therefore: $\begin{pmatrix} \downarrow \psi \\ \vdots \\ \downarrow \phi \end{pmatrix} \leftrightarrow \underbrace{\left(\quad \right)}_{\dim A} \} \dim B$

so $\begin{pmatrix} \downarrow \psi \\ \vdots \\ \downarrow \phi \end{pmatrix}^A$ are column vectors, $\begin{pmatrix} \downarrow \psi \\ \vdots \\ \downarrow \phi \end{pmatrix}^B$ row vectors, $\begin{pmatrix} \downarrow \psi \\ \vdots \\ \downarrow \phi \end{pmatrix}^C \leftrightarrow (\lambda)$

What's special about \mathbb{C} ?

— In general, we can write a process in matrix form:

$$\begin{pmatrix} \downarrow \psi \\ \vdots \\ \downarrow \phi \end{pmatrix}^B = \sum_{ij} m_{ij}^B \begin{pmatrix} \downarrow \psi \\ \vdots \\ \downarrow \phi \end{pmatrix}^A$$

M-N complex numbers

Thm (diagonalisation) if $\frac{f}{A} = \frac{F}{A}$, then there exists some ONB $\{\frac{1}{i}\}_{i=1}^n$ such that:

$$\text{eigenbasis} \quad \frac{f}{A} = \sum_j r_j \frac{1}{i} \quad \Leftrightarrow \quad \begin{pmatrix} r_1 & & 0 \\ & \ddots & 0 \\ 0 & & r_n \end{pmatrix}$$

$\dim(A)$ parameters

$$\frac{f}{A} = \left(\sum_j r_j \frac{1}{i} \right) = r_i \frac{1}{i} = r_i \uparrow \quad \text{eigenvalue.}$$

eigenvectors of f

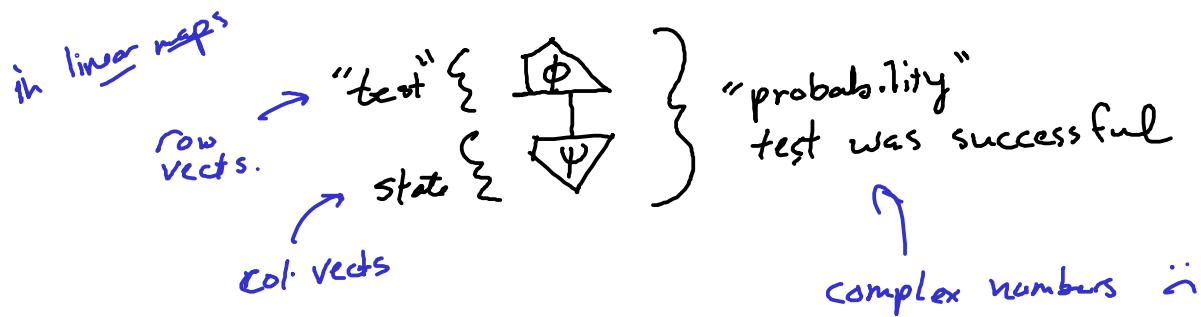
FURTHERMORE: f self-adjoint $\Rightarrow r_i = \bar{r}_i$ (real numbers)

if f is positive $\Rightarrow r_i \geq 0$ (+-ve real numbers)

Thm In linear maps, $\exists \frac{1}{\lambda} \cdot \lambda = \frac{1}{\psi}$ iff λ is real, ≥ 0 .

Pf (book)

Recall the Generalised Born rule:



Fix: for any complex number γ , $\bar{\gamma}\gamma = |\gamma|^2$ is a +ive number

$$0 \leq \begin{array}{c} \hat{\psi} \\ \downarrow \\ \psi \end{array} \quad \begin{array}{c} \hat{\phi} \\ \downarrow \\ \phi \end{array} \leq 1 \quad \left. \begin{array}{c} \uparrow \\ \psi, \phi \\ \text{normalised} \end{array} \right\} \text{probability}$$

...but we've lost the elegance of the Born rule.

Fix let: $\hat{A} := \begin{array}{c} \hat{A} \\ \downarrow \\ \psi \end{array} := \begin{array}{c} \hat{A} \\ \downarrow \\ \psi \end{array} \begin{array}{c} \downarrow \\ \psi \end{array} \quad \hat{A} = \begin{bmatrix} \cdot & \cdot \\ \cdot & \cdot \end{bmatrix} \leftarrow A \otimes A$

$$\hat{\phi} := \begin{array}{c} \hat{\phi} \\ \downarrow \\ \phi \end{array} = \begin{array}{c} \hat{\phi} \\ \downarrow \\ \phi \end{array} \begin{array}{c} \downarrow \\ \phi \end{array}$$

$$\text{test} \left\{ \begin{array}{c} \text{top} \\ \downarrow \end{array} \right. \phi =: \begin{array}{c} \text{top} \quad \text{top} \\ \downarrow \quad \downarrow \\ \psi \quad \psi \end{array} \left. \begin{array}{c} \text{probability} \end{array} \right\}$$

more generally, we define:

$$\text{double} \left(\begin{array}{c} \text{top} \\ \downarrow \\ f \\ \downarrow \\ \text{bottom} \end{array} \right) = \frac{\text{top}}{\|f\|} := \begin{array}{c} \text{top} \\ \downarrow \\ f \quad f \\ \downarrow \quad \downarrow \\ \text{bottom} \end{array}$$

DEF The theory of pure quantum maps has:

- * types \hat{A} for every Hilbert space A (i.e. $\widehat{\mathbb{C}^n}$ for all n)
- * processes $\hat{f}: \hat{A} \rightarrow \hat{B}$ for every linear map $f: A \rightarrow B$.

NOTE, our convention for $\begin{array}{c} \text{top} \\ \downarrow \\ \psi \\ \downarrow \\ \text{bottom} \end{array}$ will be:

$$\text{double} \left(\begin{array}{c} \text{top} \\ \downarrow \\ \psi \\ \downarrow \\ \text{bottom} \end{array} \right) =: \begin{array}{c} \text{top} \quad \text{top} \\ \downarrow \quad \downarrow \\ \hat{A} \quad \hat{B} \\ \downarrow \quad \downarrow \\ \psi \quad \psi \end{array} := \begin{array}{c} \text{top} \quad \text{top} \quad \text{top} \quad \text{top} \\ \downarrow \quad \downarrow \quad \downarrow \quad \downarrow \\ \hat{A} \quad \hat{A} \quad \hat{A} \quad \hat{B} \\ \downarrow \quad \downarrow \quad \downarrow \quad \downarrow \\ \psi \quad \psi \quad \psi \quad \psi \end{array}$$

$$\text{double}(\cup) = \cup \circ \text{double} = \hat{\cup} = \cup \hat{A}$$

$$\text{double}(\cap) = \cap \circ \text{double} = \cap \hat{A}$$

$$\cap = \cap \cup = \parallel = \parallel \hat{A}$$

Thm Doubling preserves string diagrams.

$$\text{double}\left(\begin{array}{c} \text{a} \\ \text{b} \\ \text{c} \\ \text{d} \end{array}\right) = \begin{array}{c} \hat{a} \\ \hat{b} \\ \hat{c} \\ \hat{d} \end{array}$$

PF (sketch)

* \circ and \otimes (e.g. $\text{double}(g \circ f) = \text{double}(g) \circ \text{double}(f)$)

* swaps, cups, caps

* adjoints

(book)



* on states, $\overline{\psi} = \overline{\phi} \Rightarrow \overline{\hat{\psi}} = \overline{\hat{\phi}}$, but
 $\overline{\psi} \overline{\psi} = \overline{\phi} \overline{\phi}$

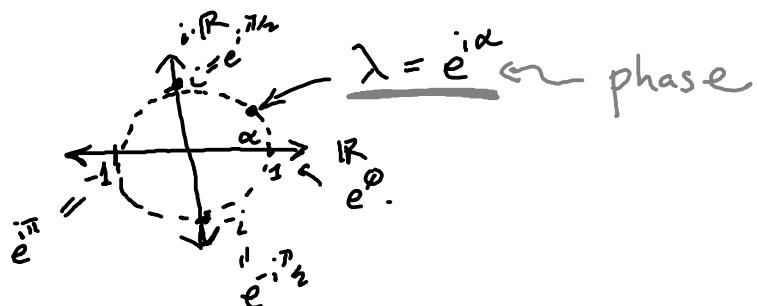
What about the converse?

* let λ be a complex # s.t. $\bar{\lambda} \cdot \lambda = 1$ (e.g. $\lambda = i$)

$$\text{double}(\lambda \cdot \begin{array}{c} \downarrow \\ \psi \end{array}) = \begin{array}{c} \downarrow \\ \psi \end{array} \bar{\lambda} \lambda \begin{array}{c} \downarrow \\ \psi \end{array} = \begin{array}{c} \downarrow \\ \psi \end{array} \begin{array}{c} \downarrow \\ \psi \end{array} = \text{double}\left(\begin{array}{c} \downarrow \\ \psi \end{array}\right)$$

"doubling eats numbers where $\bar{\lambda} \cdot \lambda = 1$."

* What numbers $\lambda \in \mathbb{C}$ satisfy $\bar{\lambda} \lambda = |\lambda|^2 = 1$?



global phase is redundant data.

$$\hat{\frac{1}{\Phi}} = \text{double} \left(\frac{1}{\Psi} \right) = \text{double} \left(e^{i\alpha} \cdot \frac{1}{\Psi} \right)$$

$$\text{Thm } \frac{1}{f} = \frac{1}{g} \iff \frac{1}{f} = e^{i\alpha} \frac{1}{g}$$

↑ global phase

Pf already showed \Leftarrow , so n.t.s. \Rightarrow . Let

$$\begin{array}{c} \overset{B}{\square} \overset{B}{\square} \\ f \quad f \\ \sqcup \quad \sqcup \\ \overset{B}{\square} \overset{B}{\square} \\ g \quad g \\ \sqcup \quad \sqcup \\ \end{array} = \begin{array}{c} \overset{B}{\square} \overset{B}{\square} \\ g \quad g \\ \sqcup \quad \sqcup \\ \end{array}.$$

and define:

$$\lambda := \begin{array}{c} \text{circle} \\ f \quad f \\ \sqcup \quad \sqcup \end{array} \quad \text{and} \quad \mu := \begin{array}{c} \text{circle} \\ g \quad g \\ \sqcup \quad \sqcup \end{array}$$

$$\bar{\lambda} \lambda = \begin{array}{c} \text{circle} \\ f \quad f \quad f \quad f \\ \sqcup \quad \sqcup \quad \sqcup \quad \sqcup \end{array} = \underbrace{\begin{array}{c} \text{circle} \\ f \quad g \quad g \quad f \\ \sqcup \quad \sqcup \quad \sqcup \quad \sqcup \end{array}}_{\bar{\mu} \mu} = \bar{\mu} \mu$$

$$2 \cdot \frac{1}{f} = \begin{array}{c} \text{circle} \\ f \quad f \quad f \\ \sqcup \quad \sqcup \quad \sqcup \end{array} = \begin{array}{c} \text{circle} \\ f \quad g \quad g \\ \sqcup \quad \sqcup \quad \sqcup \end{array} = \bar{\mu} \cdot \frac{1}{g}$$

$$\Rightarrow \frac{1}{f} = \underbrace{\frac{\bar{\mu}}{\lambda}}_{\text{for some } \lambda} \cdot \frac{1}{g} = e^{i\alpha} \cdot \frac{1}{g} \quad (\text{for some } \alpha). \quad \blacksquare$$

$$\left| \frac{\bar{\mu}}{\lambda} \right|^2 = \frac{\bar{\mu}}{\lambda} \cdot \frac{\mu}{\lambda} = \frac{\bar{\mu} \mu}{\lambda \lambda} = 1$$

Lecture 10

Pure quantum states $\left| \psi \right\rangle$ are for qubits

$$\left| \psi \right\rangle = \text{double} \left(a \left| \downarrow \right\rangle + b \left| \uparrow \right\rangle \right)$$

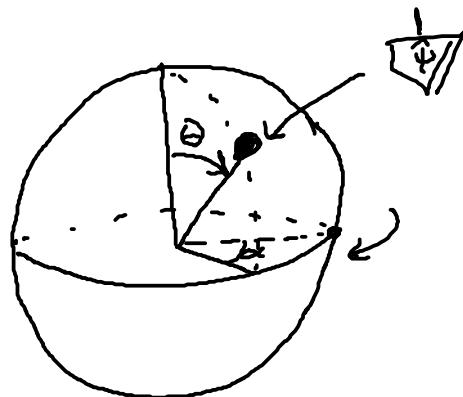
$x+iy$ $x'+iy'$

↓ ↓

$$\Psi \text{ normalised} \iff \underbrace{|a|^2 + |b|^2 = 1}_{\text{Cartesian}} \iff \begin{aligned} a &= \cos \frac{\theta}{2} \cdot e^{i\phi} \\ b &= \sin \frac{\theta}{2} \cdot e^{i\alpha} \end{aligned}$$

$\underbrace{\qquad\qquad\qquad}_{\text{polar}}$

$$\text{double} \left(\left| \psi \right\rangle \right) = \text{double} \left(e^{-i\phi} \cdot \left| \psi \right\rangle \right) = \text{double} \left(\cos \frac{\theta}{2} \left| \phi \right\rangle + e^{i\alpha} \sin \frac{\theta}{2} \left| \alpha \right\rangle \right)$$



Bloch sphere.

Things not preserved by doubling.

▷ numbers in the thing of pure quantum maps are positive real numbers, not \mathbb{C} .
(probabilities)

$$\boxed{\lambda} = \text{double}(\lambda) = \overline{\lambda} \cdot \lambda \geq 0.$$

▷ doubled ONB's are not bases!

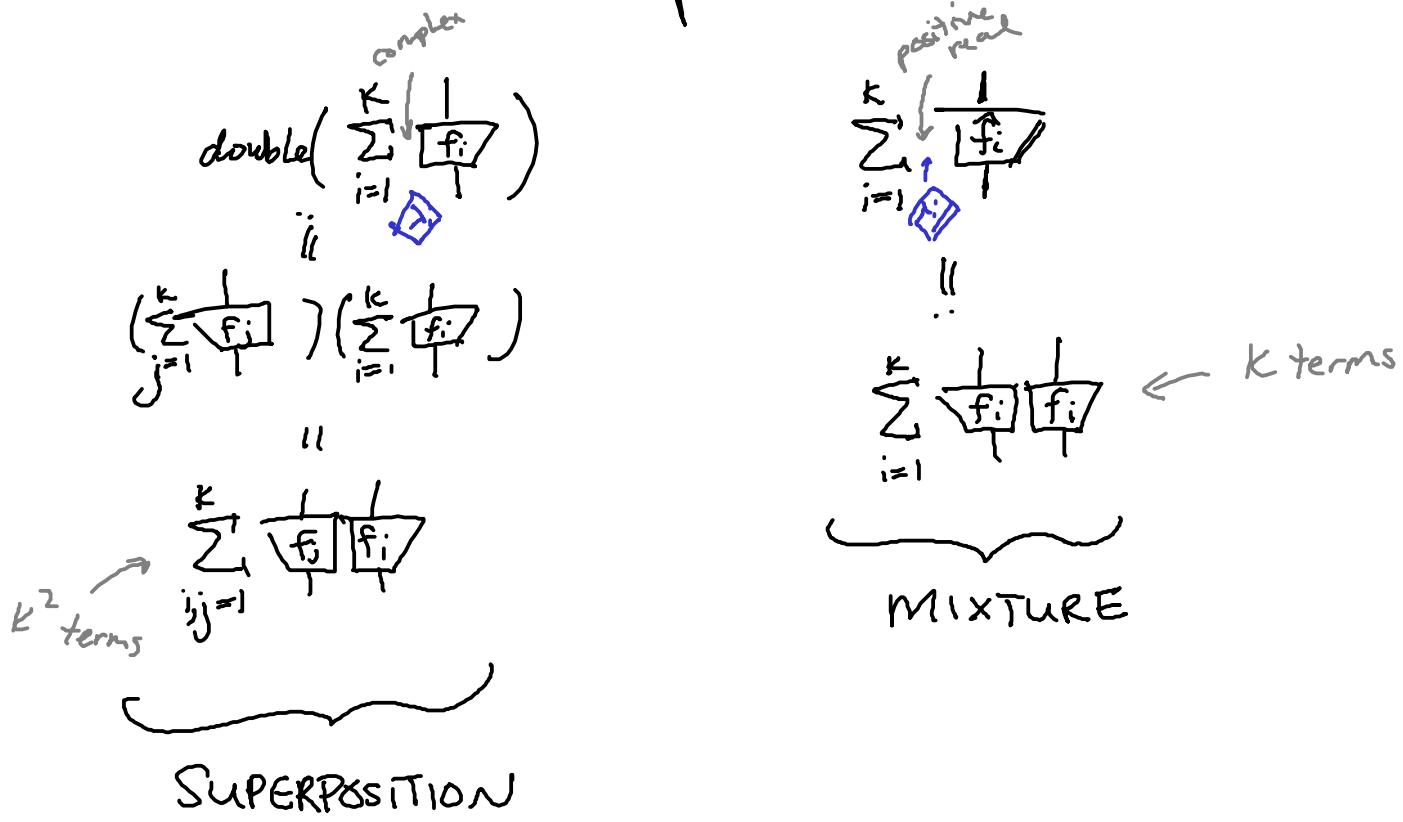
e.g. $\begin{array}{c} \uparrow \\ \boxed{\Psi} \end{array} = \sum_j \begin{array}{c} \downarrow \\ j \end{array}$ vs. $\begin{array}{c} \uparrow \\ \boxed{\phi} \end{array} = \sum_j e^{i\alpha_j} \begin{array}{c} \downarrow \\ j \end{array}$

in linear maps, $\begin{array}{c} \uparrow \\ \boxed{\psi} \end{array} = 1 \neq e^{i\alpha_i} = \begin{array}{c} \uparrow \\ \boxed{\phi} \end{array} \Rightarrow \begin{array}{c} \downarrow \\ \boxed{\psi} \end{array} \neq \begin{array}{c} \downarrow \\ \boxed{\phi} \end{array}$

in quantum maps, $\begin{array}{c} \uparrow \\ \boxed{\psi} \end{array} = 1 = 1 = \begin{array}{c} \uparrow \\ \boxed{\phi} \end{array}$, but $\begin{array}{c} \uparrow \\ \boxed{\psi} \end{array} \neq \begin{array}{c} \uparrow \\ \boxed{\phi} \end{array}$

The set $\{\begin{array}{c} \uparrow \\ i \end{array}\}_i$ describes an ONB measurement.

▷ Sums are not preserved.



6.2 Doubling "makes space" for an interesting new process.

DEF The discarding process is the effect that "discards" any normalised state $\hat{\Psi}$:

$$\text{disc.} \rightarrow \hat{\Psi} = \begin{cases} \vdots & \vdots \\ \psi & \vdash \\ \vdash & \vdash \\ \vdots & \vdots \end{cases}$$

n.b. for linear maps there is no such thing as discarding!

Pf Suppose $\exists \tilde{T}$ s.t.

$$\forall \psi : \begin{array}{c} \text{top} \\ \text{---} \\ \text{bottom} \end{array} = 1 \Rightarrow \begin{array}{c} \text{top} \\ \text{---} \\ \text{bottom} \end{array} = 1$$

$$\begin{aligned} \text{Then: } 1 &= \tilde{T} \circ \left[\frac{1}{\sqrt{2}} (\downarrow_0 + \downarrow_1) \right] \\ &= \frac{1}{\sqrt{2}} \left[\begin{array}{c} \text{top} \\ \text{---} \\ \text{bottom} \end{array} + \begin{array}{c} \text{top} \\ \text{---} \\ \text{bottom} \end{array} \right] = \frac{1}{\sqrt{2}} [1+1] = \sqrt{2} \quad \square \end{aligned}$$

Doubling to the rescue!

$$\begin{array}{c} \text{top} \\ \text{---} \\ \text{bottom} \end{array} = \begin{array}{c} \text{top} \\ \text{---} \\ \text{bottom} \end{array} = \begin{array}{c} \text{top} \\ \text{---} \\ \text{bottom} \end{array} = 1 \quad \dots$$

normalised

(n.b. $\begin{array}{c} \text{top} \\ \text{---} \\ \text{bottom} \end{array} = 1$ is uniquely fixed by $\begin{array}{c} \text{top} \\ \text{---} \\ \text{bottom} \end{array} = 1$)

Thm 6.31

$$\hat{\tilde{T}}_{\hat{A}} := \bigcap_{\hat{B}_n} \neq \hat{A} \cap_{\hat{A}} = \bigcap_{\hat{A}, \hat{B}_n}$$

In particular: $\tilde{\mathbb{T}}_{\hat{A} \otimes \hat{B}} = \tilde{\mathbb{T}}_{\hat{A}} \tilde{\mathbb{T}}_{\hat{B}} + \tilde{\mathbb{T}}_{\hat{I}} = \boxed{\quad}$

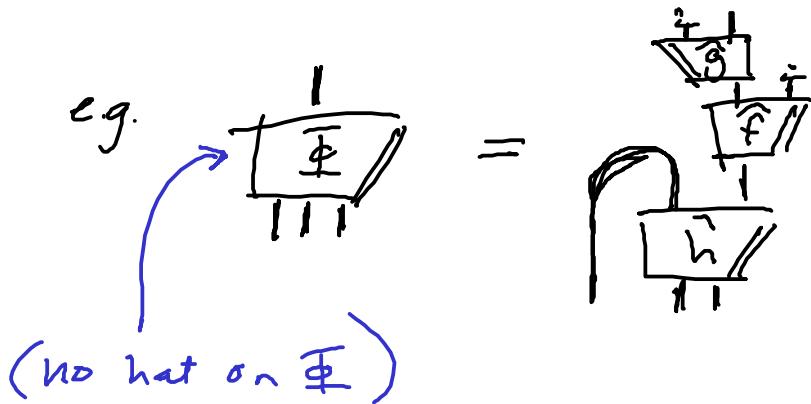
Discarding is not a pure q. map! \leftarrow Prop 6.27.

IDEA: $\cap \neq \text{double}(\triangle) = \triangle \uparrow \triangle \uparrow$

DEF The process theory of quantum maps has:

types: \hat{A} doubled Hilbert spaces

procs: ^{string} diagrams of $\boxed{\hat{f}^B_A}$ + $\tilde{\mathbb{T}}$.



CAUSALITY

DEF A process is called causal if:

$$\boxed{\Psi} = \dot{\tau}_{\hat{A}}$$

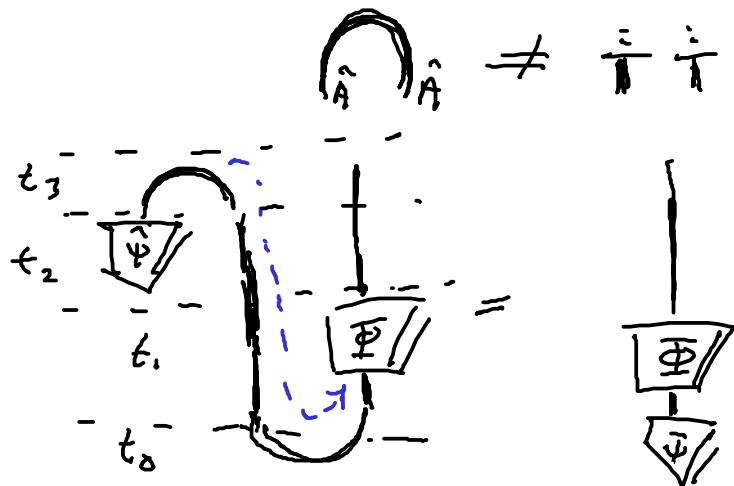
"If the output of a process is discarded, it doesn't matter which process occurred."

"A process can only influence its future."

In particular, causal effects:

$$\boxed{\Psi} = \dot{\tau}_{\hat{A}}$$

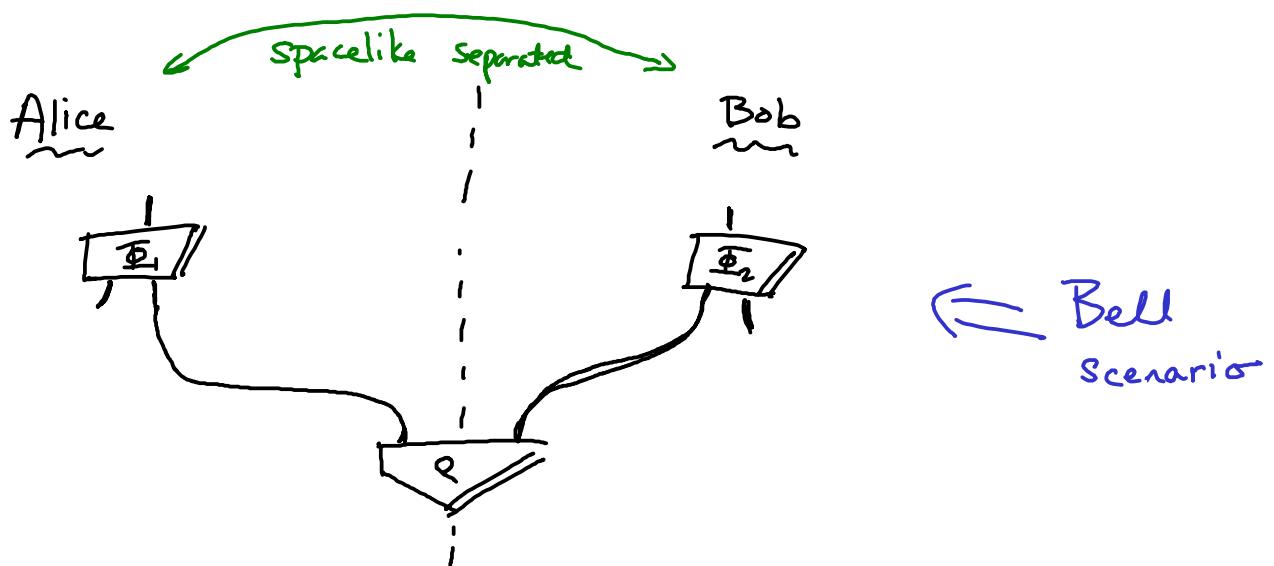
EXAMPLE OF A PROCESS WHICH IS NOT CAUSAL:



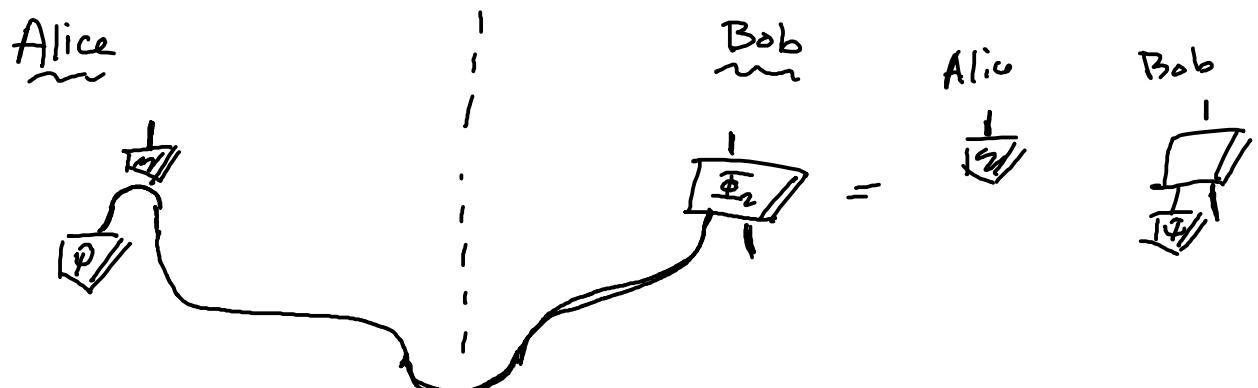
Lecture 11

6.3 Causality + non-signalling.

No signalling (special relativity)
travel faster than the speed of light." Nothing can



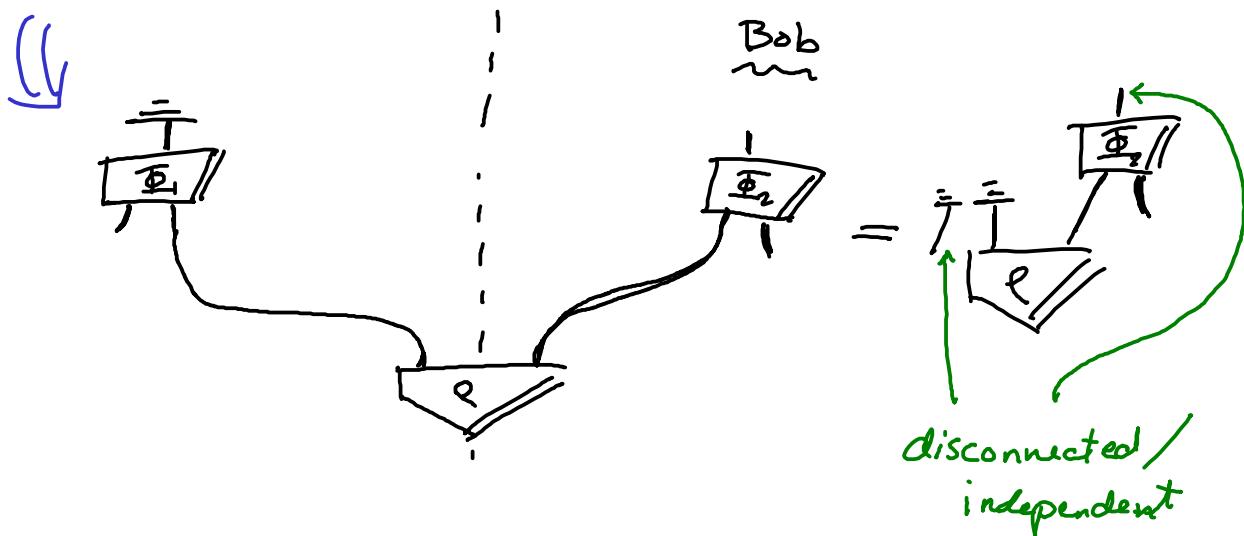
Spse .. processes didn't need to be causal.



Violates special relativity! X

Processes in "our universe" (as predicted by QT) do need to be causal.

From Bob's perspective, Bob can't access Alice's input.



\Rightarrow Bob cannot access Alice's input.

CAUSALITY \Leftrightarrow NO-SIGNALLING

* CAUSALITY RULES OUT \cap .

* So why did we learn about \cap ? ... and $\cap U = I$, etc.

"A process is $\overbrace{\text{deterministically physically realisable}}$ \Leftrightarrow it is causal."

6.4 (Non-deterministic) Quantum processes

Aside. Mixing.

NOTE: $\bar{T} = I = \sum_i f_i \otimes \bar{f}_i = \sum_i f_i \otimes f_i^* = \sum_i f_i f_i^*$

$$\text{So: } \begin{array}{c} \text{box with } \frac{1}{\sqrt{2}} \\ \text{and } \frac{1}{\sqrt{2}} \end{array} = \begin{array}{c} \text{box with } \frac{1}{\sqrt{2}} \\ \text{and } \frac{1}{\sqrt{2}} \end{array} = \boxed{\frac{1}{\sqrt{2}} f \otimes \bar{f}}$$

$$\boxed{\frac{1}{\sqrt{2}} f} = \boxed{\frac{1}{\sqrt{2}} f} = \sum_i \boxed{\frac{1}{\sqrt{2}} f_i} = \sum_i \boxed{\frac{1}{\sqrt{2}} f_i} \quad \text{MIXTURE}$$

\Rightarrow Disc. can be replaced by mixing.

Conversely:

$$\boxed{\frac{1}{\sqrt{2}} f} = \sum_i \boxed{\frac{1}{\sqrt{2}} f_i}$$

let $\boxed{\frac{1}{\sqrt{2}} f} := \sum_j \boxed{\frac{1}{\sqrt{2}} f_j}$, then: $\boxed{\frac{1}{\sqrt{2}} f} = \sum_j f_j \boxed{\frac{1}{\sqrt{2}} f_j} = \boxed{\frac{1}{\sqrt{2}} f}$

$$\text{So: } \boxed{\frac{1}{\sqrt{2}} f} = \boxed{\frac{1}{\sqrt{2}} f}$$

$$\boxed{\frac{1}{\sqrt{2}} f} = \sum_i \boxed{\frac{1}{\sqrt{2}} f_i} = \sum_i \boxed{\frac{1}{\sqrt{2}} f_i} = \boxed{\frac{1}{\sqrt{2}} f} = \boxed{\frac{1}{\sqrt{2}} f}$$

\Rightarrow mixture can be replaced by \bar{T} .

/ASSIDE ON MIXTURE... For Now.
(possibly non-deterministic)

DEF A quantum process is a set $\left\{ \frac{\tilde{\tau}}{\|\Phi_i\|} \right\}_i$ of quantum maps such that:

$$\sum_i \frac{\tilde{\tau}}{\|\Phi_i\|} = \tilde{\tau}$$

Ex1 CAUSAL QUANTUM MAPS. $\rightarrow \left\{ \frac{\tilde{\tau}}{\|\Phi\|} \right\}$. a.k.a.
deterministic quantum processes.

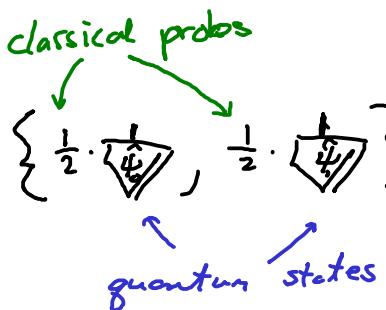
Ex2 Random state preparation,

* Alice flips a fair coin and prepares:

— $\frac{\tilde{\tau}}{\|\Phi\|}$ if heads.

— $\frac{\tilde{\tau}}{\|\Psi\|}$ if tails.

This is repr. by the quantum process: $\left\{ \frac{1}{2} \cdot \frac{\tilde{\tau}}{\|\Phi\|}, \frac{1}{2} \cdot \frac{\tilde{\tau}}{\|\Psi\|} \right\}$



Non-deterministic process \rightarrow we do know which branch happened once the process happens.

vs. a mixture \rightarrow we don't know which branch happened.

If Alice does $\{\frac{1}{2}|\psi_1\rangle\langle\psi_1|, \frac{1}{2}|\psi_2\rangle\langle\psi_2|\}$ and sends her state to Bob, but does not reveal the com flip, Bob sees a mixed state:

$$|\psi\rangle = \sum_i \frac{1}{2} |\psi_i\rangle$$

Ex3 Quantum ONB measurement.

$$\left\{ \begin{array}{c} \downarrow \\ i \end{array} \right\}_i \xrightarrow{\text{any ONB}} \left\{ \begin{array}{c} \nearrow \\ i \end{array} \right\}_i \xrightarrow{\text{adjoint}} \left\{ \begin{array}{c} \nearrow \\ i \\ \downarrow \end{array} \right\}_i \xrightarrow{\text{double}}$$

not an ONB

$$\sum_i \begin{array}{c} \nearrow \\ i \\ \downarrow \end{array} = \mathbb{I} \quad \checkmark \text{ quantum process}$$

So if Alice has a state $|\psi\rangle$, then she can measure $\left\{ \begin{array}{c} \nearrow \\ i \\ \downarrow \end{array} \right\}_i$.

$$\text{Prob}(i | \hat{\psi}) := \frac{1}{\langle \hat{\psi} | \hat{\psi} \rangle} \quad \text{Born rule.}$$

causal

$$\sum_i \text{Prob}(i | \hat{\psi}) = \sum_i \frac{1}{\langle \hat{\psi} | \hat{\psi} \rangle} = \frac{1}{\langle \hat{\psi} | \hat{\psi} \rangle} = 1 \quad \Rightarrow \{\text{Prob}(i | \hat{\psi})\}_i \text{ is a prob. distr.}$$

In Ex2, Bob's state was $\frac{1}{\sqrt{2}}|\psi_0\rangle + \frac{1}{\sqrt{2}}|\psi_1\rangle$.

So, Bob can also use the Born rule:

$$\begin{aligned} \text{Prob}(i|\rho) &= \langle \overset{i}{\triangle} | \rho | \overset{i}{\triangle} \rangle = \frac{1}{2} \cdot \langle \overset{i}{\triangle} | \rho | \overset{i}{\triangle} \rangle + \frac{1}{2} \cdot \langle \overset{i}{\triangle} | \rho | \overset{i}{\triangle} \rangle \\ &= \frac{1}{2} \cdot \text{Prob}(i|\hat{\psi}_0) + \frac{1}{2} \cdot \text{Prob}(i|\hat{\psi}_1) \end{aligned}$$

Eg. if $\langle \overset{i}{\triangle} | \rho | \overset{i}{\triangle} \rangle = |\overset{i}{\triangle}\rangle$ and $\langle \overset{i}{\triangle} | \rho | \overset{i}{\triangle} \rangle = \emptyset$, then:

$$\text{Prob}(i|\rho) = \frac{1}{2} \cdot \langle \overset{i}{\triangle} | \rho | \overset{i}{\triangle} \rangle + \frac{1}{2} \cdot \langle \overset{i}{\triangle} | \rho | \overset{i}{\triangle} \rangle = \frac{1}{2}.$$